



Dynamics of nuclear and globular clusters, relativistic dynamics of compact objects

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Jan. 8, 2025

ISGC Seminar

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<https://astro-silkroad.eu>



- 1) Introduction, Globular Star Clusters**
- 2) Nuclear Star Clusters
- 3) Code(s) and Hardware

University of Heidelberg: ARI/ZAH, connected with IMPRS

Every year graduate school application deadline (November)

<https://zah.uni-heidelberg.de/institutes>

<https://www.imprs-hd.mpg.de/>

The **Zentrum für Astronomie der Universität Heidelberg (ZAH)** was established in 2005 by joining the **Astronomisches Rechen-Institut (ARI)**, the **Landessternwarte Königstuhl (LSW)** and the **Institut für Theoretische Astrophysik (ITA)**.

Astronomisches Rechen-Institut (ARI)



Mönchhofstr. 12-14
69120 Heidelberg
Tel: 06221 54 1801

Institut für Theoretische Astrophysik (ITA)



Albert-Ueberle-Str. 2 and
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International Max Planck Research School for
Astronomy & Cosmic Physics at the University of Heidelberg



Astronomisches
Rechen-Institut (ARI) at
U.f Heidelberg;
founded May 10, 1700



Siemens 2002
Computer in 1964
At ARI



Kavli Institute for Astronomy and Astrophysics, Peking Univ.



Total ~25 postdocs; 2/3 are non-Chinese.

Founded 2008
Regular Openings:

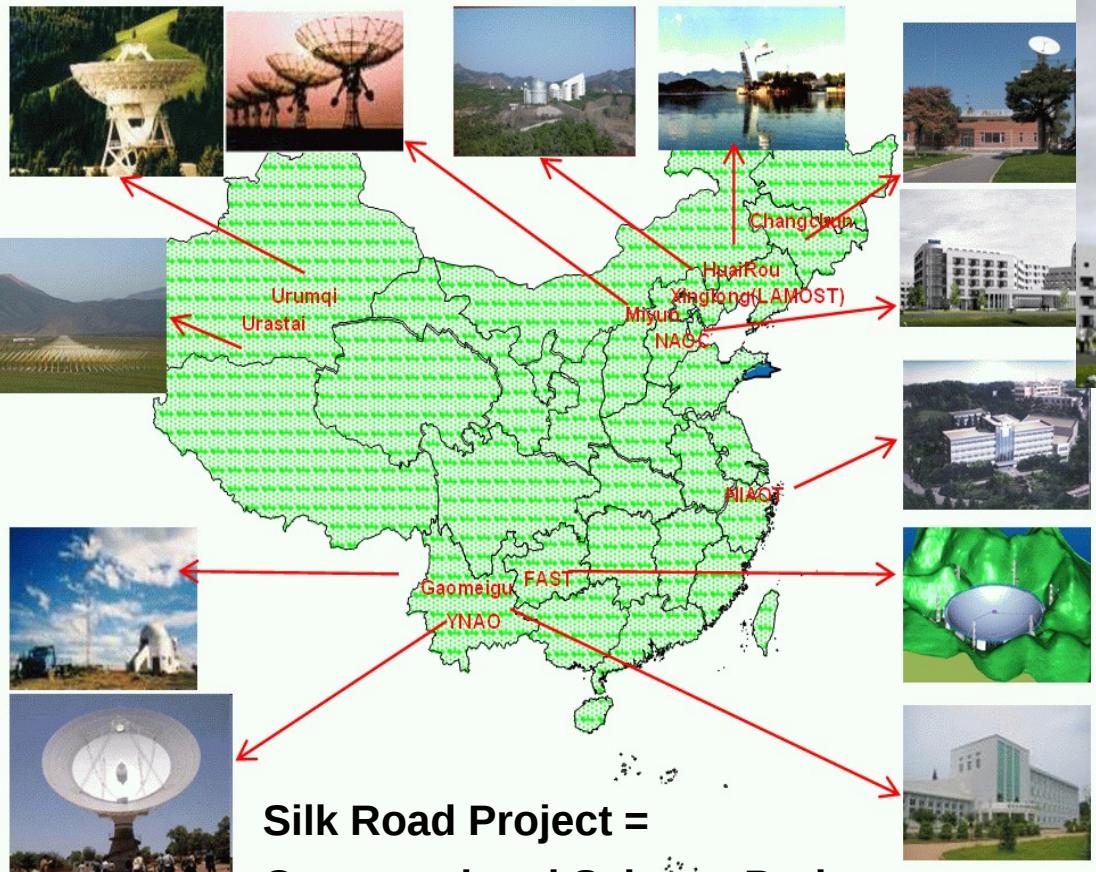
Postdoc
Faculty
Visitors
(see AAS Job Reg,
search KIAA)



中国科学院国家天文台

NATIONAL ASTRONOMICAL OBSERVATORIES , CHINESE ACADEMY OF SCIENCES

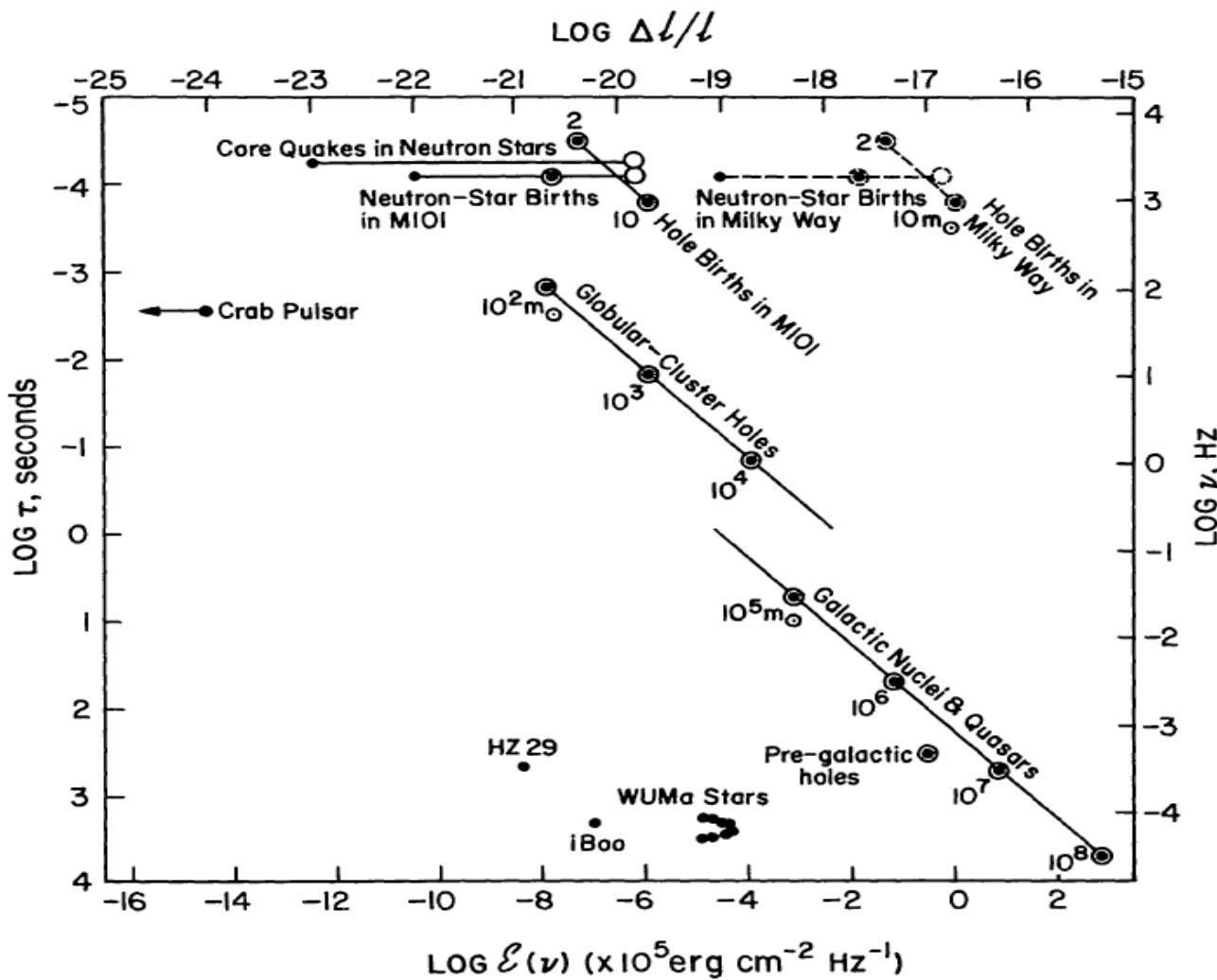
NAOC/
CAS



Top: NAOC Headquarter Beijing
Bottom: LAMOST Site

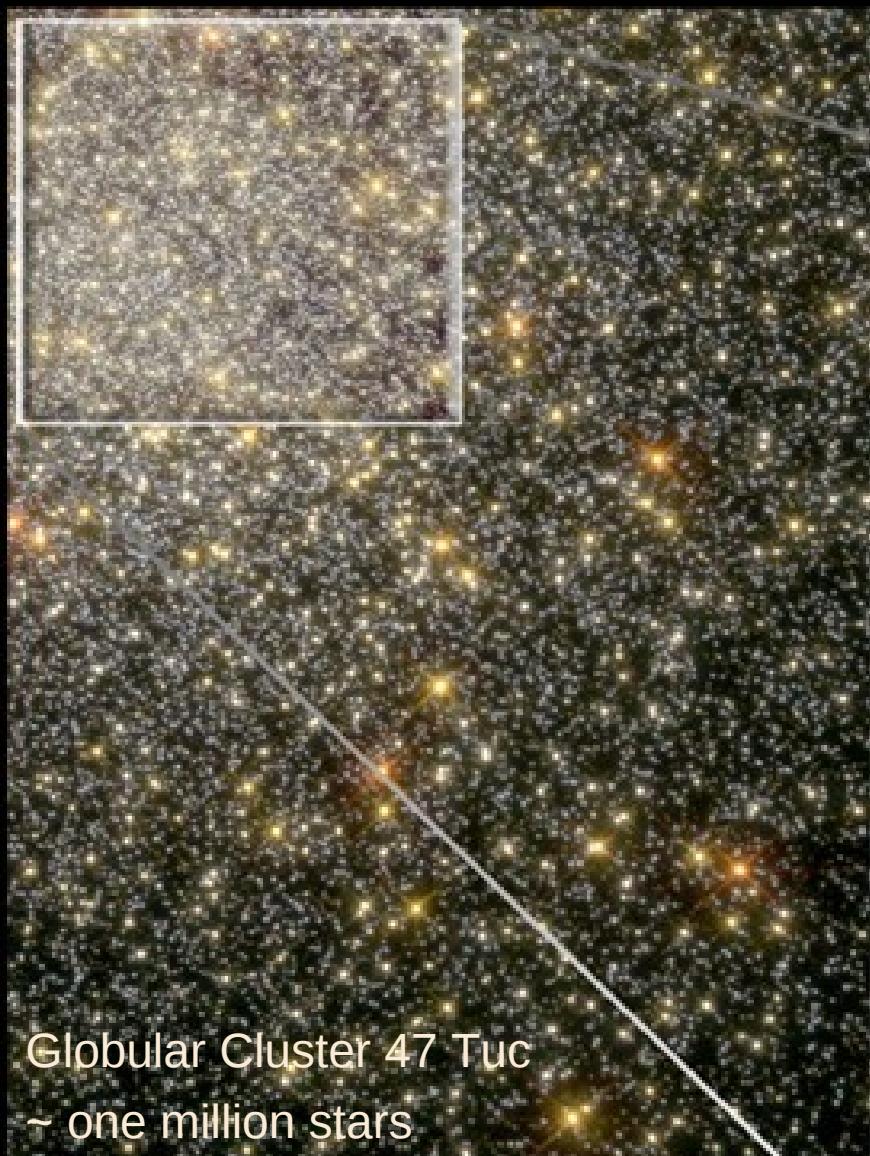


Sources In Luminosity



Tyson
&
Giffard
Ann. Rev.
Astr. Astroph.
1978 !!

Figure 2 Numerical estimates of the strongest gravitational radiation signals that might be reaching the earth (from Thorne 1977). The strength of each signal is plotted as a function of its dominant spectral characteristics. Continuous monochromatic signals are denoted by dots corresponding to their dimensionless strain amplitudes. Dots surrounded by circles indicate the strain and spectral energy density of broad-band bursts with frequency ranging from 0.5ν to 1.5ν and duration $\tau \sim \frac{1}{2}\nu$. A damped ringing wave is denoted by a dot at the appropriate frequency giving the amplitude and an open-circle giving the total energy spectral density. With the exception of those connected by dotted lines, the distances chosen would result in a few events per year of the strength shown.

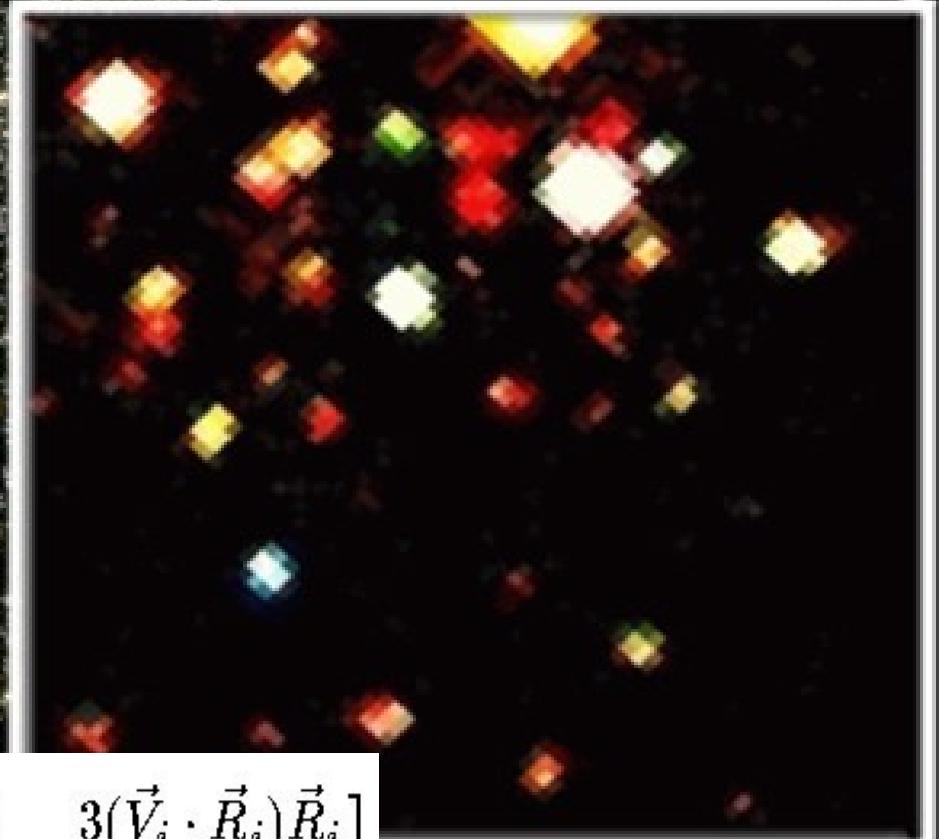


Globular Cluster 47 Tuc
~ one million stars

(Credit: X-ray: NASA/CfA/J. Grindlay et al.,
Optical: NASA/STScI/R. Gilliland et al.)

X-ray binaries

with neutron stars
and black holes



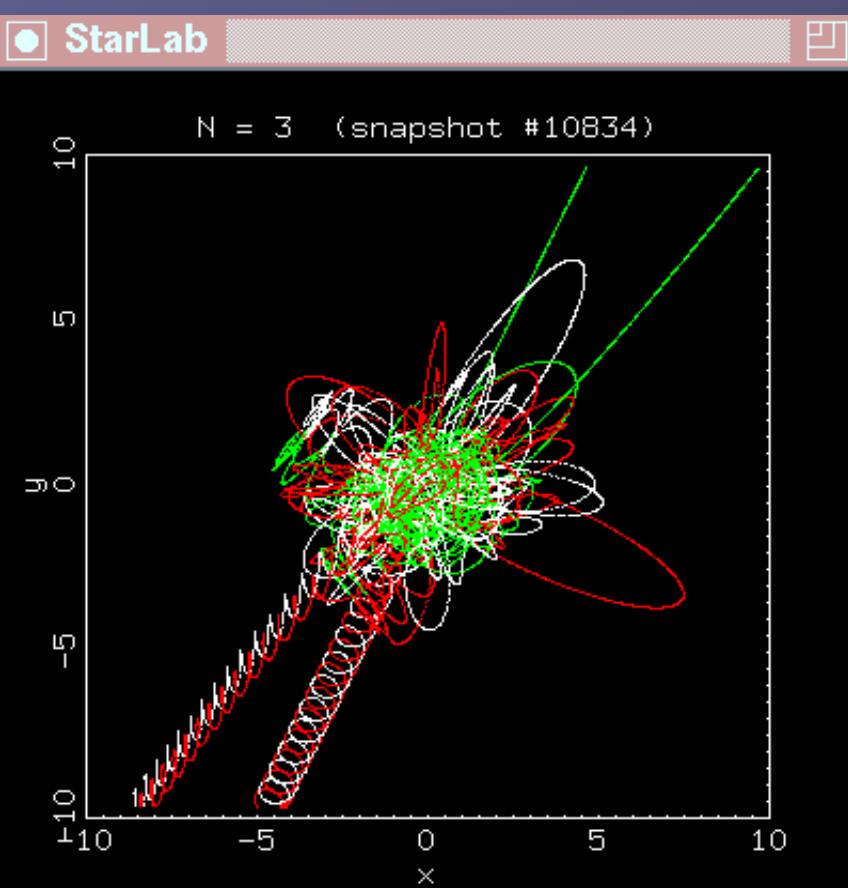
$$\vec{a}_0 = \sum_j G m_j \frac{\vec{R}_j}{R_j^3} ; \quad \vec{a}_0 = \sum_j G m_j \left[\frac{\vec{V}_j}{R_j^3} - \frac{3(\vec{V}_j \cdot \vec{R}_j)\vec{R}_j}{R_j^5} \right]$$

Star Clusters: Modelling the Dynamics

(Manage deterministic chaos)

Chaos in 3-body Encounters
Starlab Simulation (S.L.W. McMillan)

<http://www.physics.drexel.edu/~steve/>



Gravothermal Oscillations - Attractor in Phase Space

Spurzem 1994, Giersz & Spurzem 1994
Amaro-Seoane, Freitag & Sp. 2004

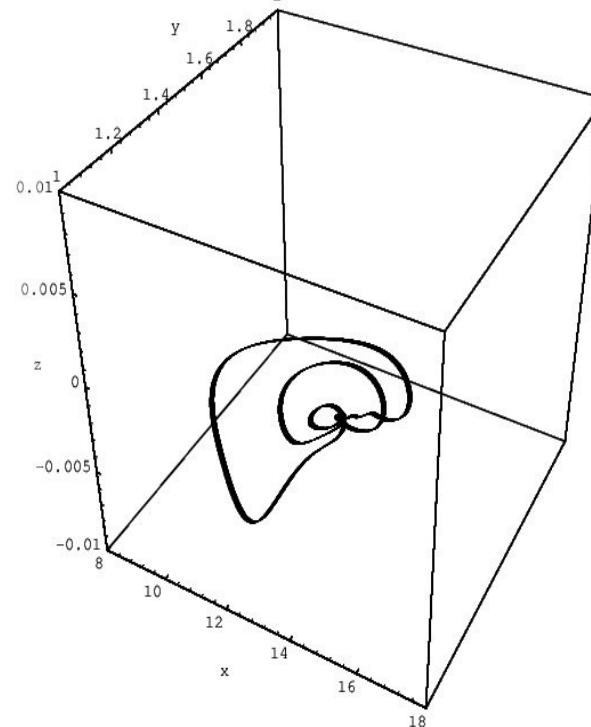


Fig. 3:
Projected three-dimensional attractor for $N = 100,000$ system, $x = \log \rho'_c$, $y = \log \sigma'_c$, $z = \xi$.

Star Clusters: Modelling the Dynamics

(coupling of scales)

$$\vec{a}_0 = \sum_j G m_j \frac{\vec{R}_j}{R_j^3} ; \quad \ddot{\vec{a}}_0 = \sum_j G m_j \left[\frac{\vec{V}_j}{R_j^3} - \frac{3(\vec{V}_j \cdot \vec{R}_j)\vec{R}_j}{R_j^5} \right]$$

- $N = \infty$

negative specific Heat

gravothermal Collapse

gravothermal Oscillations

- $N = 3$ ($N = 2, \dots, \approx 100$)

History

Exponential Instability

Chaos and Resonance

Regularisation

- $N = 10^6$ ($N = 10^4, 10^5$)

Post-Kollaps-Evolution

Binaries

Globular Clusters

Star Clusters: Modelling the Dynamics

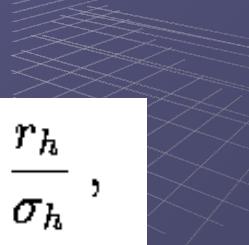
(multiple time scales)

● Millisecond Pulsars

● Dynamical Time Scale

● Relaxation Time Scale

● Age of Universe



$$t_{\text{cr}} = \frac{r_h}{\sigma_h} , \quad 10^{-3} \text{ sec}$$

$$t_{\text{rx}} = \frac{9}{16\sqrt{\pi}} \frac{\sigma^3}{G^2 m \rho \ln(\gamma N)} . \quad 10^6 \text{ yrs}$$

$$10^8 \text{ yrs}^{**}$$

$$10^{10} \text{ yrs}$$

Laboratories for gravothermal N-Body Systems!

