

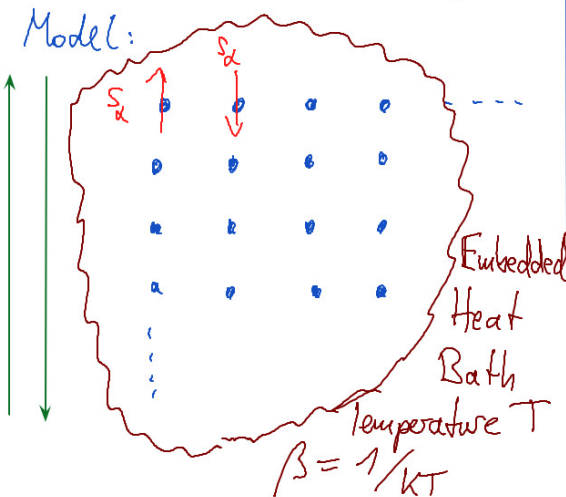
9.1.3. Ising Model

Ernst Ising, 1900 - 1998 Germany / UK / USA

Simple model of ferromagnetism \Leftrightarrow spontaneous ordering processes

Model:

Magnetic
Field
 B



- 2D mesh of Fe atoms $S_\alpha = \pm 1$
 $\alpha = 1, \dots$ number of atoms N
 $N = n \times n$ (in Tutorial: $n = 30$ $N = 900$)
- (method of scaling \Leftrightarrow dimensionless numbers)
- Every spin creates a magnetic moment
 $m = \pm 1 = S_\alpha$

We have in our tutorial $30 \times 30 = 900$ atoms

A spin state: $S = \{s_1, \dots, s_N\}$ describes the physical state of our system.

All possible states of the system: $S_i; i = 1, \dots, 2^N \sim 2^{900}$
(forms an ensemble) $\sim 8 \cdot 10^{270}$ states

2^{900} is a really crazy amount of states!

- Number of elementary particles in universe: $\sim 10^{80}$
- Number of Planck time (10^{-44} s) since "Big Bang" $\sim 10^{60}$

Fundamental Quantities:

- Energy of the system H (Hamilton function)

$$H(S_i) = -B \sum_{\alpha} S_{\alpha} - J \sum_{\langle \alpha, \beta \rangle} S_{\alpha} S_{\beta} \quad S_i = \{S_1, \dots, S_N\}_i$$

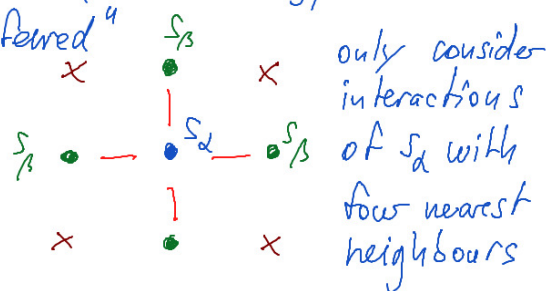
$\alpha = 1, \dots, N$

J : spin-spin interaction energy
 $J > 0$: ferromagnetic material, parallel spins are "preferred" (smaller energy)

$B > 0$: spins parallel to B "preferred"

$$\sum_{\langle \alpha, \beta \rangle} S_{\alpha} S_{\beta} = \sum_{\alpha} S_{\alpha} \left(\sum_{\beta} S_{\beta} \right)$$

β
4 nearest neighbours



- Magnetic Moment of the system:

$$M(S_i) = \sum_{\alpha} S_{\alpha}$$

(remember: every spin $S_{\alpha} = \pm 1$
corresponds to a magn. mom. $m = \pm 1$)

Note: $-N \leq M \leq N$

$M=0$: equal number of spin-up and -down
 $M=N$: all spins +1 - $M=-N$ all spins -1

- Magnetic moment per Fe atom (per spin):

$$m(S_i) = \frac{1}{N} M(S_i) \quad ; \quad -1 \leq m(S_i) \leq +1$$

Note: $m(S_i)$ can take any real value between $-1 \dots +1$

- Energy per Fe atom (per spin):

$$e(S_i) = \frac{1}{N} H(S_i)$$

Thermodynamic Quantities / Statistical Mechanics

- grand canonical ensemble (heat bath, not isolated)
- probability of state S , S_i ; $w(S)$, probability density function PDF $p(x)$
- $w(S) = \exp(-\beta H(S)) / \mathcal{Z}$

if $H(S)$ is "very" negative \Leftrightarrow probability is high ("preferred")
is "very" positive \Leftrightarrow " is small

