

Fr 14.6.19

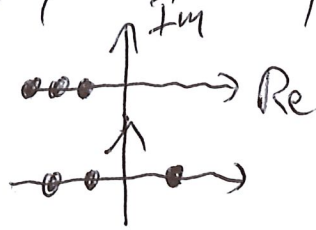
Lorentz dynamical system

①

$$FP_1: (0, 0, 0); \text{ char. Polyn. } P(\lambda) = (\lambda + 1)(\lambda^2 + (1+\sigma)\lambda + \sigma(1+r))$$

$$0 \leq r < 1$$

$$r > 1$$

All $\lambda_{1,2,3} < 0$ stableOne $\lambda_{1,2,3} > 0$ unstable

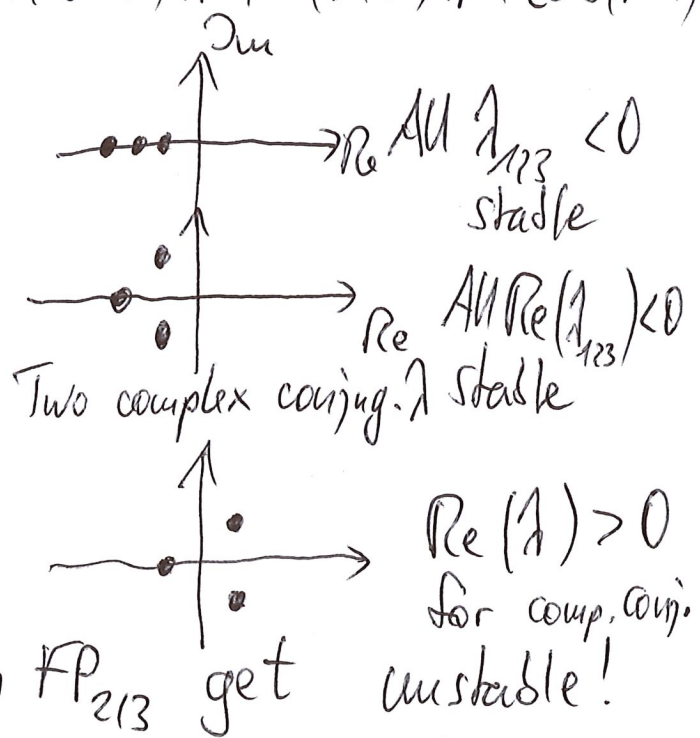
$$FP_{2/3}: (\pm \sqrt{b(r-1)}, \pm \sqrt{b(r-1)}, r-1) \quad \text{Only } r > 1$$

$$\text{char. Polyn. } -P(\lambda) = \lambda^3 + (1+b+\sigma)\lambda^2 + b(r+\sigma)\lambda + 2\sigma b(r-1)$$

$$1 < r < r_{\text{crit}1} = 1.346$$

$$r_{\text{crit}1} < r < r_{\text{crit}2} = 24.74$$

$$r > r_{\text{crit}2}$$

All $\lambda_{1,2,3} < 0$
stableTwo complex conj. λ stable
All $\text{Re}(\lambda_{1,2,3}) < 0$ both $FP_{2/3}$ get unstable!
 $\text{Re}(\lambda) > 0$
for comp. conj.

char. Polyn. ansatz: (Theorem of Vieta):

$$-P(\lambda) = (\lambda - \lambda_0)(\lambda - \lambda_r - i\lambda_i)(\lambda - \lambda_r + i\lambda_i)$$

with three EW $\lambda_1 = \lambda_0$; $\lambda_2 = \lambda_r + i\lambda_i$; $\lambda_3 = \lambda_r - i\lambda_i$

$$-P(\lambda) = \lambda^3 + \lambda^2(-\lambda_0 - 2\lambda_r) + \lambda[\lambda_r^2 + \lambda_i^2 + 2\lambda_0\lambda_r] - \lambda_0(\lambda_r^2 + \lambda_i^2)$$