Note: Cosmological and Galactic N-Body Simulations need few crossing times, and less than a relaxation time, while gravothermal systems need multiples of N crossing times, several Complexity goes approx $N^{**2.3}$!

$$t_{\text{cr}} \approx \sqrt{\frac{r_h^3}{GM_h}} .$$

← Virial Equilibrium →

$$\frac{t_{\text{rx}}}{t_{\text{dyn}}} \propto \frac{N}{\log(\gamma N)} .$$

Relaxation time definitions: Here: Larson, R.B. (1970) – theoretically sound, also see Spurzem & Takahashi (1995); in N-Body Code use: Eq. (8-71) of Binney/Tremaine (BT) Galactic Dynamics (from Spitzer & Hart 1971)

(Kustaanheimo and Stiefel 1965). We modified this scheme to allow for relativistic corrections to the Newtonian forces by expanding the acceleration in a series of powers of $1/c$ in the following way (Damour & Deruelle 1981; Soffel 1989):

$$\underline{a} = \underbrace{\underline{a}_0}_{\text{Newt.}} + \underbrace{c^{-2}\underline{a}_2}_{1\mathcal{PN}} + \underbrace{c^{-4}\underline{a}_4}_{2\mathcal{PN}} + \underbrace{c^{-5}\underline{a}_5}_{2.5\mathcal{PN}} + \mathcal{O}(c^{-6}), \quad (1)$$

periast. shift grav. rad.

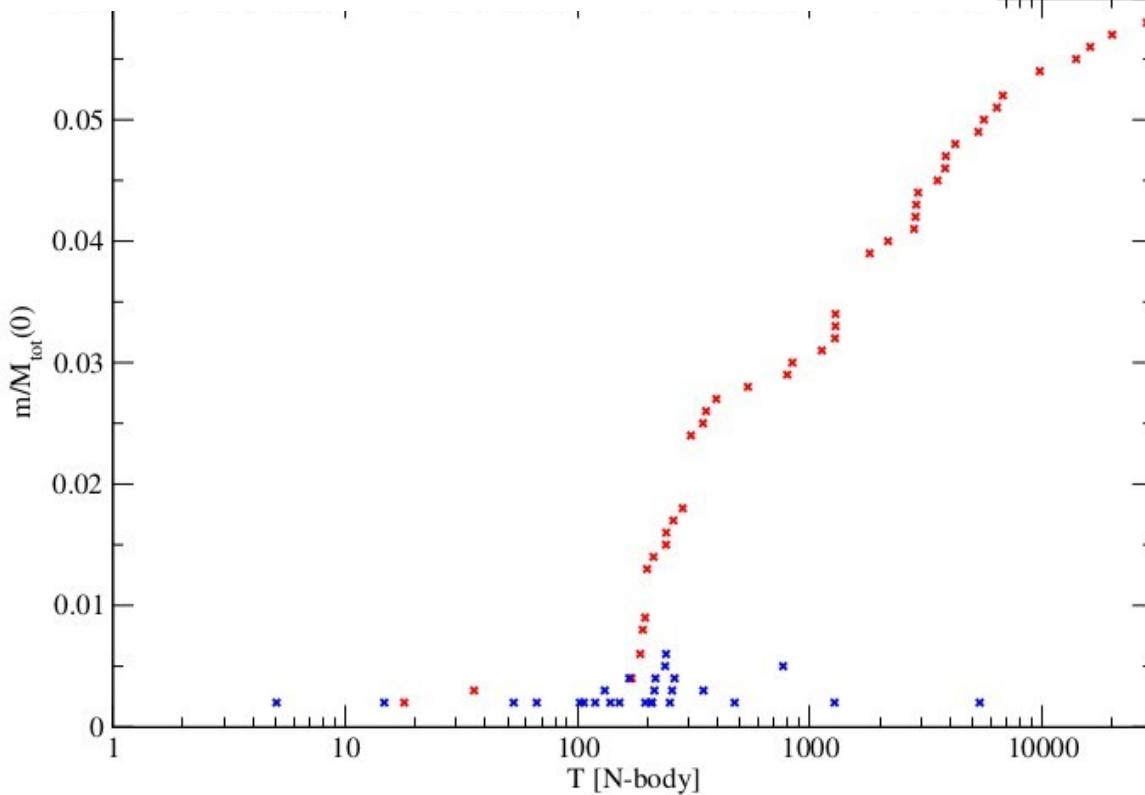


Figure 1. Time evolution of merging masses. The formation of the runaway particle is about the time of the cluster core collapse. For more details see text

History: Runaway Merger Of Black Holes Cluster

Kupi, Amaro-Seoane,
Spurzem.. 2006

“Perturbation”
In two-body
KS regularization

Switch on if:

Perihel Shift > 10⁻⁴

AND

v/c > 10-3

Black Holes were retained in globular clusters simulations (and formed binaries, gravitational wave emission predicted 2010):

- Before Strader et al. detection
- Before Breen & Heggie
- Before LIGO detection

Downing 2012, Downing, Benacquista, Giersz & Spurzem, 2010, 2011:
(see also Banerjee, Baumgardt & Kroupa 2010 but ...)



Compact Binaries in Star Clusters I - Black Hole Binaries Inside Globular Clusters

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⁵Nicolaus Copernicus Astronomical Center, Polish Academy of Sciences, ul. Bartycka 18, 00-716 Warsaw, Poland

Compact Binaries in Star Clusters II - Escapers and Detection Rates

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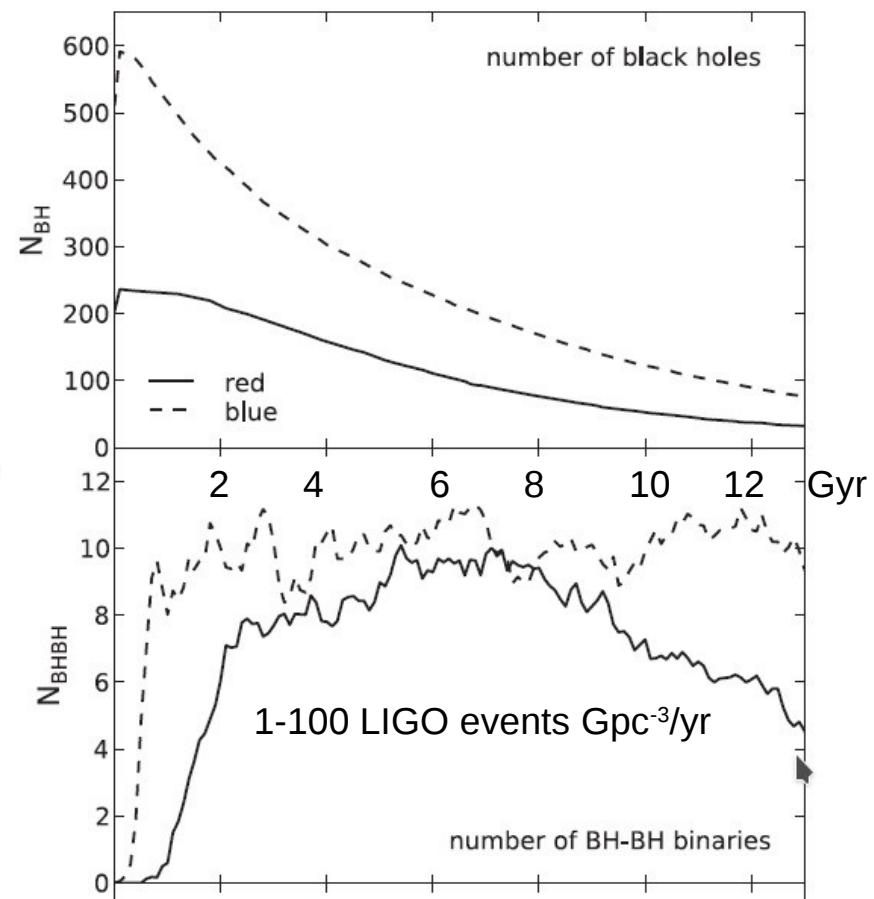
²Fellow of the International Max-Planck Research School for Astronomy and Cosmic Physics at the University of Heidelberg, Heidelberg, Germany

³Center for Gravitational Wave Astronomy, University of Texas at Brownsville, Brownsville, TX 78520, USA

⁴Nicolaus Copernicus Astronomical Center, Polish Academy of Sciences, ul. Bartycka 18, 00-716 Warsaw, Poland

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DRAGON I Simulation

DRAGON Simulation

<https://astro-silkroad.eu>
<https://github.com/nbody6ppgpu>
Also in: <https://www.punch4nfdi.de/>

One million stars direct simulation,

biggest and most realistic direct N-Body simulation of globular star clusters.

With stellar mass function, single and binary stellar evolution, regularization of close encounters, tidal field (NBODY6++GPU).

(NAOC/Silk Road/MPA collaboration).

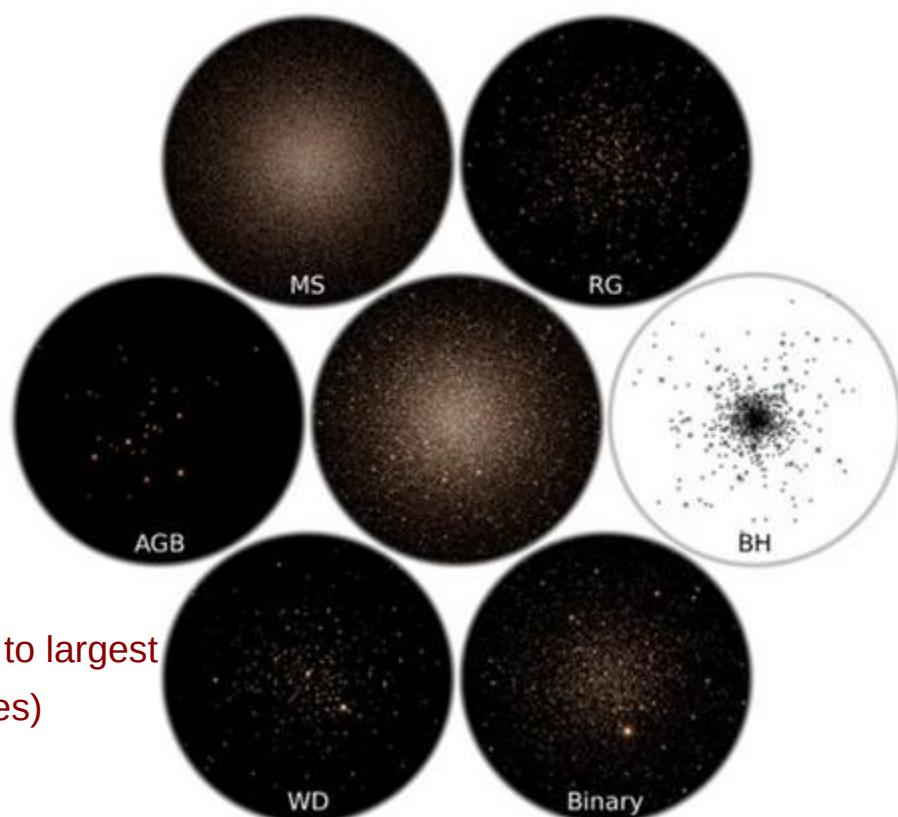
Wang, Spurzem, Aarseth, Naab et al.

MNRAS, 2015

Wang, Spurzem, Aarseth Naab, et al.

MNRAS 2016

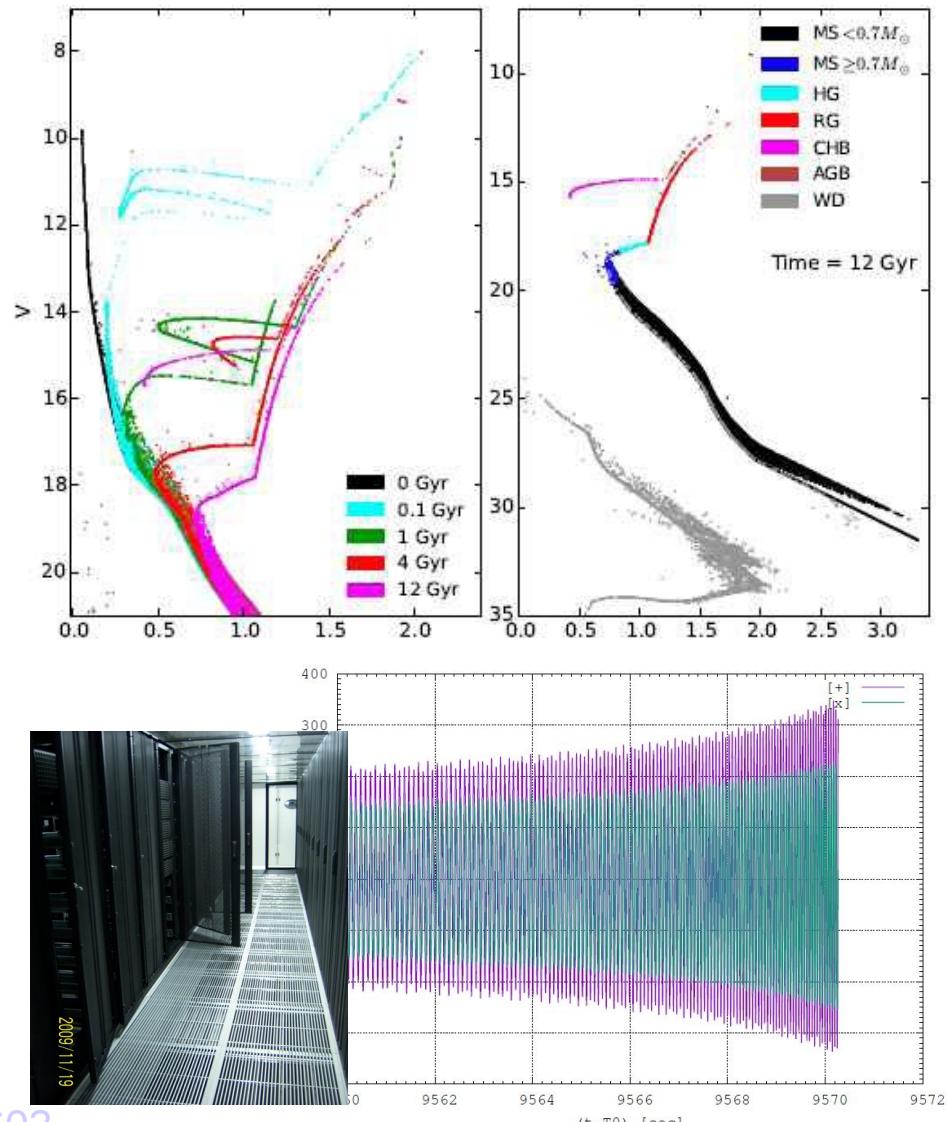
Number of Floating Point Operations (~1M bodies) similar to largest Cosmological simulations (Millennium, Illustris, ~ 10^{10} bodies)



天龙星团模拟：百万数量级恒星、黑洞和引力波

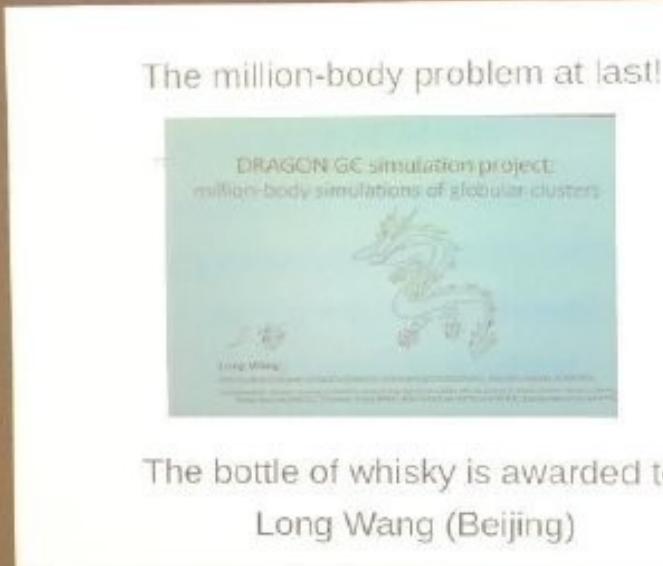
Dragon Star Cluster Simulations: Millions of Stars; black holes and gravitational waves

- First realistic globular star cluster model with million stars (*Wang, Spurzem, Aarseth, ..., Berczik, Kouwenhoven, ... MNRAS 2015, 2016*)
- Synthetic CMD (right side) with zero photometric errors, different ages shown
- Black hole binary mergers occur as observed by LIGO. Our grav. waveforms computed from simulation (right side). (Only inspiral plotted not ringdown.)
- GPU accelerated supercomputers Iaohu in NAOC and hydra of Max-Planck (MPCDF) in Germany needed!



CPU/GPU N-body6++

Long Wang, Ph.D. Peking University 2016:
Million-Body Award by MODEST community
And IAU Ph.D. prize



Key Question 1. When will we see the first star-by-star N -body model of a globular cluster?

- Honest N-body simulation
- Reasonable mass at 12 Gyr ($\sim 5 \times 10^4 M_\odot$)
- Reasonable tide (circular galactic orbit will do)
- Reasonable IMF (e.g. Kroupa)
- Reasonable binary fraction (a few percent)
- Any initial model you like (Plummer will do)
- A submitted paper (astro-ph will do)

An inducement: a bottle of single malt Scotch whisky worth €50



What about black holes?
Black Hole Binaries?

Post-Newtonian Dynamics

$r; v$: relative distance, velocity ; $\mu = m_1 m_2 / M$: reduced mass ($M = m_1 + m_2$)

$v = \mu / M$: mass ratio; $\mathbf{n} = \mathbf{r} / r$: unit vector in radial direction

$$\frac{dv^i}{dt} = -\frac{Gm}{r^2} [(1 + \mathcal{A}) n^i + \mathcal{B} v^i] + \mathcal{O}\left(\frac{1}{c^8}\right), \quad (181)$$

and find [43] that the coefficients \mathcal{A} and \mathcal{B} are

$$\begin{aligned} \mathcal{A} = & \frac{1}{c^2} \left\{ -\frac{3\dot{r}^2\nu}{2} + v^2 + 3\nu v^2 - \frac{Gm}{r} (4 + 2\nu) \right\} && \text{Perihel shift} \\ & + \frac{1}{c^4} \left\{ \frac{15\dot{r}^4\nu}{8} - \frac{45\dot{r}^4\nu^2}{8} - \frac{9\dot{r}^2\nu v^2}{2} + 6\dot{r}^2\nu^2 v^2 + 3\nu v^4 - 4\nu^2 v^4 \right. \\ & \quad \left. + \frac{Gm}{r} \left(-2\dot{r}^2 - 25\dot{r}^2\nu - 2\dot{r}^2\nu^2 - \frac{13\nu v^2}{2} + 2\nu^2 v^2 \right) + \frac{G^2 m^2}{r^2} \left(9 + \frac{87\nu}{4} \right) \right\} \\ & + \frac{1}{c^5} \left\{ -\frac{24\dot{r}\nu v^2}{5} \frac{Gm}{r} - \frac{136\dot{r}\nu}{15} \frac{G^2 m^2}{r^2} \right\} \end{aligned}$$

$$\begin{aligned} \mathcal{B} = & \frac{1}{c^2} \{-4\dot{r} + 2\dot{r}\nu\} && \text{Grav. Radiation} \\ & + \frac{1}{c^4} \left\{ \frac{9\dot{r}^3\nu}{2} + 3\dot{r}^3\nu^2 - \frac{15\dot{r}\nu v^2}{2} - 2\dot{r}\nu^2 v^2 + \frac{Gm}{r} \left(2\dot{r} + \frac{41\dot{r}\nu}{2} + 4\dot{r}\nu^2 \right) \right\} \\ & + \frac{1}{c^5} \left\{ \frac{8\nu v^2}{5} \frac{Gm}{r} + \frac{24\nu}{5} \frac{G^2 m^2}{r^2} \right\} \end{aligned}$$