Normalization factor \mathcal{Z} partition sum $\bullet \sum_{S_i} w(S_i) = 1$

$$\mathcal{Z} = \sum_{S_i} \exp(-\beta H(S_i)) \Rightarrow$$

$$\sum_{S_i} w(S_i) = \frac{1}{\mathcal{Z}} \underbrace{\sum_{S_i} \exp(-\beta H(S_i))}_{\mathcal{Z}} = 1$$

- Free Energy $F = -kT \log \mathcal{Z} = -\frac{1}{\beta} \log \mathcal{Z}$
- Internal Energy $U = kT^2 \frac{\partial \log \mathcal{Z}}{\partial T} = - \frac{\partial \log \mathcal{Z}}{\partial \beta}$
(equivalent to H)
 $\beta = 1/kT$
- • Mean Magnetization $M = - \frac{\partial F}{\partial B} = \frac{1}{\beta} \frac{\partial \log \mathcal{Z}}{\partial B}$

$$M = \frac{1}{\beta \mathcal{Z}} \cdot \frac{\partial \mathcal{Z}}{\partial B} = \frac{1}{\beta \mathcal{Z}} \frac{\partial}{\partial B} \sum_{S_i} \exp(-\beta H(S_i))$$

$$= \frac{1}{\mathcal{Z}} \sum_{S_i} \left(\sum_{\alpha} S_{\alpha} \right) \exp(-\beta H(S_i)) = \sum_{S_i} \underline{w(S_i)} \underline{M_i} =$$

Reminders: $H(S_i) = -B \sum_{\alpha} S_{\alpha} - J \sum_{(\alpha, \beta)} S_{\alpha} S_{\beta} \quad | \quad M_i = \sum_{\alpha} S_{\alpha}$

Average
over
ensemble!

$$M = \sum_{S_i} w(S_i) M_i$$

mean magnetization
average over all possible spin states
Weighted with known prob. $w(S_i)$

$$M_i = \sum_{\alpha} S_{\alpha}$$

One state!

M_i is magnetization of
one single ^{spin} state

$$U = - \frac{1}{Z} \frac{\partial}{\partial \beta} Z = - \frac{1}{Z} \frac{\partial}{\partial \beta} \sum_{S_i} \exp(-\beta H(S_i))$$

$$= \frac{1}{Z} \sum_{S_i} H(S_i) \exp(-\beta H(S_i))$$

$$= \sum_{S_i} w(S_i) H(S_i) \quad \text{with } H(S_i) = U(S_i)$$

Weighted
average of $H(S_i)$ over our ensemble!

$$M = \sum_{S_i} w(S_i) \underline{M_i} ; U = \sum_{S_i} w(S_i) \underline{H(S_i)}$$

- Tutorial:
- start with random spin state $S_i \rightarrow w(S_i) \cdot Z$
 - choose a new state by flipping one spin \rightarrow accept or reject Metropolis algorithm
- in Metropolis we need: $w(S_{i+1}) / w(S_i)$
- sweeps : N_{spin} times spin flipping
-

$$\frac{H}{kT} = \beta H ; h = \beta H ; b = \beta B ; j = \beta J$$

$$h = -b \sum_{\alpha} S_{\alpha} - j \sum_{\langle \alpha \beta \rangle} S_{\alpha} S_{\beta} \quad w(S_i) = \frac{1}{Z} \exp(-h(S_i))$$

Varying in experiment: • b magnetic field • j temperature ($j = J / kT$)

For $j < j_{\text{crit}}$ no order

For $j > j_{\text{crit}}$ spontaneous order \rightarrow ferromagnetism

• Approximate Model: mean field model $j_{\text{crit}} = 0.25$ not real!

• Real case, our experiment: $j_{\text{crit}} = 0.4406868 \dots$

Onsager 1944, $\sinh(2j_{\text{crit}}) = 1$

$$\sinh^{-1}(1) = 2j_{\text{crit}} = \ln(1 + \sqrt{2})$$

$$\sinh^{-1}(x) = \ln(x + \sqrt{x^2 + 1}) \quad \Rightarrow \quad j_{\text{crit}} = \frac{\ln(1 + \sqrt{2})}{2} \\ \sim 0.4406868 \dots$$