Equate both forms of $P(\lambda)$; comp. coeff:

$$1 + b + \sigma = -\lambda_0 - 2\lambda_r$$

$$b(r + \sigma) = \lambda_r^2 + \lambda_i^2 + 2\lambda_0\lambda_r$$

$$2\sigma b(r - 1) = -\lambda_0(\lambda_r^2 + \lambda_i^2)$$

To find r_{crit2} we look for $\lambda_r = 0!$

$$2\sigma b(r - 1) = \underbrace{(1 + b + \sigma)}_{-\lambda_0} \underbrace{b(r + \sigma)}_{\lambda_i^2}$$

$$2\sigma r - 2\sigma = r + br + \sigma r + \sigma + b\sigma + \sigma^2$$

$$r(\sigma - b - 1) = \sigma(3 + b + \sigma)$$

$$r = r_{crit2} = \frac{\sigma(3 + b + \sigma)}{\sigma - b - 1} \sim 24.74 \dots$$

For $r < r_{crit2}$ FP_{213} stable

$r > r_{crit2}$ FP_{213} unstable

Volume contraction:

$$\frac{\partial}{\partial x} \dot{x} + \frac{\partial}{\partial y} \dot{y} + \frac{\partial}{\partial z} \dot{z} = -\sigma - 1 - b < 0$$

(Trace of Jacobian)

Globally Contracting! $\text{div}(\text{velocity}) < 0$
 $\dot{x}, \dot{y}, \dot{z}$

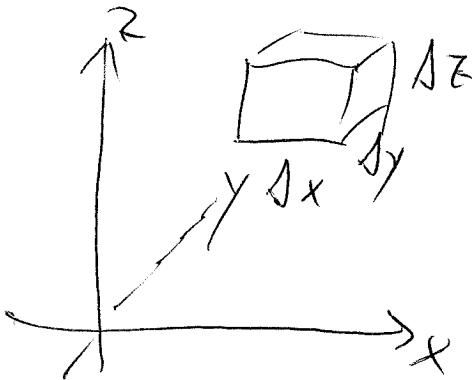
Definition of the "Attractor" A
(trajectory $\vec{x}(t)$ approaches A
and remains on it)

- if $\vec{x}_0 = \vec{x}(t_0) \in A$, $\vec{x}(t) \in A$ for all $t > t_0$
- (in some open environment of A
if $\vec{x}_0 = \vec{x}(t_0)$ starts), $\lim_{t \rightarrow \infty} \vec{x}(t) \in A$

stable FP and limit cycles are attractors

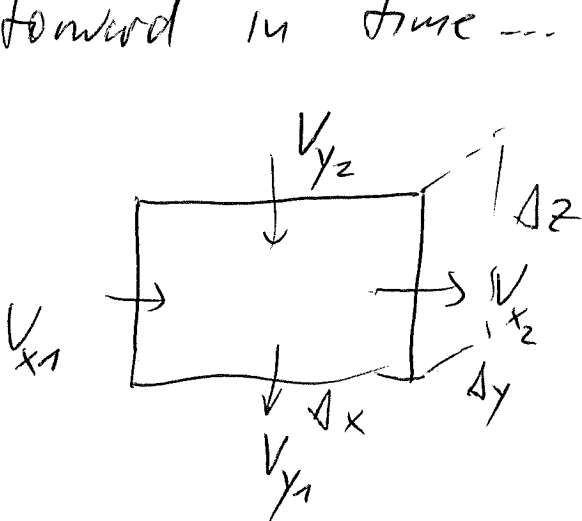
Wed 19.6.19 Lorentz dynamical system

Final Topics: (1) • Contraction



$$V = \Delta x \Delta y \Delta z$$

What happens to volume under Lorentz dynamical equation? (Every point integrated forward in time...)



$$\begin{aligned} \Delta V &= (\overbrace{v_{x2} - v_{x1}}^{\Delta v_x}) \Delta t \Delta y \Delta z \\ &+ (\overbrace{v_{y2} - v_{y1}}^{\Delta v_y}) \Delta t \Delta x \Delta z \\ &+ (\overbrace{v_{z2} - v_{z1}}^{\Delta v_z}) \Delta t \Delta y \Delta x \end{aligned}$$

Divide by $\Delta x \Delta y \Delta z \Delta t$; use $\Delta \log V = \frac{\Delta V}{V}$

$$\frac{\Delta \log V}{\Delta t} = \frac{\Delta v_x}{\Delta x} + \frac{\Delta v_y}{\Delta y} + \frac{\Delta v_z}{\Delta z}$$

Lim $\Delta x, \Delta y, \Delta z, \Delta t \rightarrow 0$

$$\frac{d \log V}{dt} = \frac{\partial v_x}{\partial x} + \frac{\partial v_y}{\partial y} + \frac{\partial v_z}{\partial z}$$