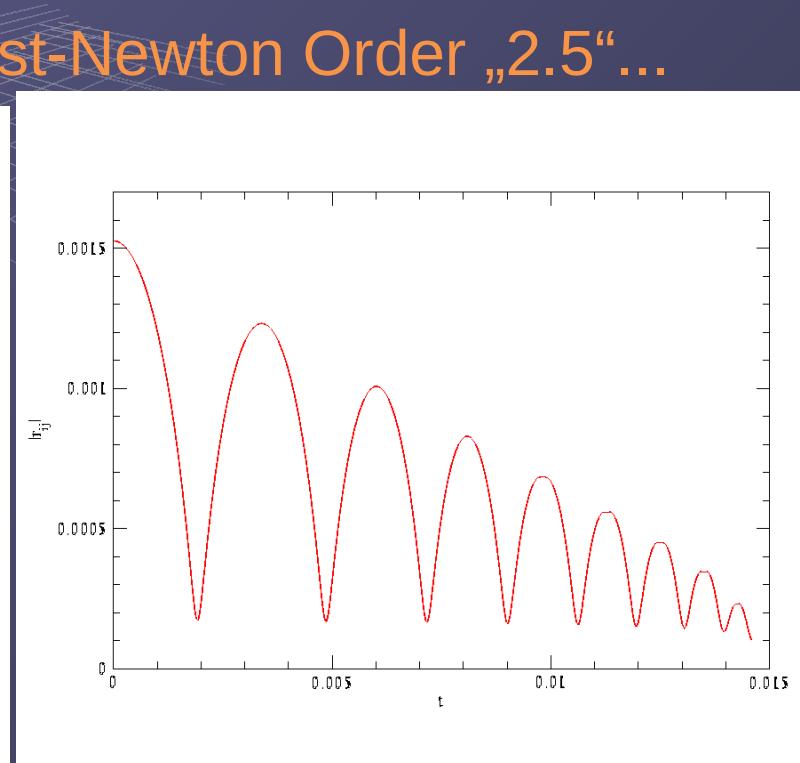
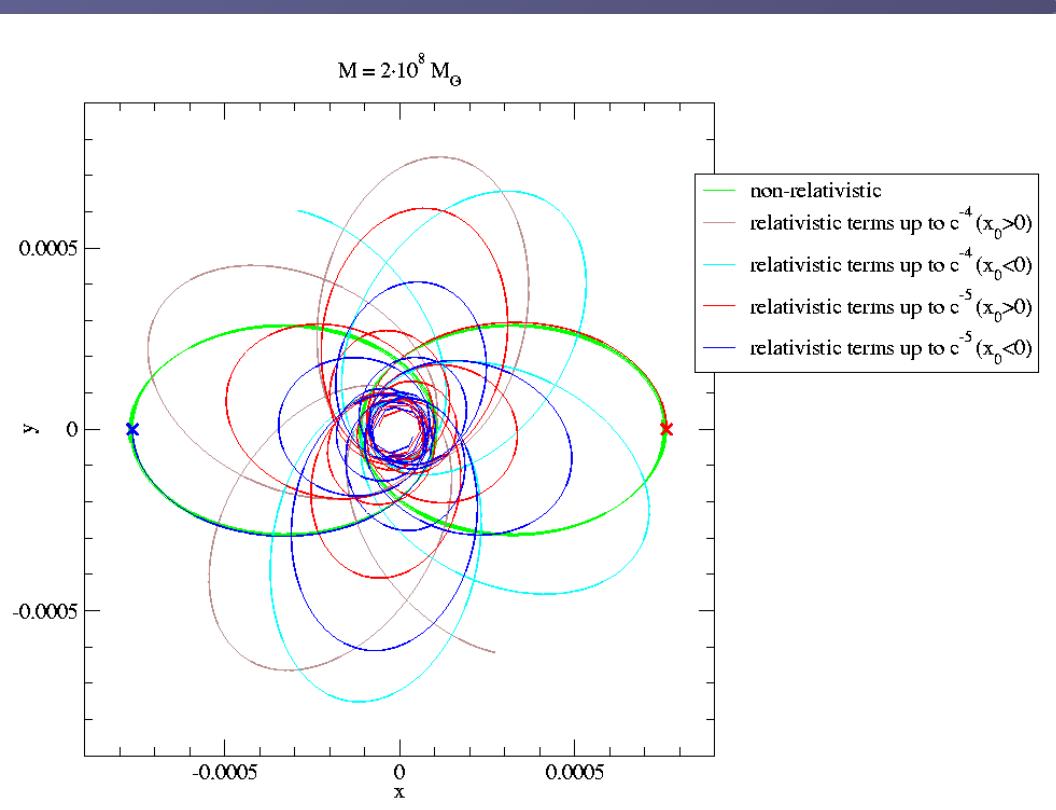
Schäfer, Gauge Theor. Grav. 36, 2223 (2004)

Memmesheimer, Gopakumar, Schäfer, Phys. Rev.D 70, 104011 (2004)

Blanchet, Luc; Living Reviews 2002, lrr-2002-3

Gravitational Waves

What happens afterwards? Post-Newton Order „2.5“...



Kupi, Amaro-Seoane & Spurzem 2006

EUROPEAN GRAVITATIONAL OBSERVATORY



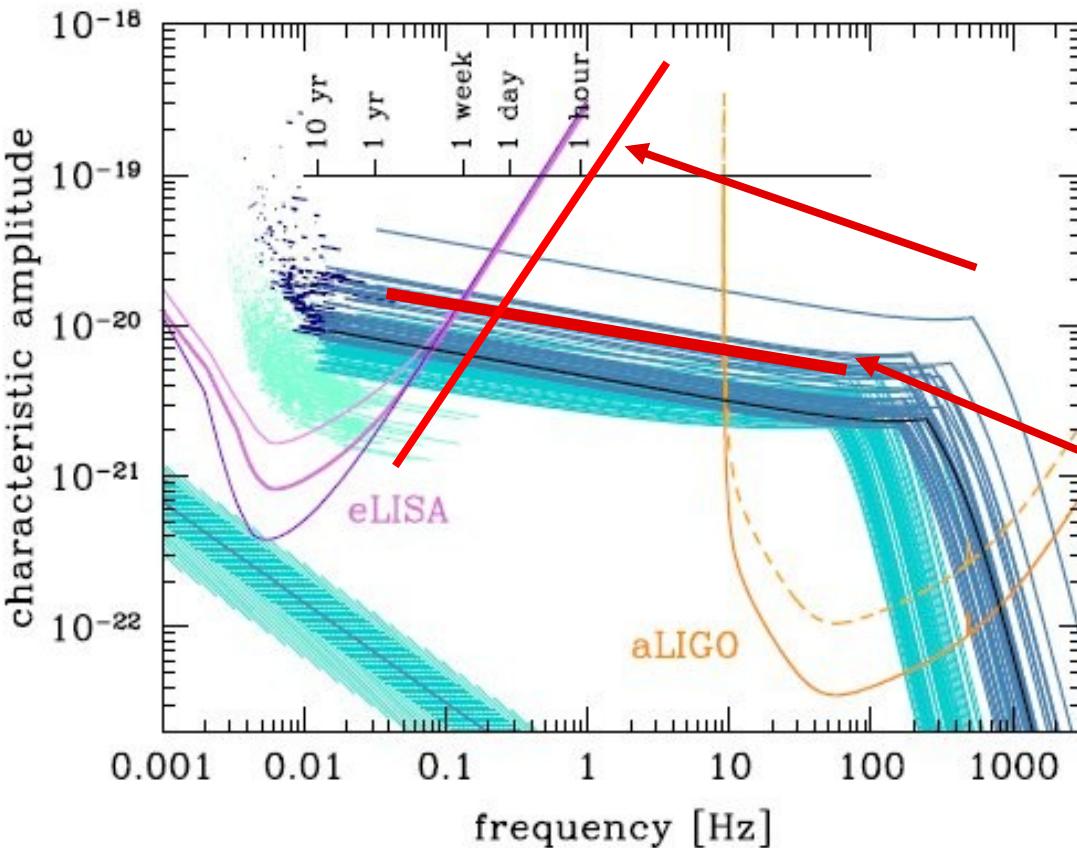
VIRGO Detector in Cascina near Pisa, Italy
(LIGO USA, KAGRA Japan)

Consortium of

LISA =
Laser Space
Interferometer Antenna



Gong, Lau,
... Spurzem ...
2015, 2011
(JphCS, CQGra)



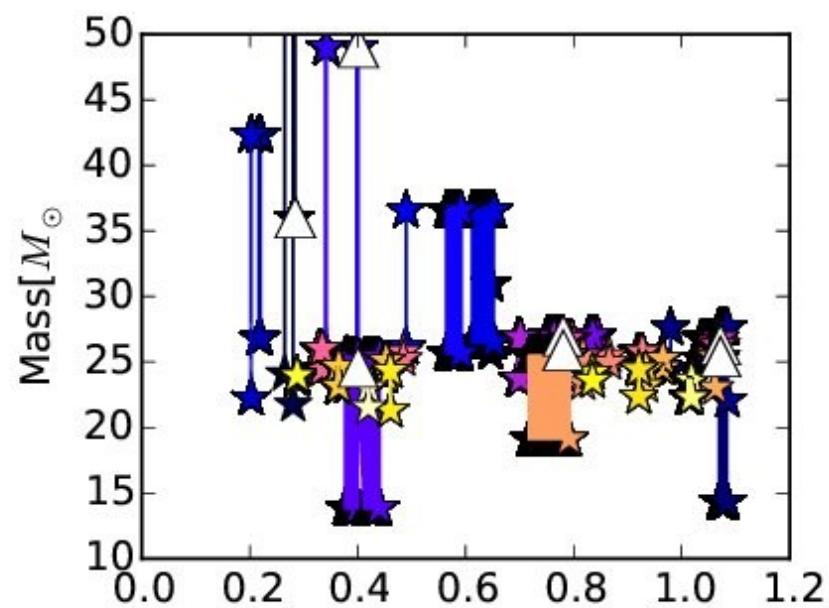
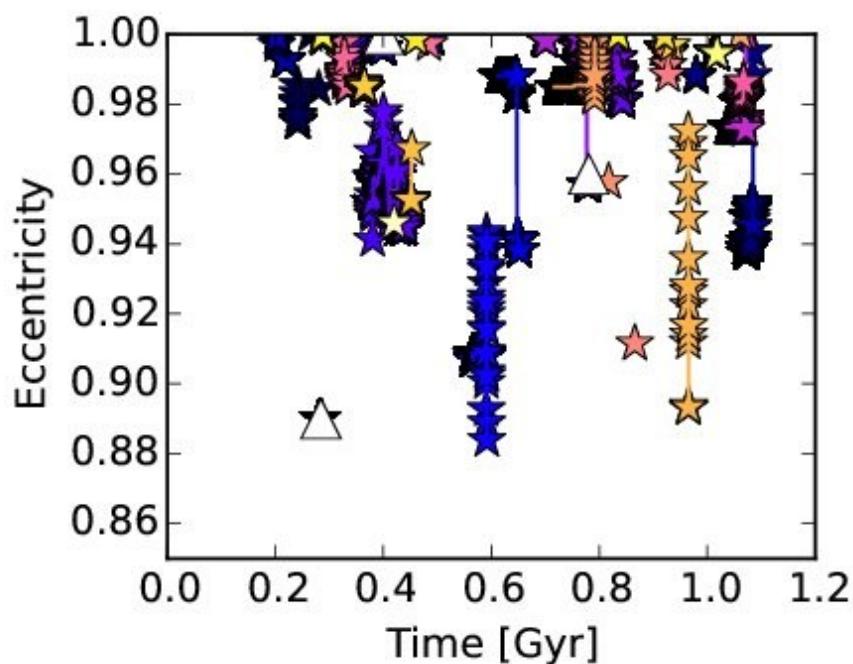
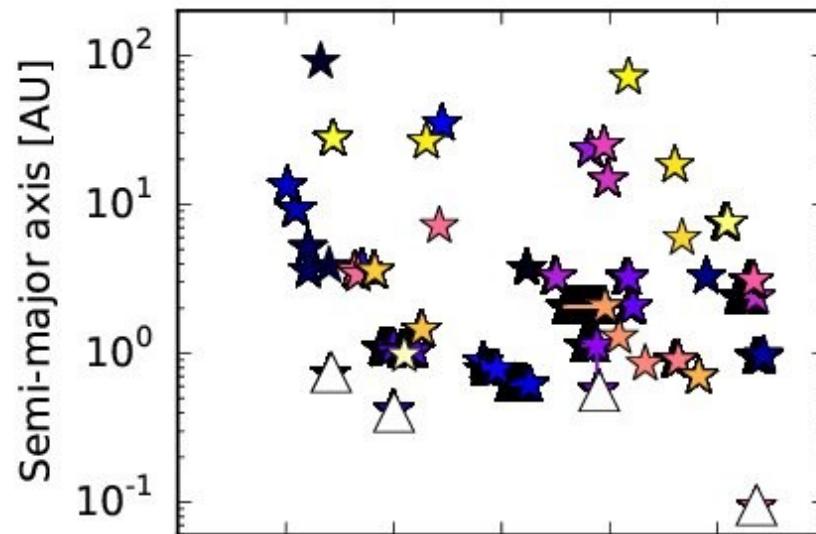
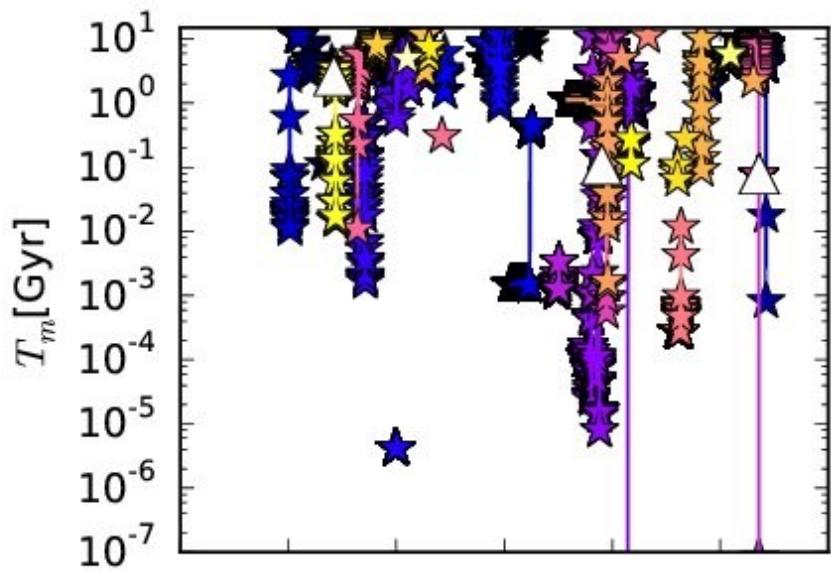
Taiji or
Tianqin
Hand
Drawn
Estimate

“Our”
DRAGON
Black Hole
Binary

Background Plot: Sesana 2016

FIG. 1: The multi-band GW astronomy concept. The violet lines are the total sensitivity curves (assuming two Michelson) of three eLISA configurations; from top to bottom N2A1, N2A2, N2A5 (from [III](#)). The orange lines are the current (dashed) and design (solid) aLIGO sensitivity curves. The lines in different blue flavours represent characteristic amplitude tracks of BHB sources for a realization of the *flat* population model (see main text) seen with $S/N > 1$ in the N2A2 configuration (highlighted as the thick eLISA middle curve), integrated assuming a five year mission lifetime. The light turquoise lines clustering around 0.01Hz are sources seen in eLISA with $S/N < 5$ (for clarity, we down-sampled them by a factor of 20 and we removed sources extending to the aLIGO band); the light and dark blue curves crossing to the aLIGO band are sources with $S/N > 5$ and $S/N > 8$ respectively in eLISA; the dark blue marks in the upper left corner are other sources with $S/N > 8$ in eLISA but not crossing to the aLIGO band within the mission lifetime. For comparison, the characteristic amplitude track completed by GW150914 is shown as a black solid line, and the chart at the top of the figure indicates the frequency progression of this particular source in the last 10 years before coalescence. The shaded area at the bottom left marks the expected confusion noise level produced by the same population model (median, 68% and 95% intervals are shown). The waveforms shown are second order post-Newtonian inspirals phenomenologically adjusted with a Lorentzian function to describe the ringdown.

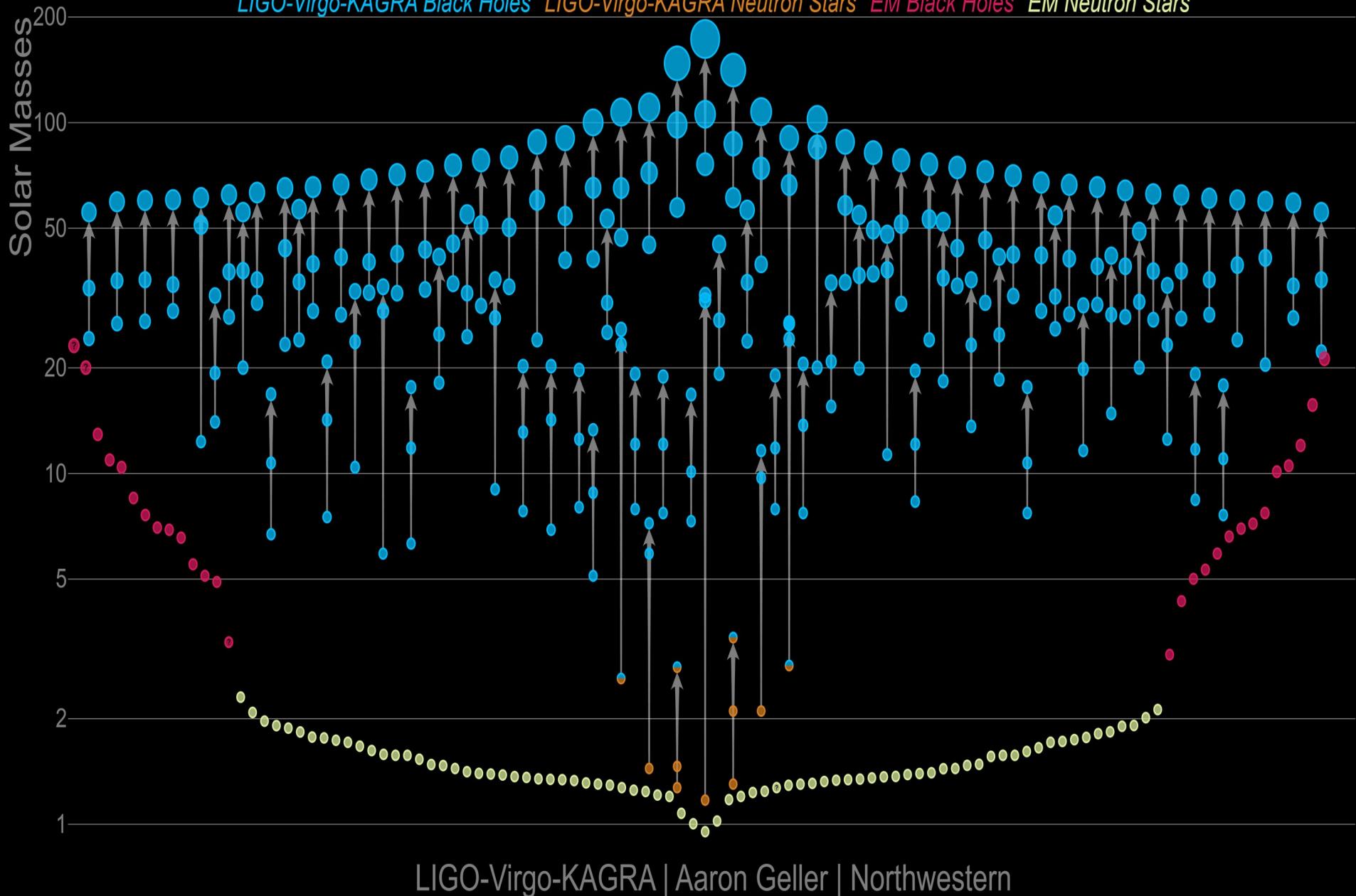
* * BH binaries in GCs ▲ ▲ BH binaries escapers



DRAGON II Simulations

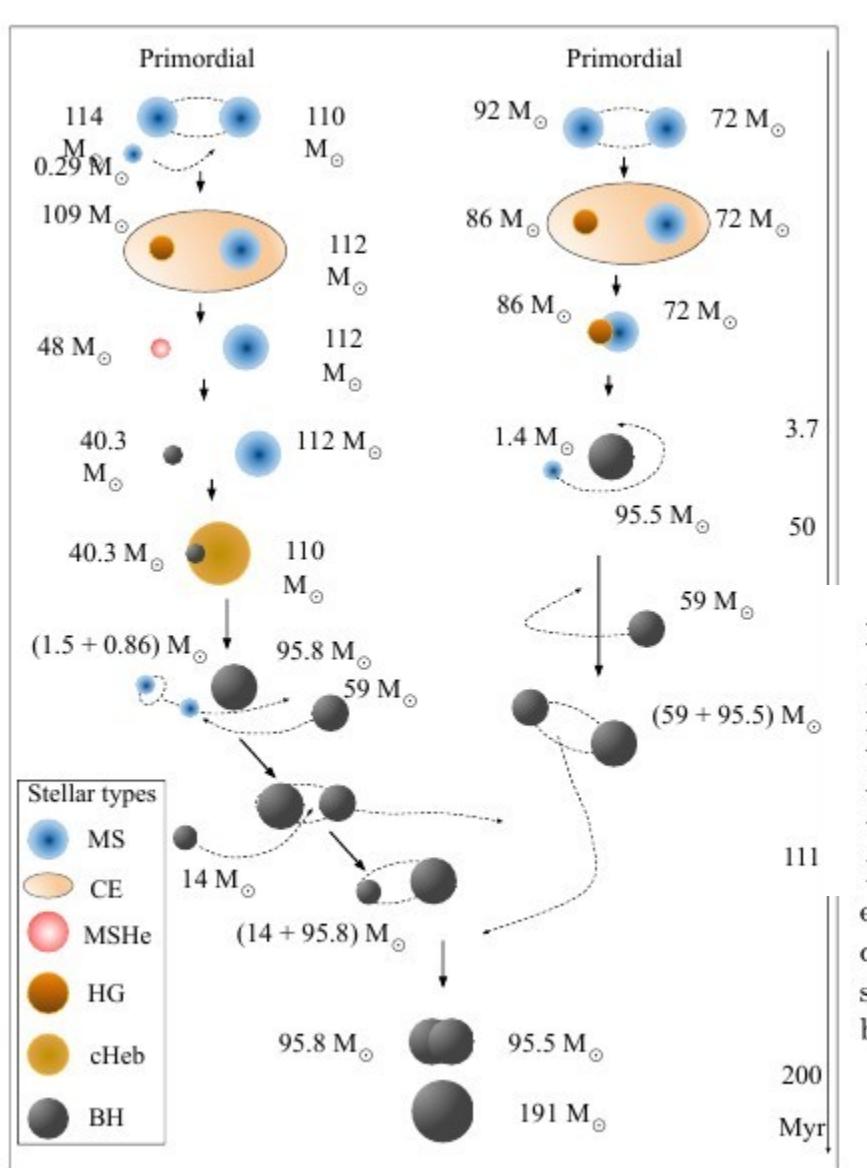
Masses in the Stellar Graveyard

LIGO-Virgo-KAGRA Black Holes LIGO-Virgo-KAGRA Neutron Stars EM Black Holes EM Neutron Stars



DRAGON-II Simulations – Paper II

using NBODY6++GPU



Arca Sedda et al. 2023abc:
MNRAS:

19 models, up to 1 million stars, up to 33% initial hard binaries

Including GR kicks for mergers!

Figure 2. Formation of an IMBH in simulation with $N = 120k$, $R_{\text{HM}} = 1.75$ pc, and $f_b = 0.2$, realization ID 0. Two massive primordial binaries undergo common envelope that eventually lead to the formation of two nearly equal mass BHs ($m_{\text{BH}} \sim 95 M_{\odot}$) that eventually find each other via a complex series of binary-binary interactions. The binary eventually merge and builds-up an IMBH with mass $m_{\text{IMBH}} \simeq 191 M_{\odot}$. The color-coded legend is ent colors correspond to different evolutionary stages: main sequence (MS), common envelope (CE), naked main sequence He star (MSHe), Hertzsprung gap (HG), core He burning (cHeb), and black hole (BH).

DRAGON-II Simulations using NBODY6++GPU

$$\left\langle \frac{da}{dt} \right\rangle = -\frac{64}{5} \frac{G^3 m_1 m_2 (m_1 + m_2)}{c^5 a^3 (1 - e^2)^{7/2}} \left(1 + \frac{73}{24} e^2 + \frac{37}{96} e^4 \right),$$

$$\left\langle \frac{de}{dt} \right\rangle = -\frac{304}{15} e \frac{G^3 m_1 m_2 (m_1 + m_2)}{c^5 a^4 (1 - e^2)^{5/2}} \left(1 + \frac{121}{304} e^2 \right).$$

Orbit Averaged Post-Newtonian
evolution of
semi-major axis a, eccentricity e
(Peters & Mathews 1963,
Peters 1964)

$$\begin{aligned} \vec{v}_{\text{GW}} &= v_m \hat{e}_{\perp,1} + v_{\perp} (\cos \xi \hat{e}_{\perp,1} + \sin \xi \hat{e}_{\perp,2}) + v_{\parallel} \hat{e}_{\parallel}, \\ v_m &= A \eta^2 \sqrt{1 - 4\eta} (1 + B\eta), \\ v_{\perp} &= \frac{H\eta^2}{1 + q_{\text{BBH}}} (S_{2,\parallel} - q_{\text{BBH}} S_{1,\parallel}), \\ v_{\parallel} &= \frac{16\eta^2}{1 + q_{\text{BBH}}} [V_{11} + V_A \Xi_{\parallel} + V_B \Xi_{\parallel}^2 + V_C \Xi_{\parallel}^3] \times \\ &\quad \times |\vec{S}_{2,\perp} - q_{\text{BBH}} \vec{S}_{1,\perp}| \cos(\phi_{\Delta} - \phi_1). \end{aligned}$$

Black Hole Binary Mergers
with relativistic kicks
including spins
(DRAGON II,
Arca Sedda et al. 2023ab, 2024)

How DRAGON-II Black Hole Binaries Grav. Wave Emission Could be observed?

(DRAGONII Papers,
ArcaSedda et al. 2023ab, 2024)

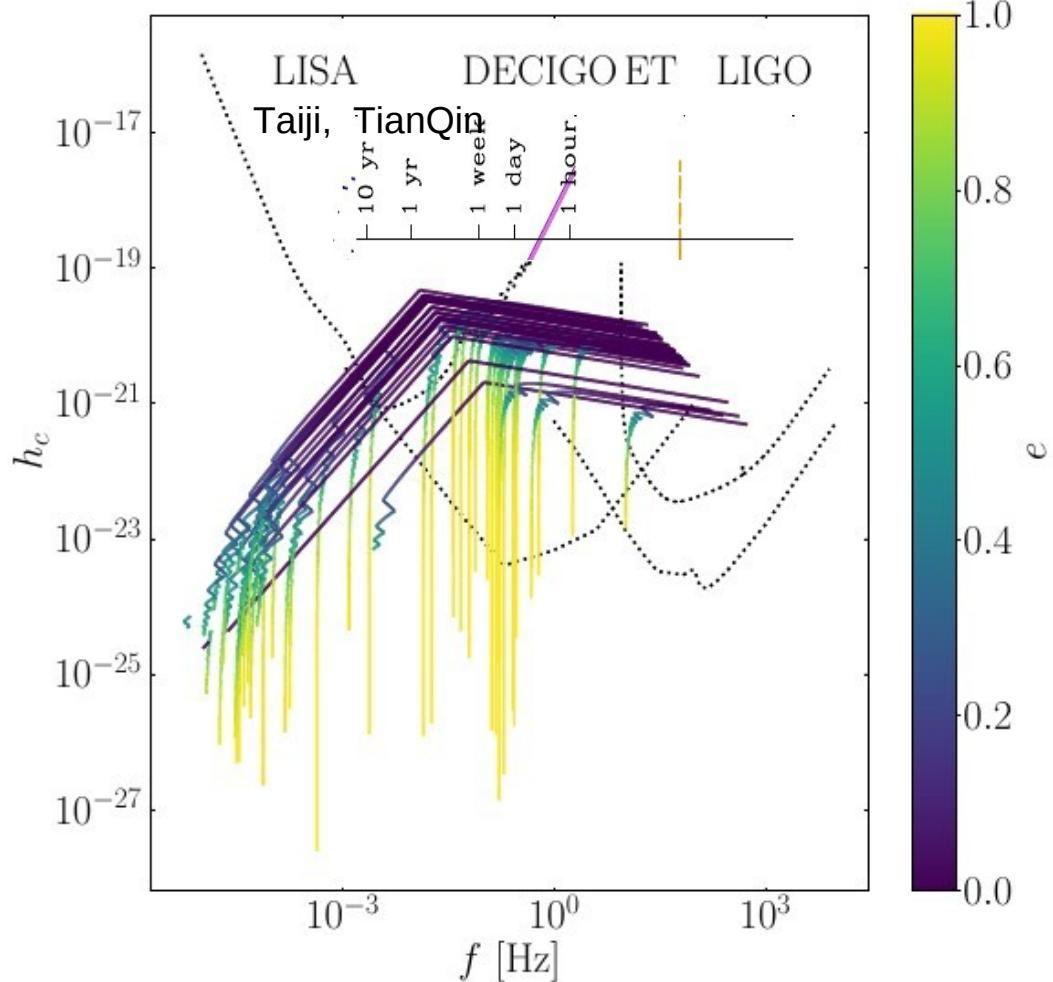


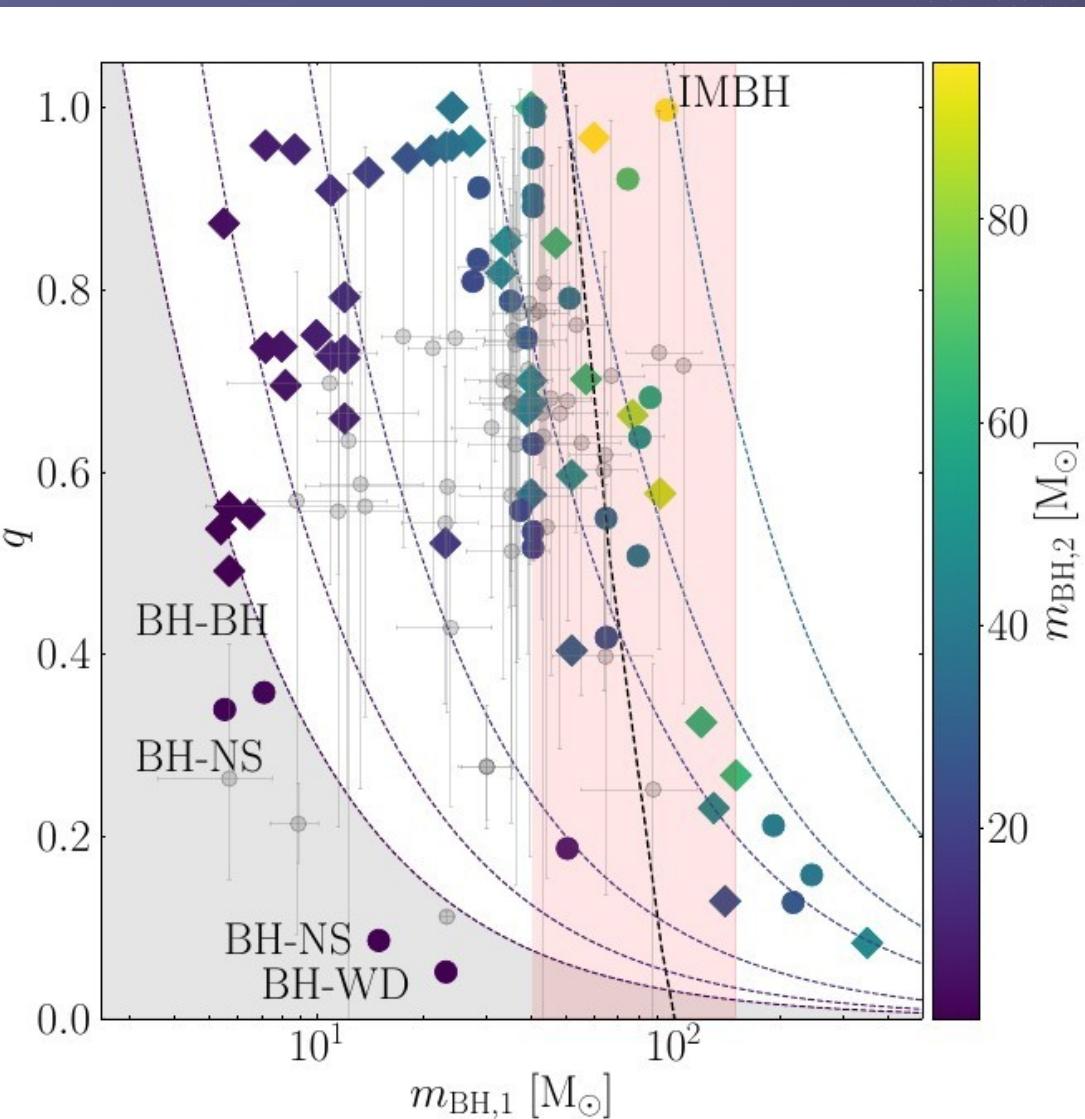
Fig.: ArcaSedda et al. 2023c;

Time Axis: Sesana et al. 2016

Figure 7. Evolution of the binary GW strain as a function of the frequency for all mergers in DRAGON-II simulations, assuming that the sources are located at a redshift $z = 0.05$. The colormap identifies the binary eccentricity along its orbit. Dotted lines represent sensitivity curves for different detectors, from left to right: LISA, DECIGO, ET, and LIGO.

DRAGON-II Simulations – Paper III

using NBODY6++GPU



*Arca Sedda et al. 2023ab,
2024: MNRAS*
*19 models, up to 1 million stars,
up to 33% initial hard binaries*

Compact Object Mergers
Compared with LIGO-Virgo
GWTC-3 catalogue (grey
symbols)

Mass ratio q vs. primary mass $m_{\text{BH},1}$; colour code:
secondary mass $m_{\text{BH},2}$

grey shade: neutron star
involved

red Shade: mass gap

What about spins?

Post-Newtonian Dynamics

Spin-Orbit Interaction S / Spin-Spin SS

$$\begin{aligned}\frac{d\mathbf{v}_1}{dt} = & \mathbf{A}_{\text{N}} + \frac{1}{c^2} \mathbf{A}_{1\text{PN}} + \frac{1}{c^3} \mathbf{A}_{1.5\text{PN}} + \frac{1}{c^4} [\mathbf{A}_{2\text{PN}} + \mathbf{A}_{2\text{PN}}] \\ & + \frac{1}{c^5} [\mathbf{A}_{2.5\text{PN}} + \mathbf{A}_{2.5\text{PN}}] + \mathcal{O}\left(\frac{1}{c^6}\right).\end{aligned}\quad (5.1)$$

Faye, Blanchet, Buonanno 2006

$$\begin{aligned}\mathbf{A}_{1.5\text{PN}} = & \frac{Gm_2}{r_{12}^3} \left\{ \left[6 \frac{(S_1, n_{12}, \mathbf{v}_{12})}{m_1} + 6 \frac{(S_2, n_{12}, \mathbf{v}_{12})}{m_2} \right] \mathbf{n}_{12} \right. \\ & + 3(n_{12}\mathbf{v}_{12}) \frac{\mathbf{n}_{12} \times \mathbf{S}_1}{m_1} + 6(n_{12}\mathbf{v}_{12}) \frac{\mathbf{n}_{12} \times \mathbf{S}_2}{m_2} \\ & \left. - 3 \frac{\mathbf{v}_{12} \times \mathbf{S}_1}{m_1} - 4 \frac{\mathbf{v}_{12} \times \mathbf{S}_2}{m_2} \right\}.\end{aligned}\quad (5.3a)$$

Post- Newtonian Dynamics Gravitational Wave Templates

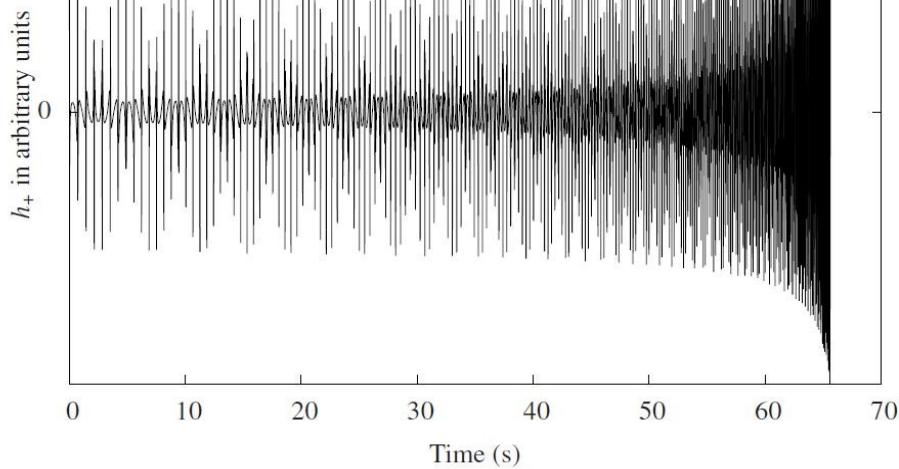


Figure 3.11: Waveform for two equal mass objects on an orbit with $e = 0.5$.

Brem,
Amaro-Seoane,
Spurzem,
MNRAS 2013

Handle spin-orbit and
spin-spin coupling
(P.Brem, R. Spurzem,
Univ. Heidelberg)

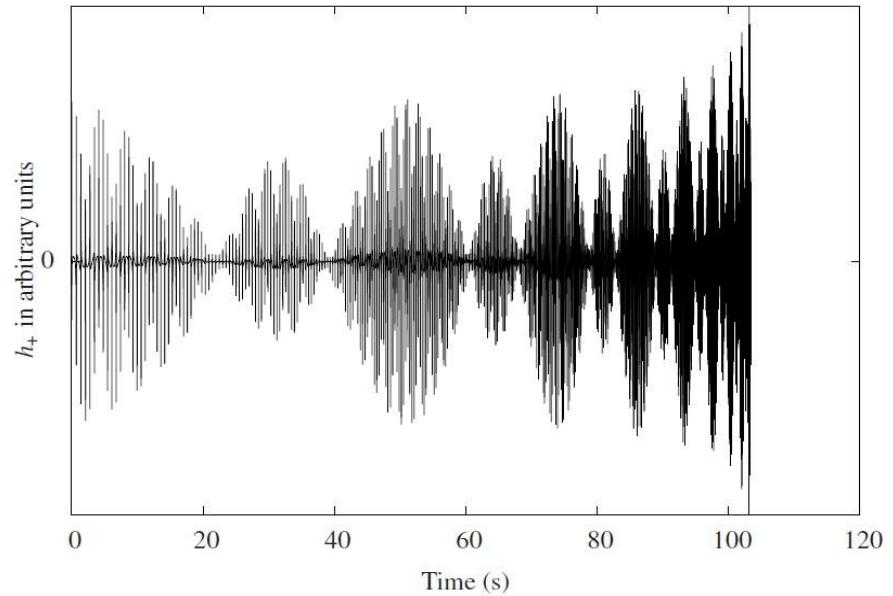
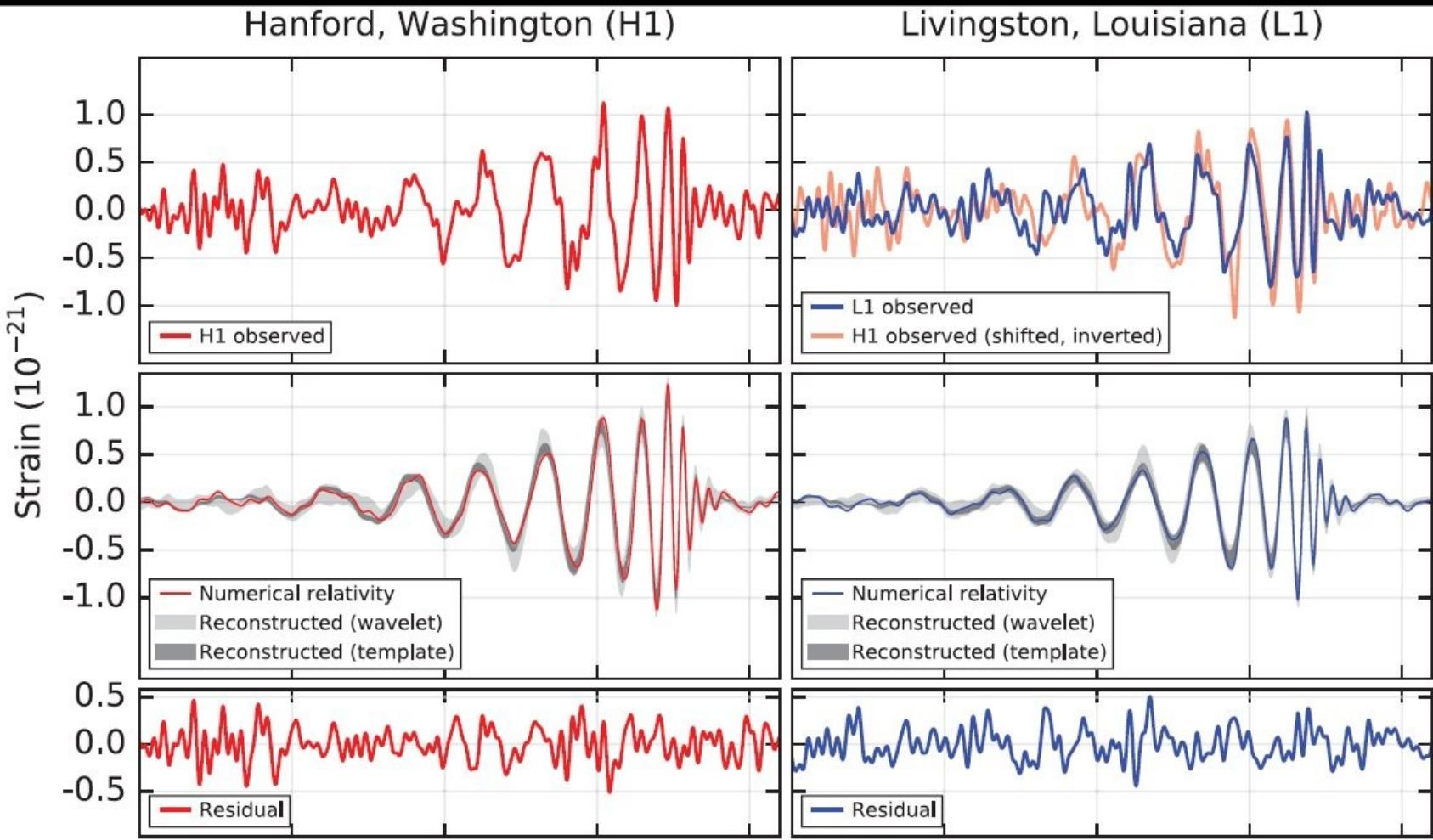


Figure 3.12: Waveform for two objects with a mass ratio of $q = 1/10$ on an orbit with $e = 0.5$ and spins $a_{1,x} = 1.0$, $a_{2,y} = 1.0$.