$v_x = x, v_y = y, v_z = z$

$$= -1 - 1 - 1 = -3 < 0!$$

(2) Dimension of the attractor A :

- $\dim A \neq 3$ (volume contraction)
- $\dim A \neq 2$ (curves would intersect,
more mathematically:
Theorem of Poincaré-Bendixson)
- $\dim A \neq 0$ (no stable FP)
- $\dim A \neq 1$? most difficult to show
could be limit cycle
exclude with theory of
discrete maps

Theorem of Poincaré - Bendixson (1D, 2D)

$$\dot{\vec{x}} = \vec{F}(\vec{x}) \qquad \begin{pmatrix} \dot{x} \\ \dot{y} \end{pmatrix} = \begin{pmatrix} F_1(x,y) \\ F_2(x,y) \end{pmatrix}$$

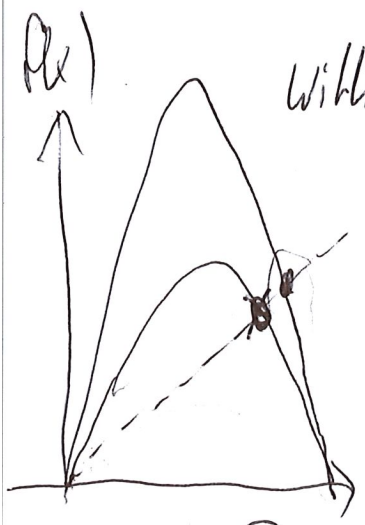
If trajectory $\vec{x}(t)$ remains within closed bounded region D for all $t \geq 0$:

$\Rightarrow \vec{x}(t)$ is

- (i) a closed orbit
- (ii) approach a periodic orbit (limit cycle)
- (iii) approach a fixed point

NB: if there is no ^{stable} fixed point,
 $\vec{x}(t)$ remains in $D \Rightarrow$ limit cycle

(4) Short Excursion Theory of discrete maps



• Set $x_0, x_1, \dots, x_n, x_{n+1}, \dots$
 with $x_{n+1} = f(x_n)$

• Fixed Points and stability
 $x^* = f(x^*)$

• Stability if $|f'(x^*)| < 1$



Example: Logistic Map, $c > 0$

$$x_{n+1} = c x_n (1 - x_n) = f(x_n)$$

$$(1 - 2x) \cdot c = c(1 - x) - cx = f'(x)$$

Fixed Point: $x = 1 - \frac{1}{c}$ $f'(x) = 2 - c$

Stable for: $1 < c < 3$

Iterated Map: $x_{n+1} = f^{(p)}(x_n) = f(f(\dots f(x_n)))$

Fixed Point of $f^{(p)}$ is p -periodic point of f

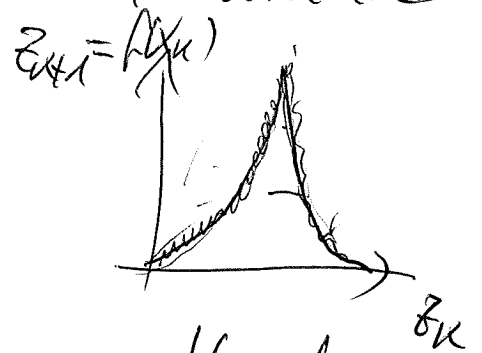
$$\dots x_n, x_{n+1}, \dots, x_{n+p-1}, \dots$$

$x_n = x_{n+p}$

Stability of FP of $f^{(p)}$: $|f^{(p)'}(x_{n+p})| = |f'(x_{n+p}) \dots f'(x_n)|$

Also x_{n+1}, x_{n+2}, \dots are FP of $f^{(p)}$, all stable or all unstable

Now look at $z_{k+1} = f(z_k)$ for Lorenz dynamical system as discrete map



Let us assume the Lorenz attractor has a ^{stable} p -periodic orbit. Then the discrete map $z_{k+p} = f^p(z_k)$

must have a stable fixed point. However, this is impossible, because

$$|f'(z_k)| > 1 \text{ everywhere!}$$

And therefore also

$$|f^p(z_k)| = \prod_{i=1}^p |f'(z_{k+i})| > 1!!$$

Therefore $\dim A \neq 1!$

Caveat: the map (the line) $f(z)$ is not steady! Also fractal ---