GW Detection Abbott et al. 2016



$$|\mathbf{a}_{\text{fin}}| = \frac{1}{(1+q)^2} [|\mathbf{a}_1|^2 + |\mathbf{a}_2|^2 q^4 + 2|\mathbf{a}_2||\mathbf{a}_1|q^2 \cos \alpha \\ + 2(|\mathbf{a}_1|\cos \beta + |\mathbf{a}_2|q^2 \cos \gamma)|\mathbf{l}|q + |\mathbf{l}|^2 q^2]^{1/2},$$

where $q = M_2/M_1$ is the mass ratio and the angles are defined as

$$\cos \alpha = \hat{\mathbf{a}}_1 \cdot \hat{\mathbf{a}}_2, \quad \cos \beta = \hat{\mathbf{a}}_1 \cdot \hat{\mathbf{l}}, \quad \cos \gamma = \hat{\mathbf{a}}_2 \cdot \hat{\mathbf{l}}.$$

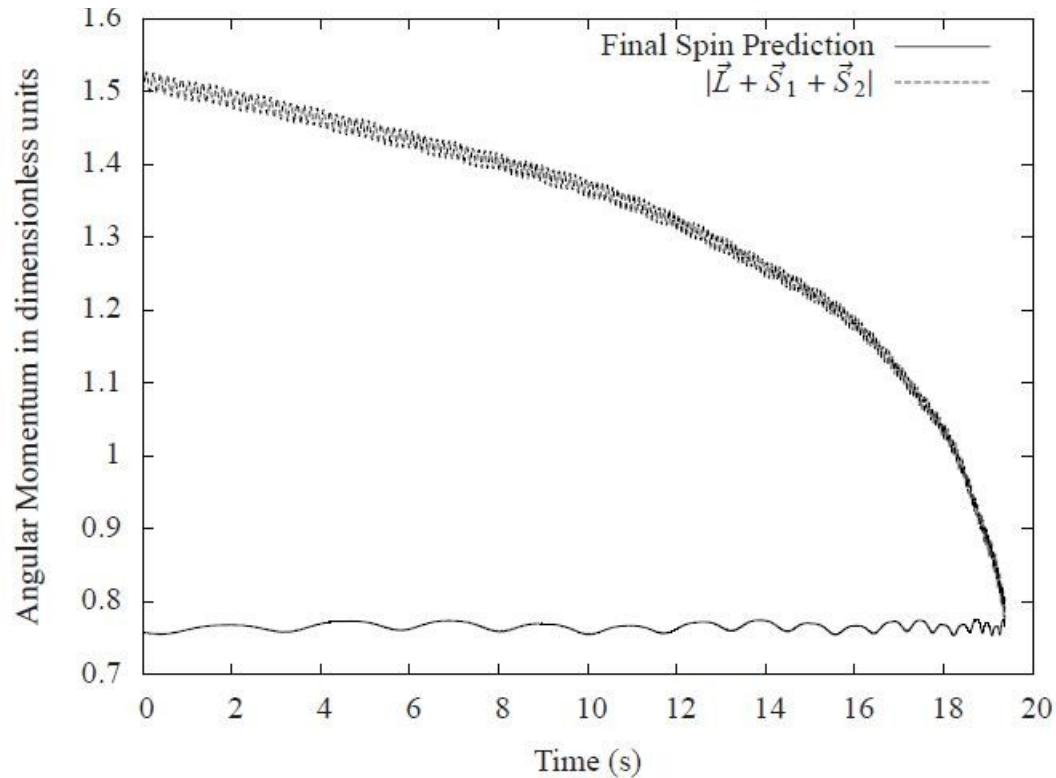


Figure 3.7: Comparison between the current final spin prediction and the actual total angular momentum of the binary system.

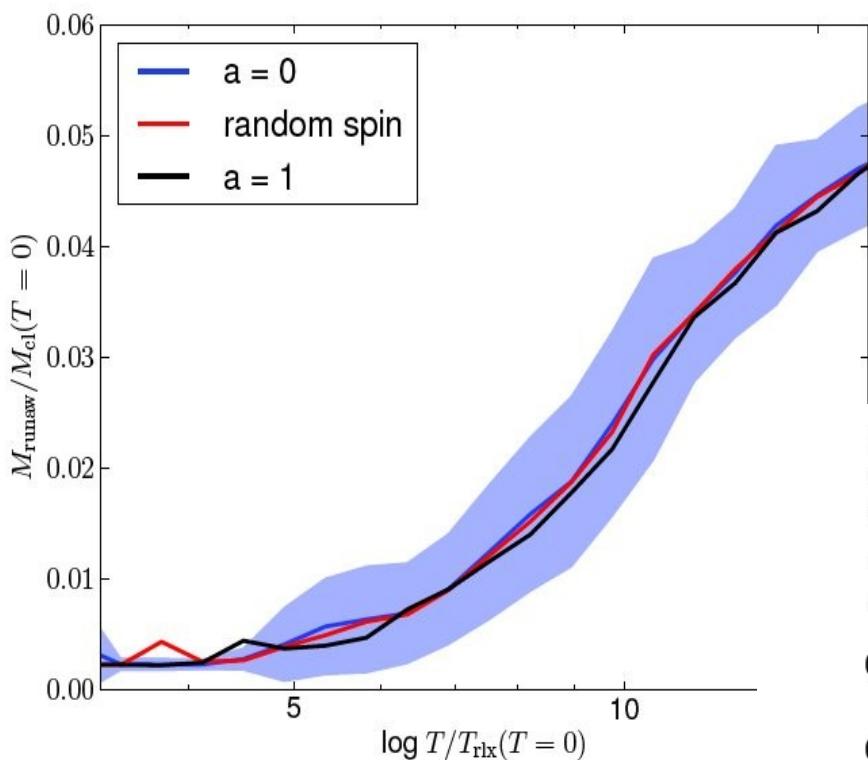


Figure 8. Mass of the runaway body, M_{runaw} , for each setup, averaged over 500 runs. $M_{\text{cl}}(T = 0)$ is the total mass of the cluster at the time $T = 0$ and $T_{\text{rlx}}(T = 0)$ the initial relaxation time of the cluster. The shaded area shows the standard deviation for the $a = 0$ case.

Brem, Amaro-Seoane,
Spurzem, MNRAS, 2013

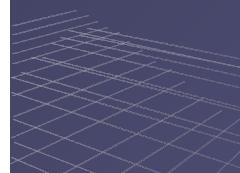
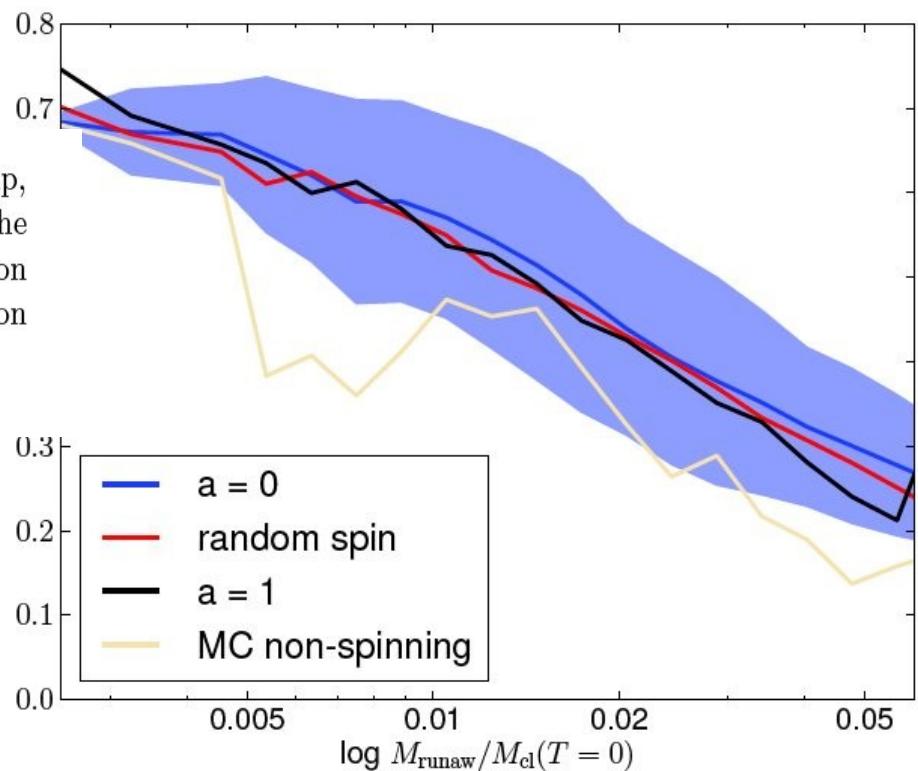


Figure 9. Spin of the runaway body in each simulation, averaged over 500 runs. The shaded area shows the standard deviation for the $a = 0$ case. All initial spin setups lead to a similar evolution, except for the very first data point which is slightly higher for the maximally spinning initial conditions.



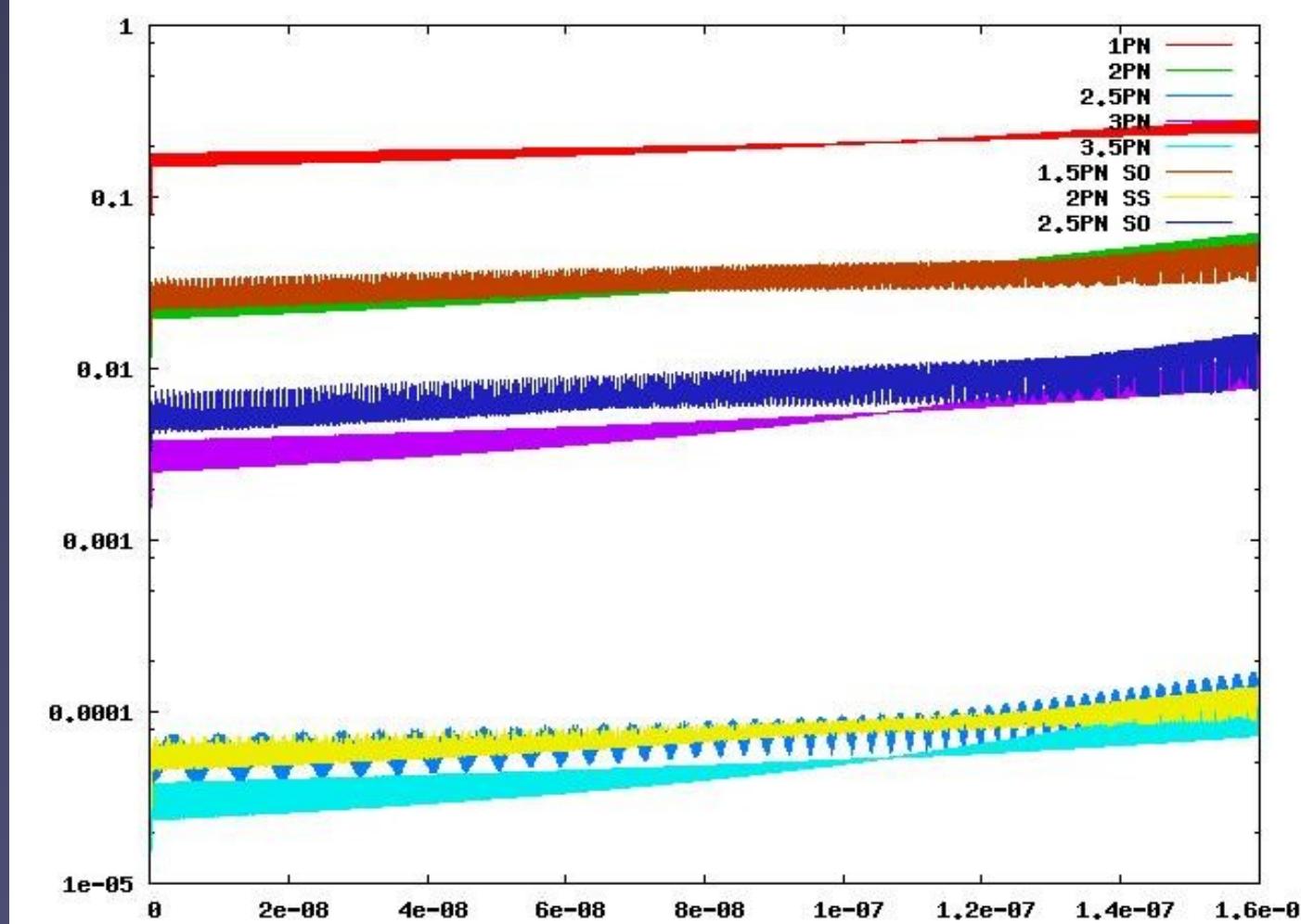
Post-Newtonian Dynamics

Brem, Amaro-Seoane,
Spurzem,
MNRAS 2013

Include
Spin-Orbit
Spin-Spin
PN3, PN3.5
Spin Dynamics

By Patrick Brem
(Diploma Thesis
Univ. Heidelberg)

1PN
2PN + 1.5PN SO
3PN + 2.5PN SO
2.5PN + 2PN SS
3.5PN

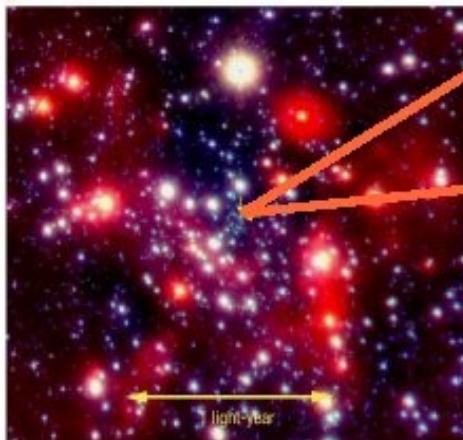


- 1) Introduction, Globular Star Clusters
- 2) Nuclear Star Clusters**
- 3) Code(s) and Hardware

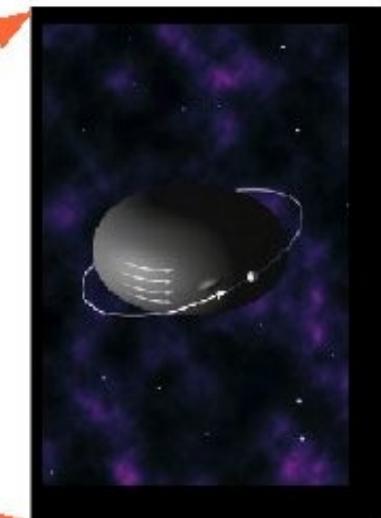
Setting the stage: the galactic nucleus



Size ~ 10 Kpc
Density ~ $0.05 M_{\text{sun}} \text{ pc}^{-3}$
Vel. Disp. ~ 40 Km/s
Relaxation time ~ 10^{15} yrs.



Size ~ 1-10 pc
Density ~ $10^{6-8} M_{\text{sun}} \text{ pc}^{-3}$
Vel. Disp. ~ 10^{2-3} Km/s
Relaxation time ~ 10^{8-9} yrs.

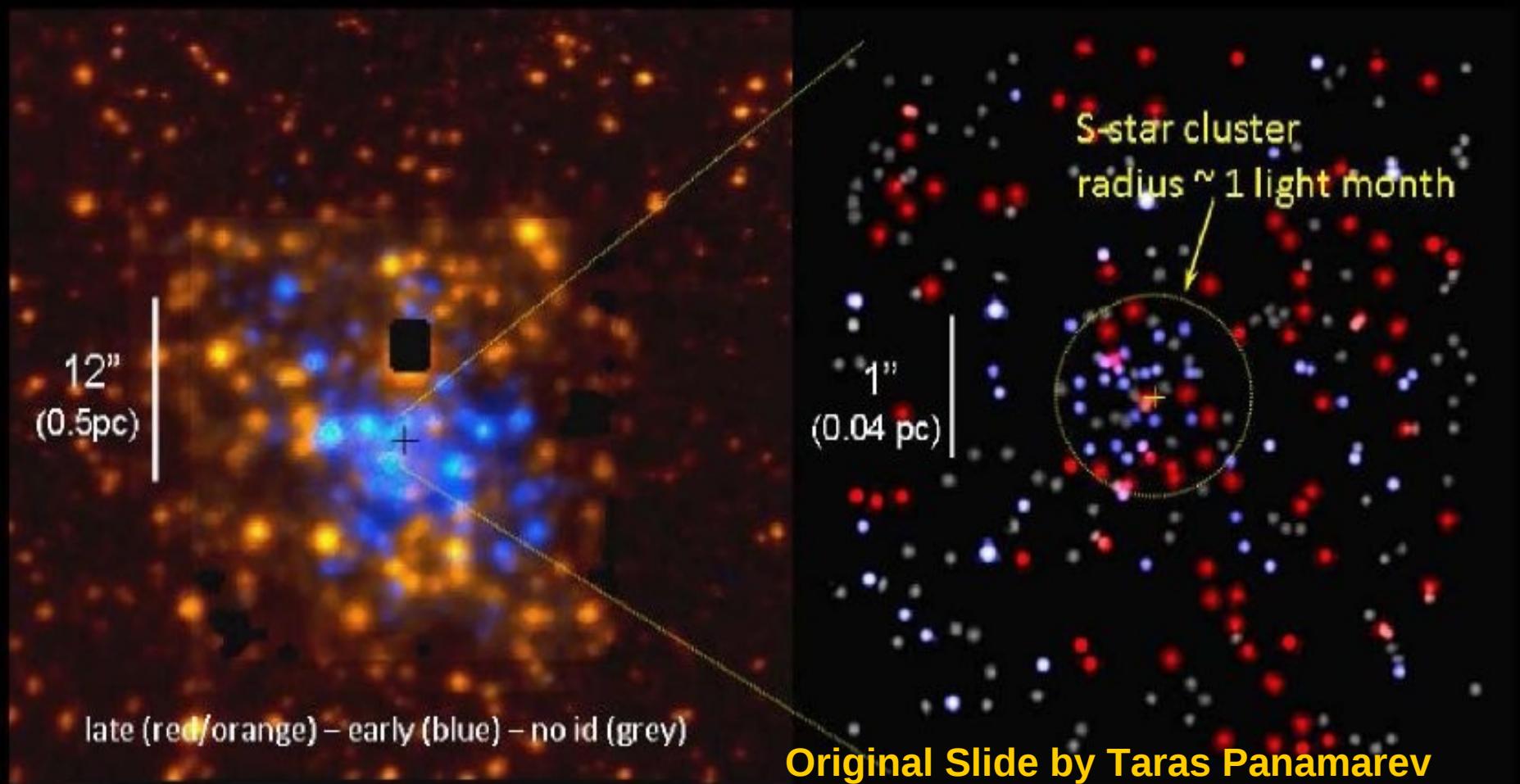


Size ~ $10^{-7}-10^{-4}$ pc
 $R_s = 2G M_{\text{BH}} / c^2$
 $R_t \sim (\alpha M_{\text{BH}} / m_*)^{1/3} R_*$
Loss cone aperture: θ

Slide:
Miguel
Preto

Distribution of stars

Model: Panamarev, Just, Spurzem, Berczik, Wang, Arca Sedda 2019, MNRAS
“DRAGON” of Galactic Center (one million bodies. SMBH. 1% hard binaries)



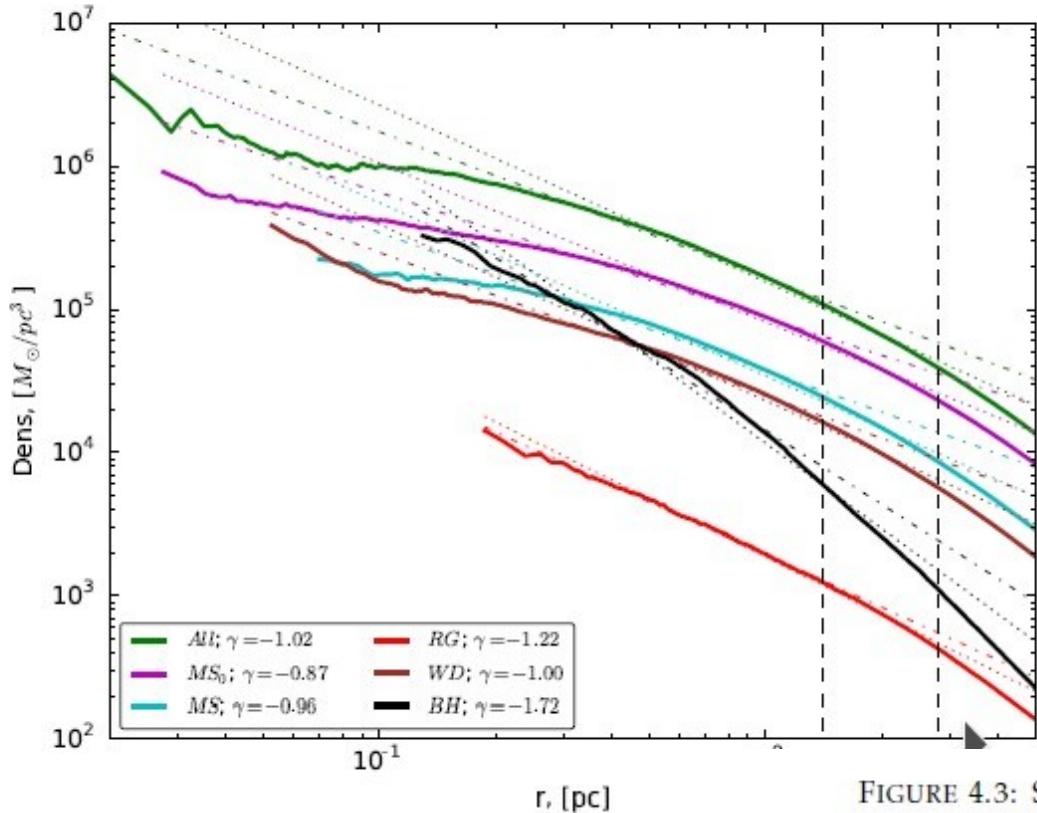
Black Holes in Galaxies



Stars moving
Around the
Central black
Hole in our Galaxy

Max-Planck Inst.
Extraterrestrische
Physik Garching
<http://www.mpe.mpg.de/ir/GC/>

DRAGON I Galactic Center Simulations



Initial Data:

1 million stars
10% fixed SMBH mass
Zero age pop, 0.8 – 100
Spherical

Density Profiles of stars and *-mass Black Holes in the Galactic Center after 5 Gyr

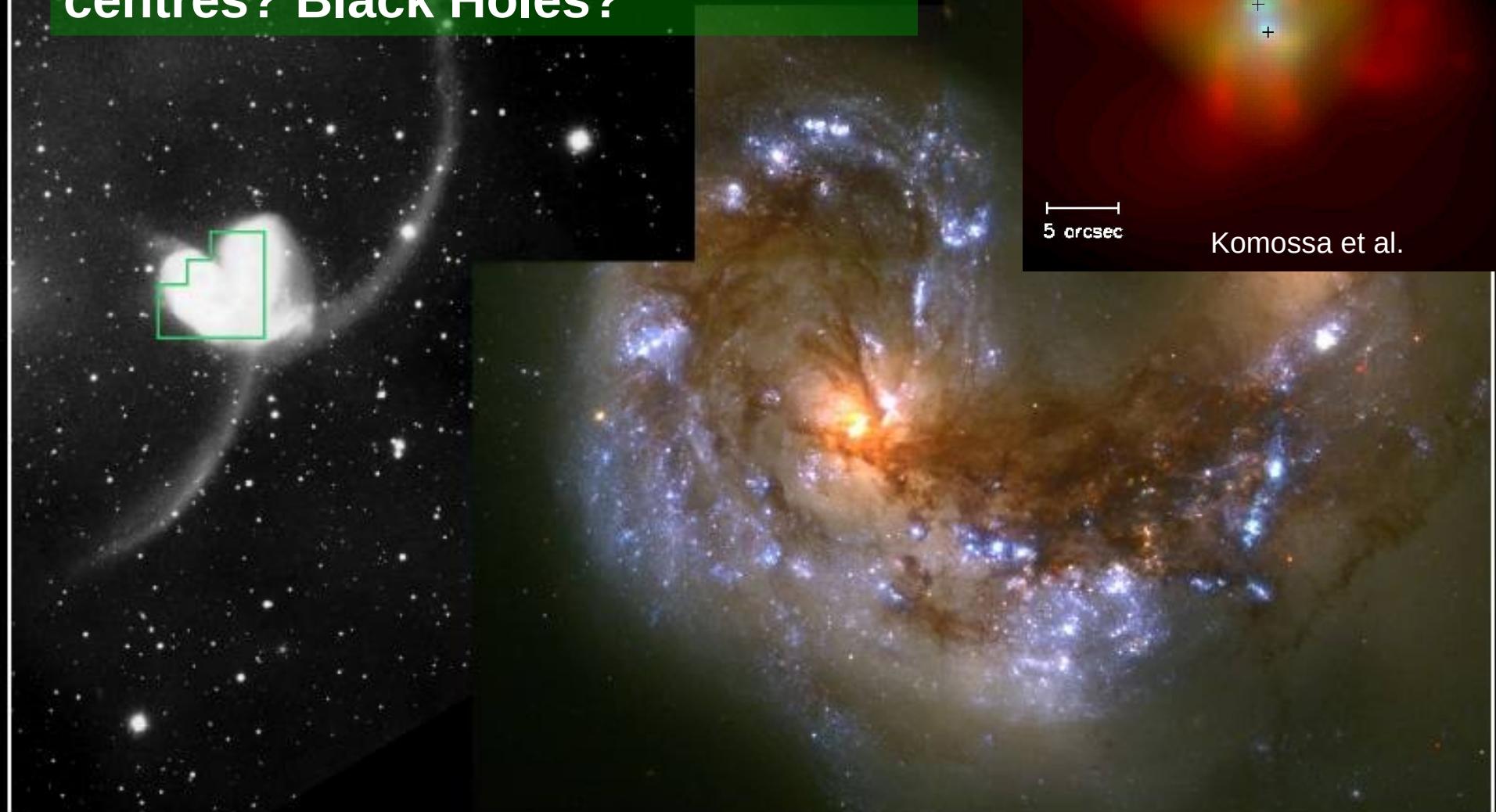
Extremely simple
“accretion radius”
For ALL objects.

FIGURE 4.3: Stellar density profiles at $t = 5$ Gyr for different stellar types. Thick solid lines correspond to: All - all stars, MS_{low} - low mass main sequence stars, MS - main sequence, RG - red giants, WD - white dwarfs, BH - black holes. Corresponding power-law slopes fitted inside the initial and final influence radii of the SMBH are shown as dash-dotted and dotted lines of the same colour. The dashed vertical lines denote the initial influence radius ($r = 1.4pc$) and the influence radius at $t = 5Gyr$ ($r \sim 2.8pc$) of the SMBH. The power-law indices fitted inside $r = 1.4pc$ are shown in the legend.

Panamarev, Just, Spurzem, et al. 2019

Panamarev, ..., Just, Spurzem, 2018, MNRAS

Galaxies merge, hierarchical Structure formation, their centres? Black Holes?

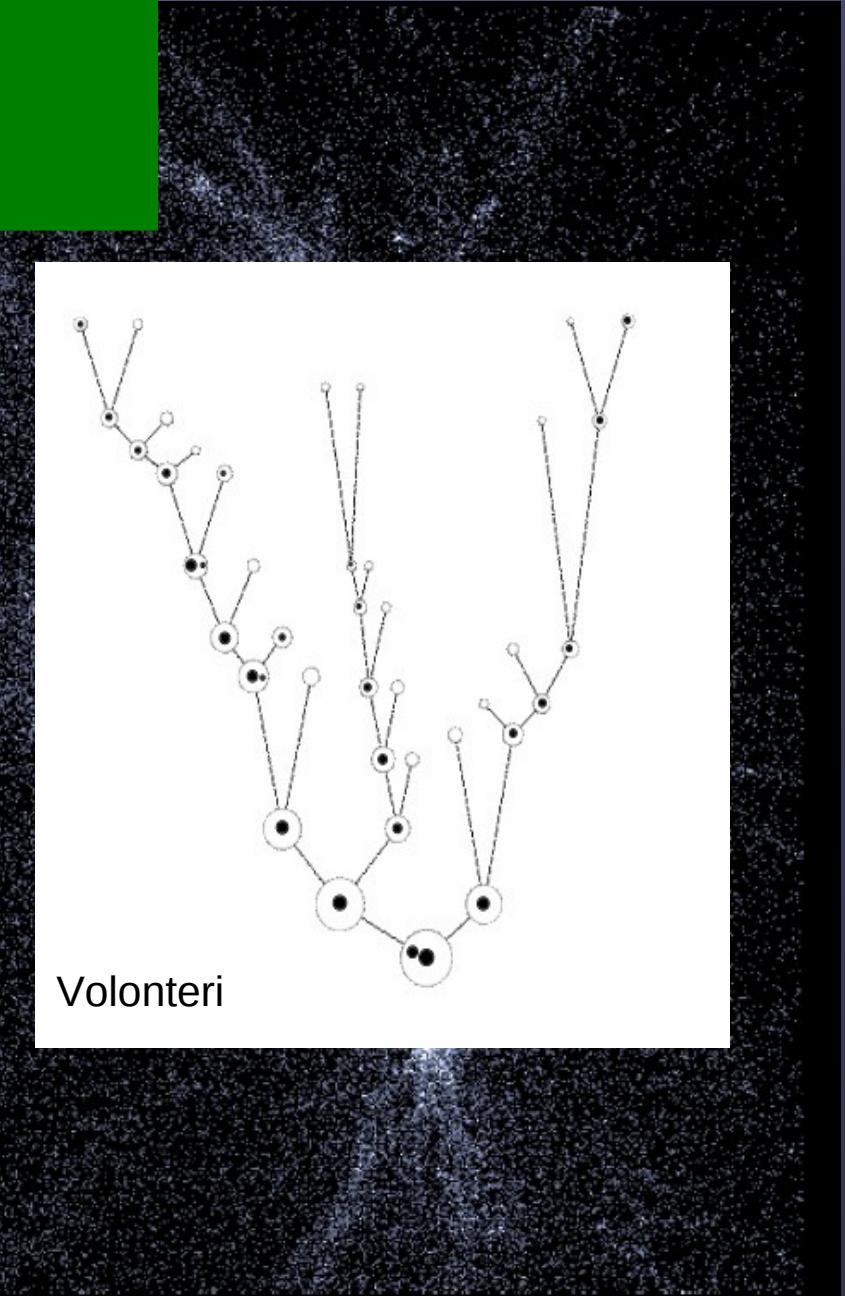
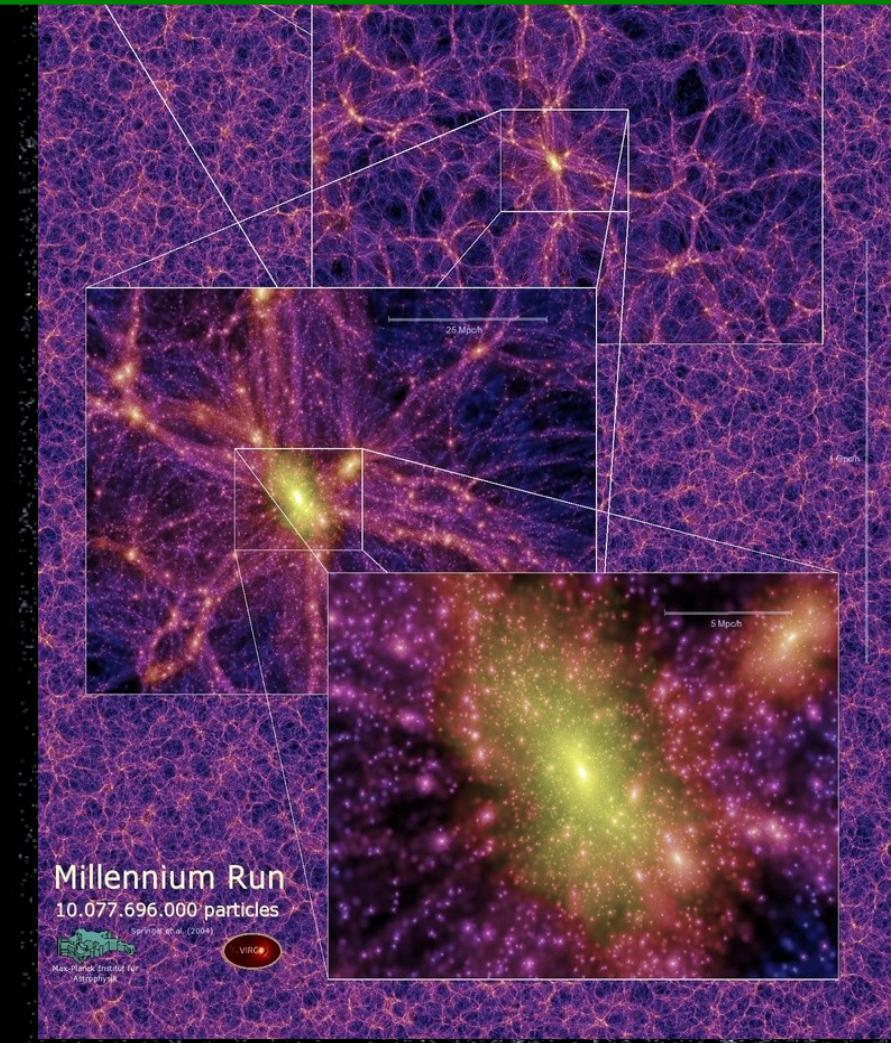


Colliding Galaxies NGC 4038 and NGC 4039

PRC97-34a • ST Scl OPO • October 21, 1997 • B, Whitmore (ST Scl) and NASA

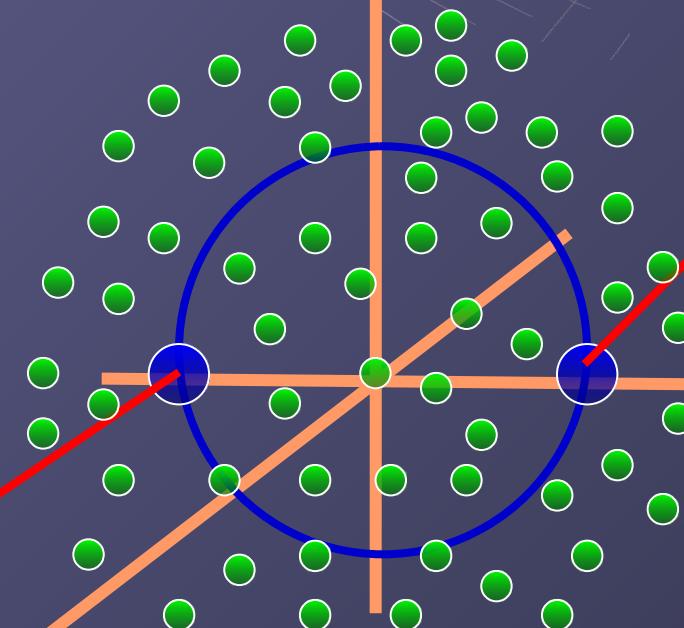
HST • WFPC2

Galaxies merge, hierarchical Structure formation, their centres? Black Holes?

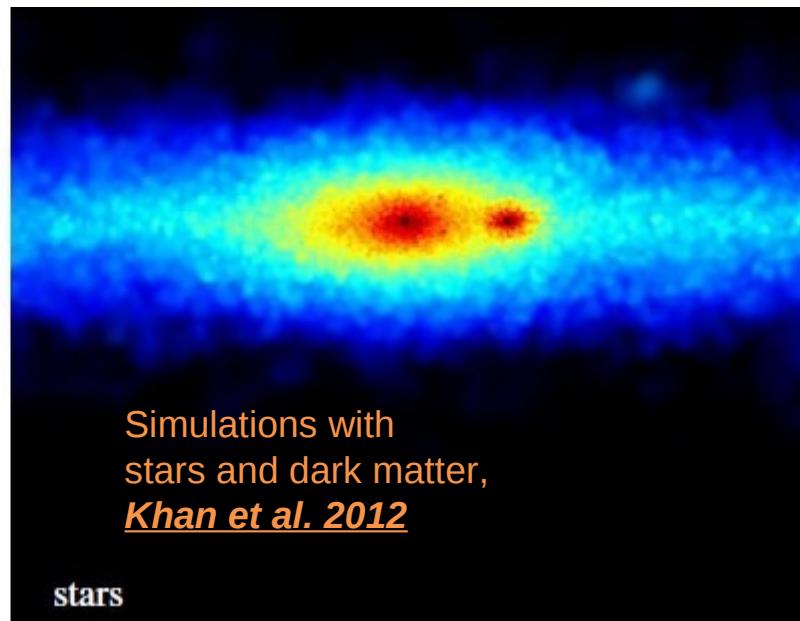
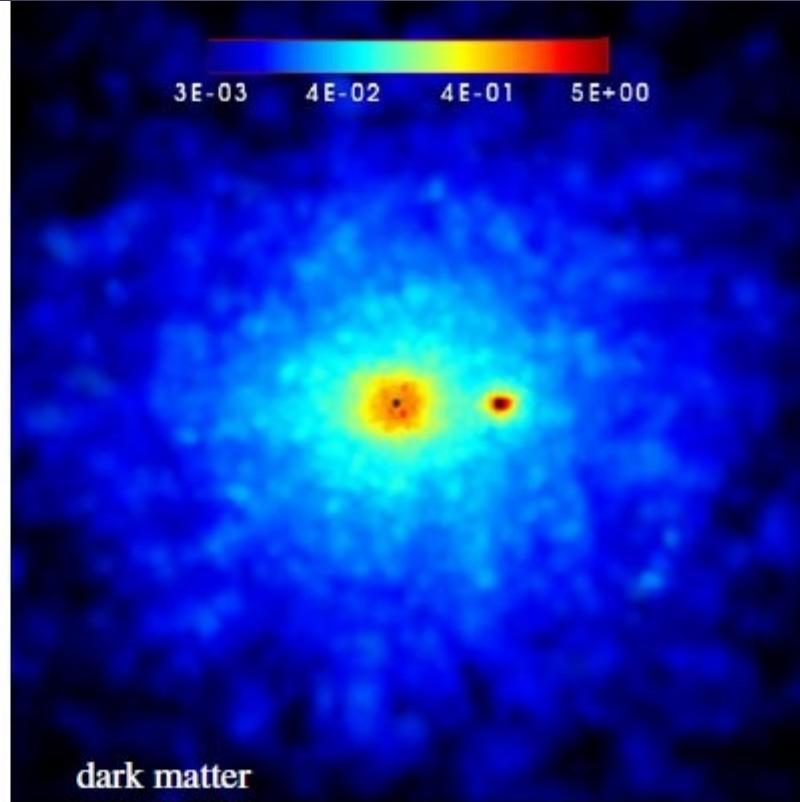
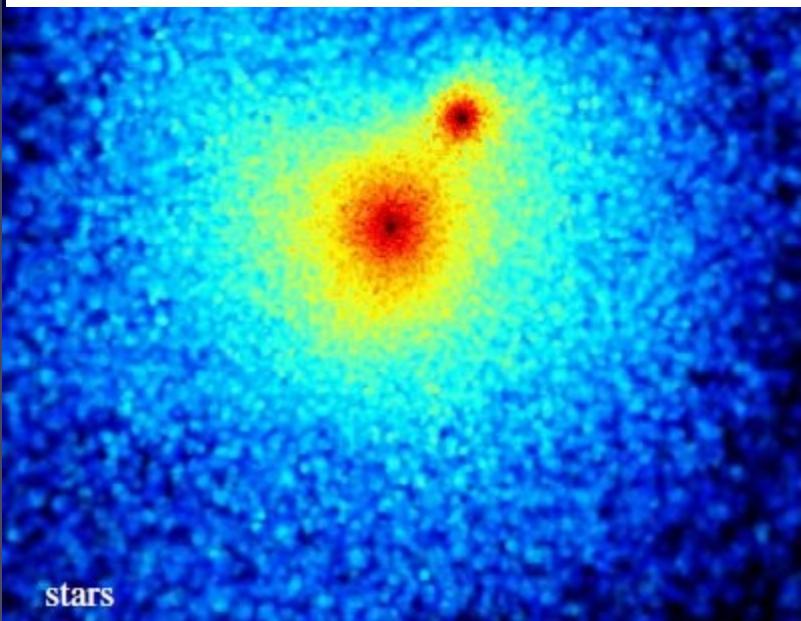
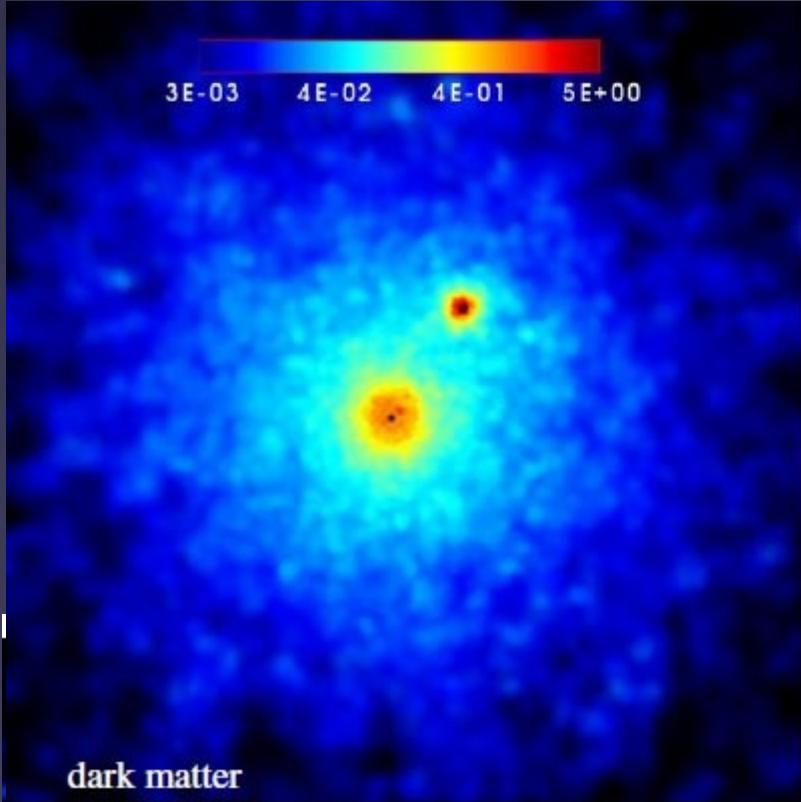


Simulations of Binary Black Holes

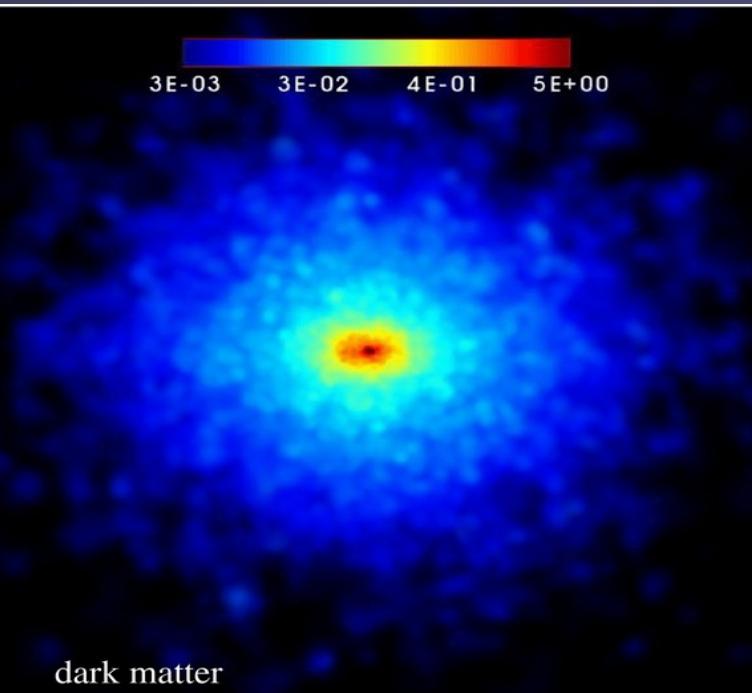
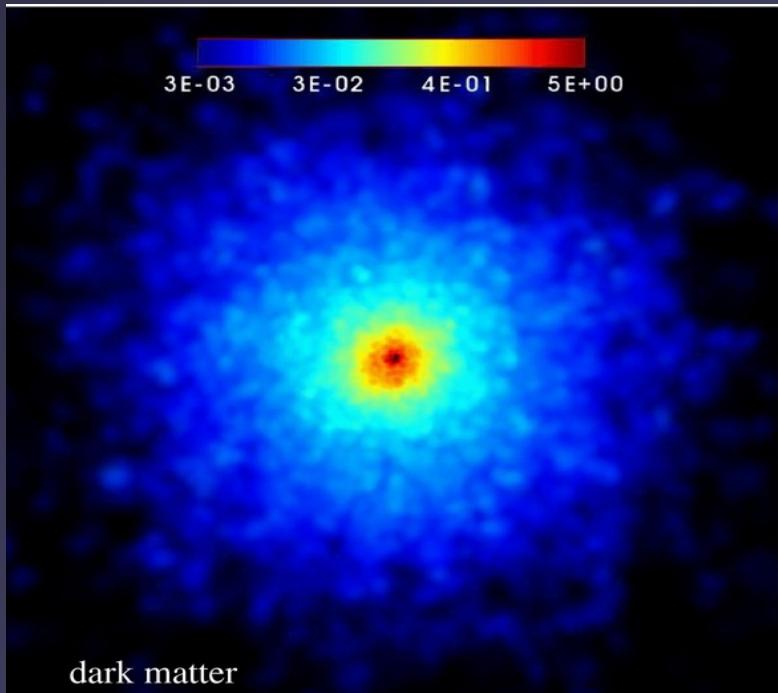
Black Holes in Dense
Stellar Clusters:
N-Body Problem
with
General Relativity
(Post-Newtonian)



Two (un-)equal-mass
black holes in central core
of simplified galaxy
(up to 6 million particles)

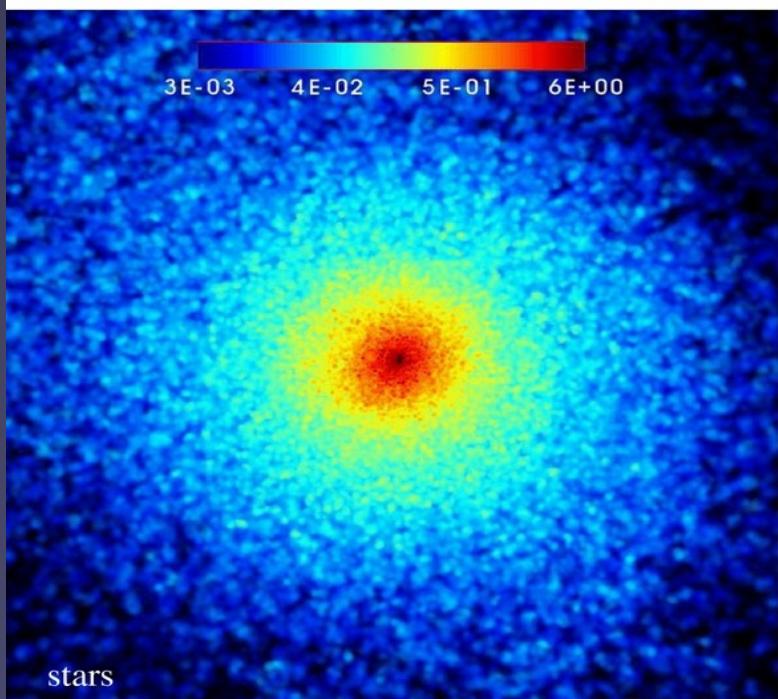


Simulations with
stars and dark matter,
Khan et al. 2012

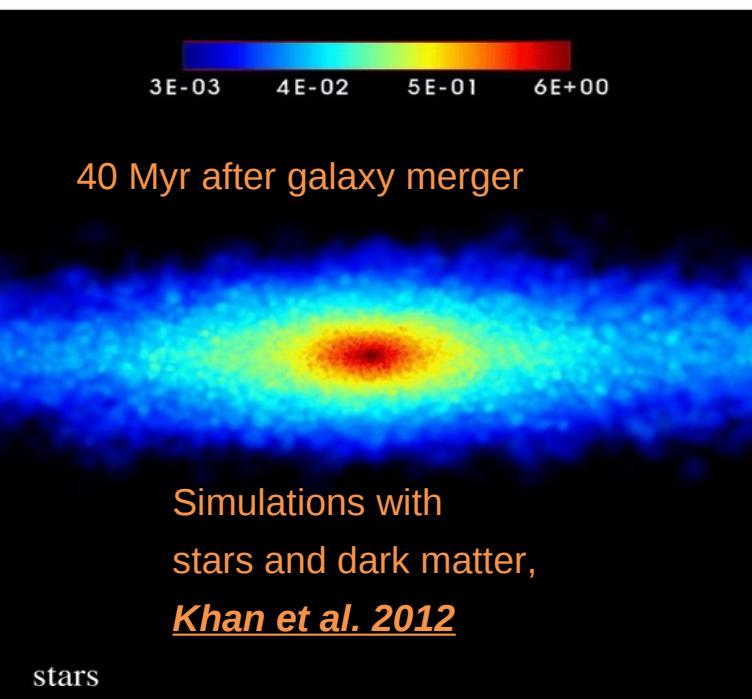


Box
Size
4 kpc

dark matter



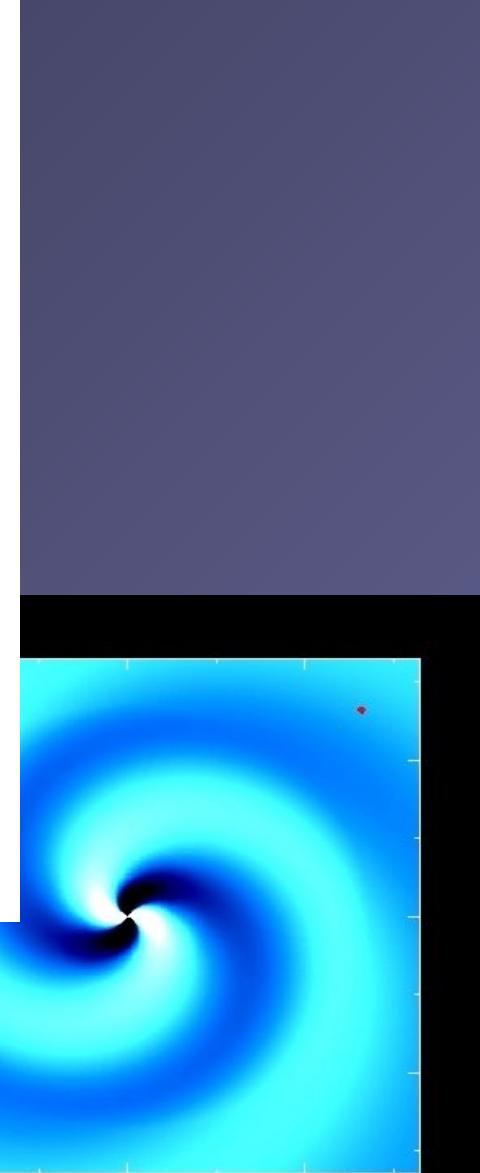
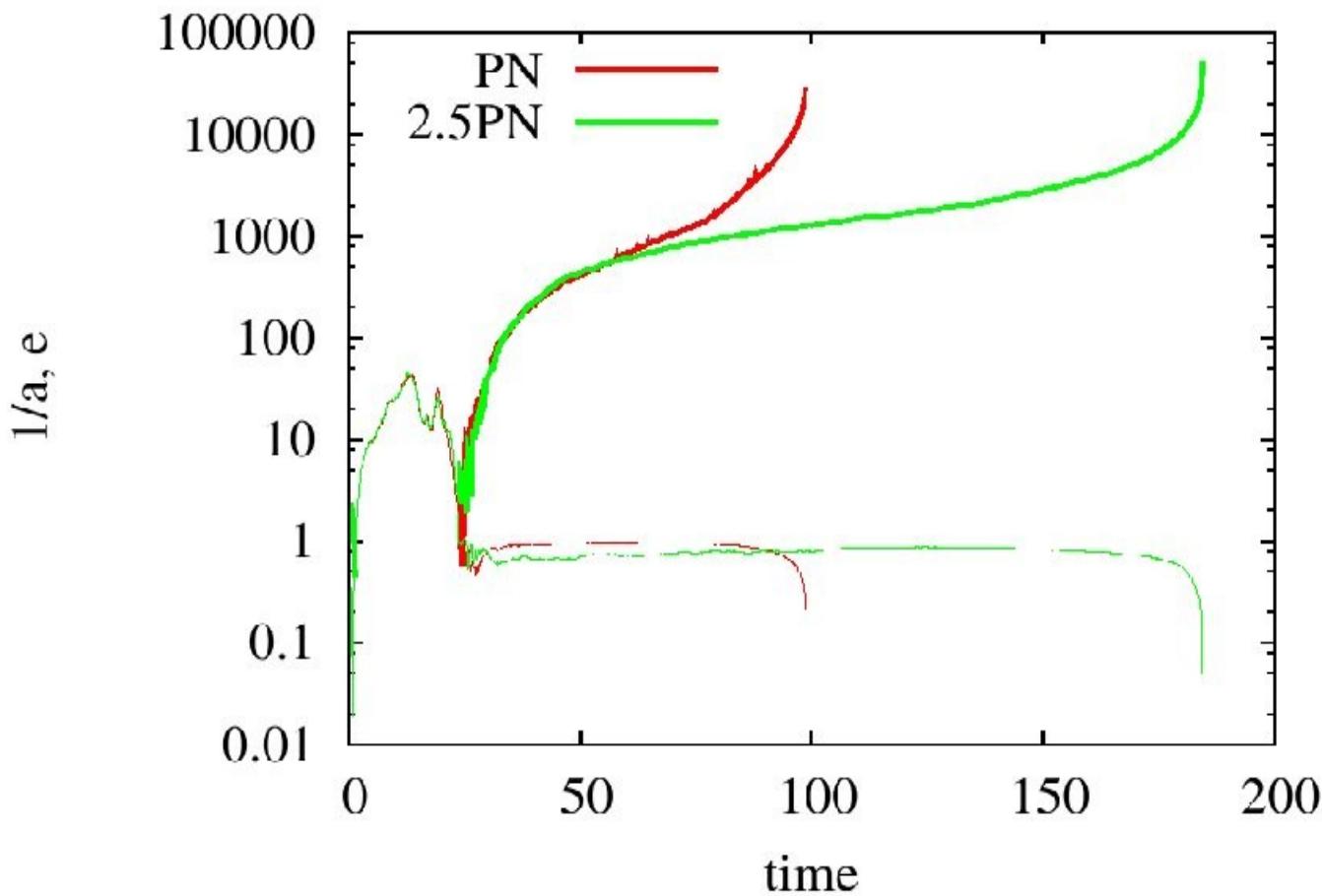
stars



40 Myr after galaxy merger

Simulations with
stars and dark matter,
Khan et al. 2012

stars



Berentzen, Preto,
Berczik, Merritt,
Spurzem,
2009, ApJ, up to
PN2.5

**Post-Newtonian
Dynamics... and they merge!**

M. Sobolenko, P. Berczik, R. Spurzem, G. Kupi, 2017;

M. Sobolenko, O. Kompaniiets, P. Berczik, V. Marchenko , A. Vasylchenko1, E. Fedorova,

B. Shukirgaliyev, 2022: NGC 6240 Supermassive Black Hole Binary dynamical evolution based on Chandra data

FAST COALESCENCE OF POST-NEWTONIAN SUPERMASSIVE BLACK HOLE

$$\mathbf{a}_{NoSpin} = \mathbf{a}_N + \frac{1}{c^2} \mathbf{a}_{1PN} + \frac{1}{c^4} \mathbf{a}_{2PN} + \frac{1}{c^5} \mathbf{a}_{2.5PN} + \frac{1}{c^6} \mathbf{a}_{3PN} + \frac{1}{c^7} \mathbf{a}_{3.5PN} + O\left(\frac{1}{c^8}\right) \quad (7)$$

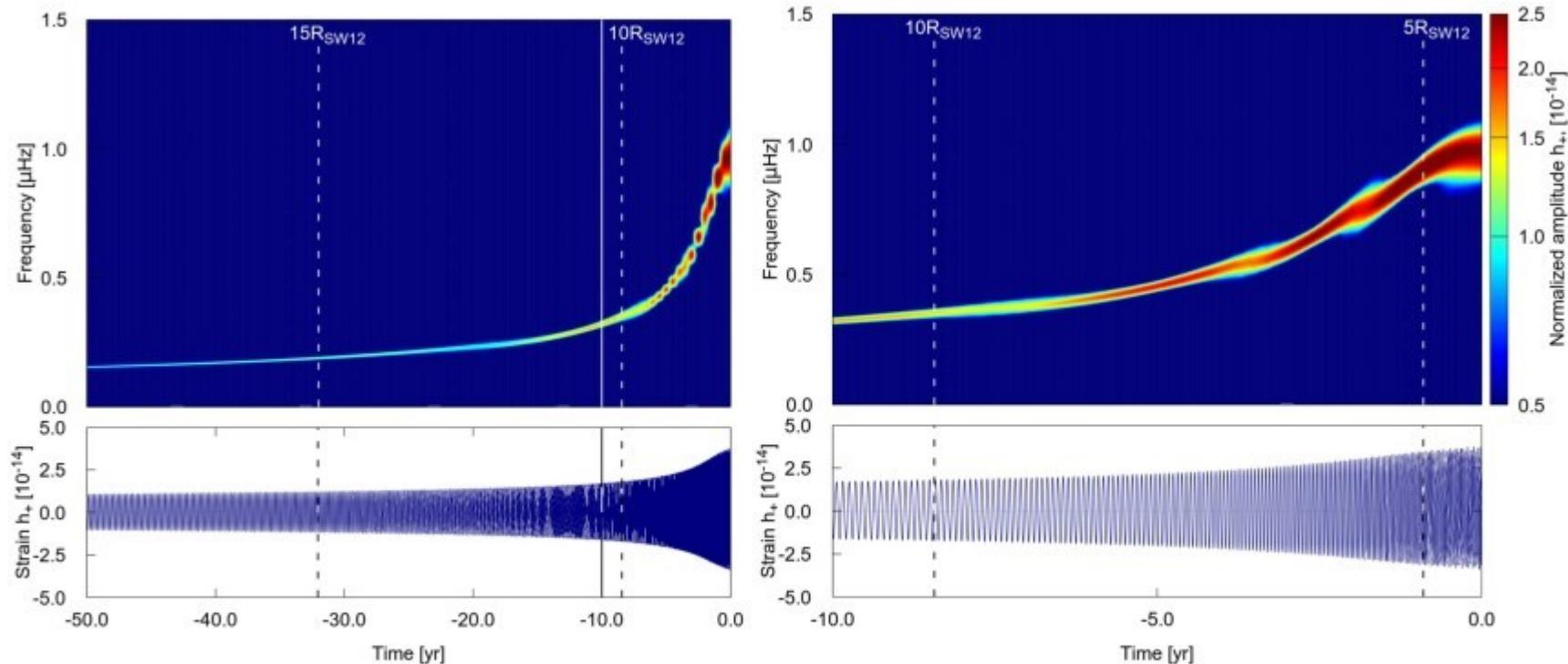


Figure 7. Time-frequency representations (top) of the strain data (bottom) for predicted gravitational waveforms of h_+ polarisation from SMBHB merging at NGC 6240 ($D_L = 111.2$ Mpc) for the last 50 yr (left) and last 10 yr (right). Major merging is represented by binary component with masses $1.36 \times 10^9 M_\odot$ and $6.8 \times 10^8 M_\odot$ and corresponding mass ratio 2:1. The final separation (due to our $\mathcal{P}N$ routine) is 0.75 mpc. The solid vertical line on the left panel indicates the last 10 yr of merging. Dashed vertical lines from left to right indicate binary separation 15, 10 and 5 Schwarzschild radii respectively.

- 1) Introduction, Globular Star Clusters
- 2) Nuclear Star Clusters
- 3) Code(s) and Hardware**

Code(s)

Nbody6++GPU

Nbody6++GPU

<https://github.com/nbody6ppgpu> (Spurzem & Kamlah 2023, LRCA)

The Hermite Scheme: 4th Order on two time points

$$\vec{a}_0 = \sum_j Gm_j \frac{\vec{R}_j}{R_j^3} ; \quad \vec{a}_0 = \sum_j Gm_j \left[\frac{\vec{V}_j}{R_j^3} - \frac{3(\vec{V}_j \cdot \vec{R}_j)\vec{R}_j}{R_j^5} \right] ,$$

$$\vec{x}_p(t) = \frac{1}{6}(t - t_0)^3 \vec{a}_0 + \frac{1}{2}(t - t_0)^2 \vec{a}_0 + (t - t_0) \vec{v} + \vec{x} ,$$

$$\vec{v}_p(t) = \frac{1}{2}(t - t_0)^2 \vec{a}_0 + (t - t_0) \vec{a}_0 + \vec{v} ,$$

Repeat Step 1 at t_1 using predicted $x, v \rightarrow a_1, \dot{a}_1$

Nbody6++GPU

$$\frac{1}{2}\vec{a}^{(2)} = -3\frac{\vec{a}_0 - \vec{a}_1}{(t - t_0)^2} - \frac{2\vec{a}_0 + \vec{a}_1}{(t - t_0)}$$

$$\frac{1}{6}\vec{a}^{(3)} = 2\frac{\vec{a}_0 - \vec{a}_1}{(t - t_0)^3} - \frac{\vec{a}_0 + \vec{a}_1}{(t - t_0)^2} ,$$

The Hermite Step
Get Higher Derivatives

$$\vec{x}(t) = \vec{x}_p(t) + \frac{1}{24}(t - t_0)^4 \vec{a}_0^{(2)} + \frac{1}{120}(t - t_0)^5 \vec{a}^{(3)} ,$$

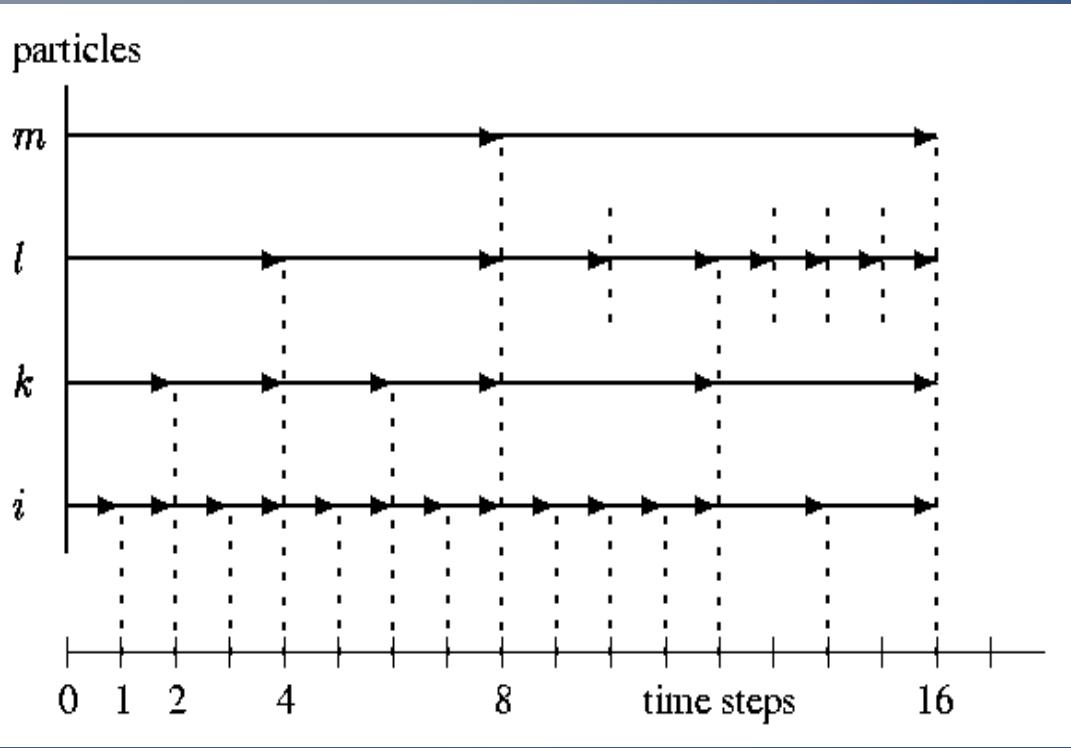
$$\vec{v}(t) = \vec{v}_p(t) + \frac{1}{6}(t - t_0)^3 \vec{a}_0^{(2)} + \frac{1}{24}(t - t_0)^4 \vec{a}_0^{(3)} .$$

The Corrector Step – this is not time symmetric!
(but can be made so by iteration if desirable)

Nbody6++GPU

*S.J.Aarseth, S. Mikkola,
J. Hurley, R. Spurzem,
L. Wang, ... (ca. 20.000 lines):*

- Hierarchical Block Time Steps
- Ahmad-Cohen Scheme
- Regularisations
- 4th order Hermite scheme
- NBODY6 (Aarseth 1999)
- NBODY6++ (Spurzem 1999) MPI
- NBODY6++GPU (Wang, Spurzem, Aarseth et al. 2015)



Hierarchical Block Time Steps

$$\Delta t = \sqrt{\eta \frac{|\vec{a}| |\vec{a}^{(2)}| + |\vec{\dot{a}}|^2}{|\vec{\dot{a}}| |\vec{a}^{(3)}| + |\vec{a}^{(2)}|^2}}.$$

NBODY1 – NBODY7: “The Growth of an Industry” (Aarseth 1999)

	ITS	ACS	KS	HITS	PN	AR	CC	MPI	GPU
NBODY1	✓								
NBODY2		✓		✓					
NBODY3	✓		✓						
NBODY4			✓	✓					
NBODY5	✓	✓	✓						
NBODY6		✓	✓	✓					
NBODY6GPU		✓	✓	✓				✓	
NBODY6++		✓	✓	✓			✓	✓	
NBODY6++GPU		✓	✓	✓	✓	✓	✓	✓	✓
NBODY7	✓	✓	✓	✓	✓	✓			

ITS: Individual time-steps [107]

ACS: Ahmad-Cohen neighbour scheme [109]

KS: KS-regularization of few-body subsystems [104]

HITS: Hermite scheme integration method combined with hierarchical block time-steps [111]

PN: Post-Newtonian [150, 125, 151]

AR: Algorithmic regularization [125]

CC: Classical chain regularization [114]

MPI: Message Passing Interface, multi-node multi-CPU parallelization [139]

GPU: use of GPU acceleration [138] (if also MPI: multi-node many GPU [144])

Berczik, Spurzem, et al., LNCS 2013; Table from: Spurzem, Kamlah 2023, LRCA

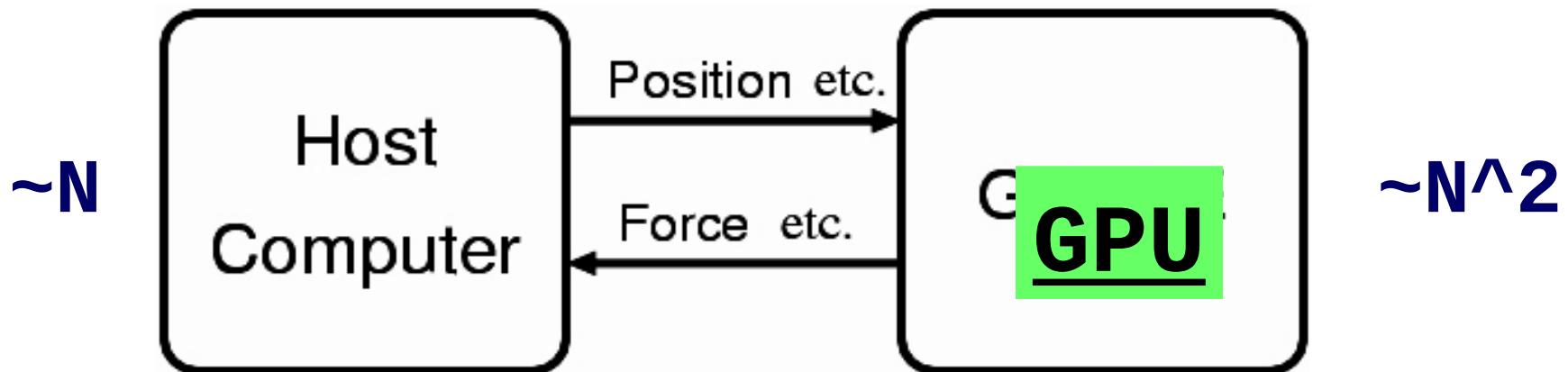
NBODY6++GPU: <https://github.com/nbody6ppgpu/>

Part of <https://www.punch4nfdi.de/> PUNCH4NFDI Consortium w. Jülich

Codes

Parallelization
Supercomputing
GPU Computing

Our own φ GRAPE/GPU N-body code



$$\vec{a}_i = \sum_{j=1; j \neq i}^N \vec{f}_{ij} \quad \vec{f}_{ij} = -\frac{G \cdot m_j}{(r_{ij}^2 + \epsilon^2)^{3/2}} \vec{r}_{ij}$$

NBODY6++GPU

<https://github.com/nbody6ppgpu/Nbody6PPGPU-beijing>

Picture by
Long Wang

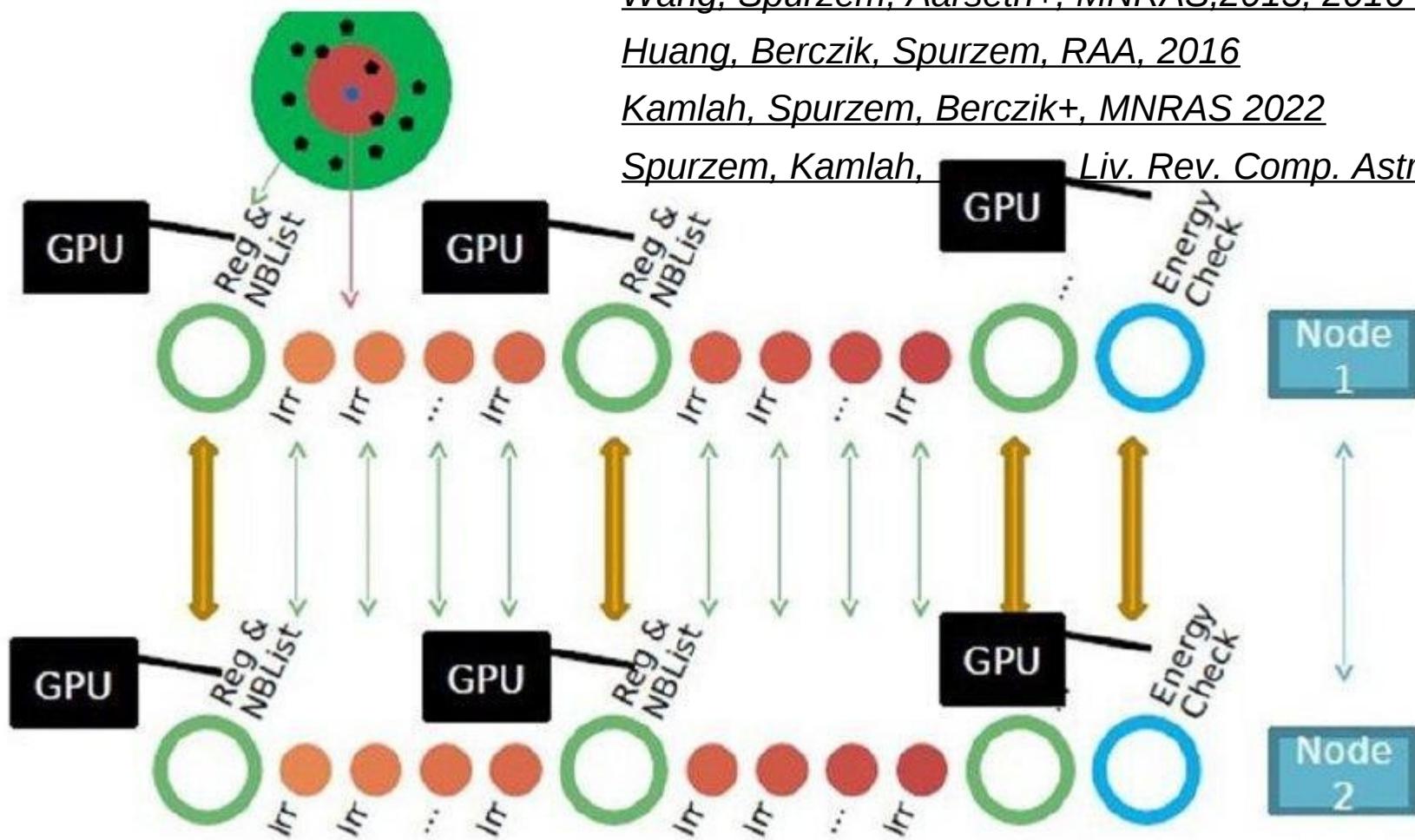
Our CPU/GPU N-body (AC) code

Wang, Spurzem, Aarseth+, MNRAS, 2015, 2016

Huang, Berczik, Spurzem, RAA, 2016

Kamlah, Spurzem, Berczik+, MNRAS 2022

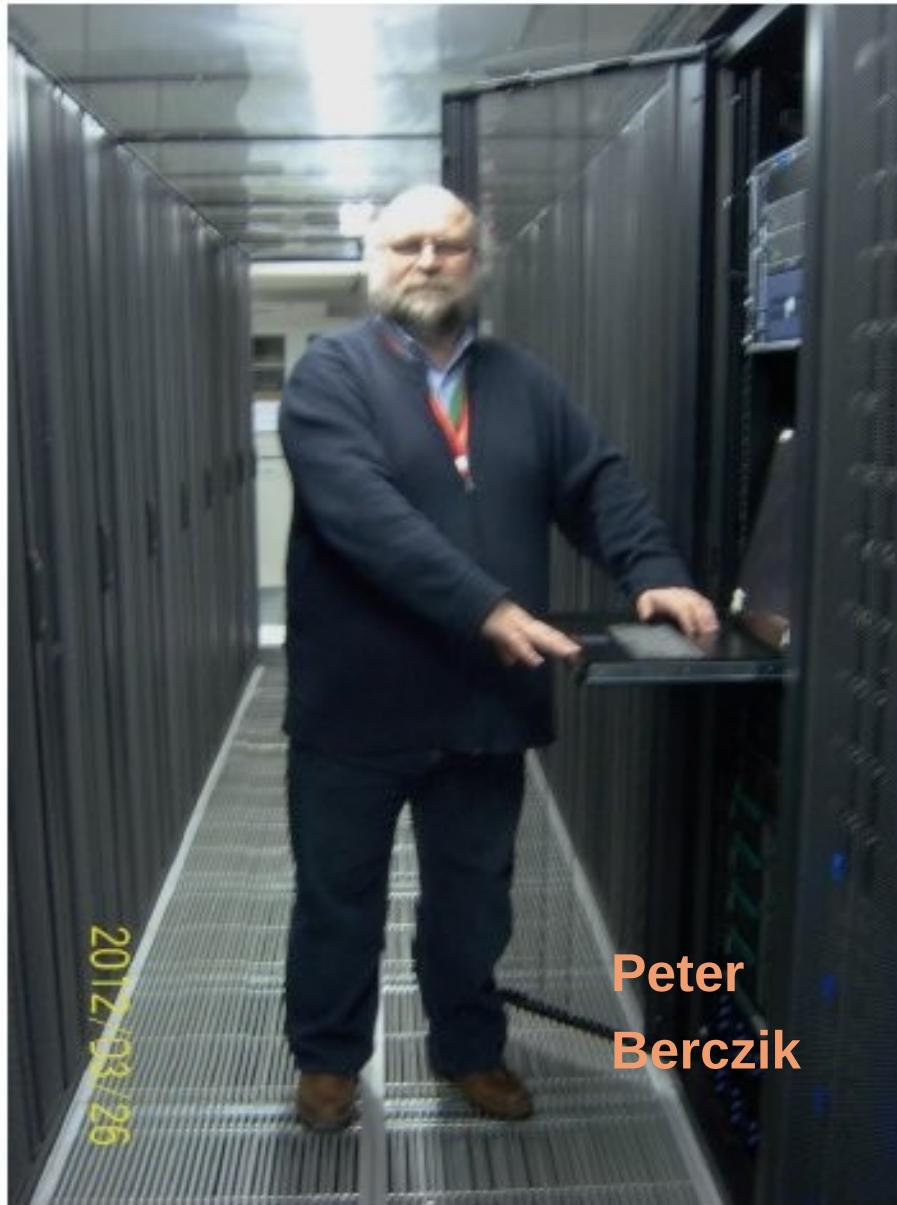
Spurzem, Kamlah, Liv. Rev. Comp. Astroph. 2023



Supercomputers

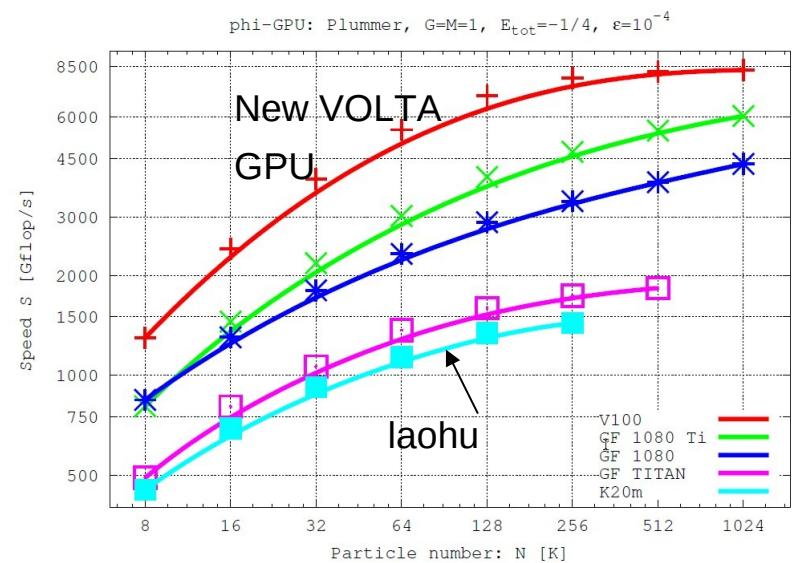
NAOC laohu cluster

64 Kepler K20



Laohu: 2009/2015
(Kepler GPU)
100 Tflop/s 150k cores

New GPUs 5-6 times
faster... (see below)



GPU Clusters used:

Kepler Cluster Heidelberg, Germany

(12x Kepler GPU, each ~2400 GPU cores)

Hansolo/Obiwan GPU Servers Beijing, China

(4x GTX 2080 GPU, each ~4300 GPU cores)

JUWELS Booster GPU Partition (Ampere A100 GPU)

(936 x 4 Ampere A100 GPU, each ~7000 GPU cores)

Raven GPU Cluster of Max-Planck Computing and Data Facility

MPCDF Garching



**Heidelberg
Germany**



**Hansolo @ NAOC
Silk Road Project**



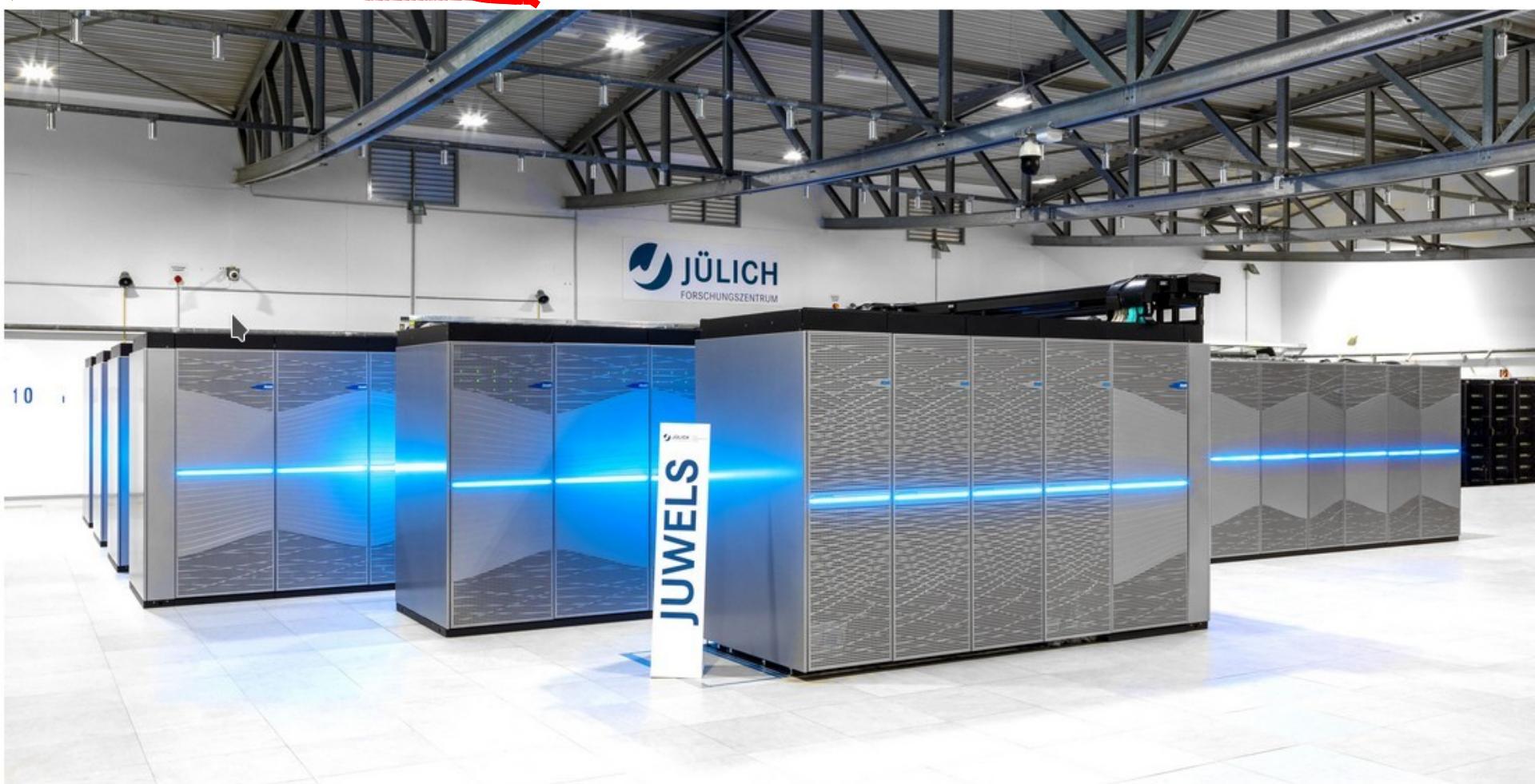
**JUWELS Juelich GPU
Cluster Germany**



**Peter
Berczik,
golowood
Cluster
Main Astron.
Observatory
Kiev, Ukraine**

JUWELS Booster 936 nodes (AMD CPU, 4x Ampere GPU)
~450.000 AMD cores, 25 million NVIDIA Ampere GPU cores
~ 70 Pflop/s SP ~ 44 Pflop/s DP
No. 12 in top500 list, No. 25 in green500 list

Jülich Wizard for European Leadership Science



LUMI

Supercomputer, Kajaani, Finland

Using only
Hydroelectric
Power and its
Heat used for heating
buildings.

No. 3 in top500
No. 7 in green500

2.2 million cores
10.000 AMD GPUs



EuroHPC and LUMI consortium:

Finland, Belgium, Czech Republic, Denmark, Estonia, Iceland, Norway,
Poland, Sweden, and Switzerland.

Summary and References

Summary Message

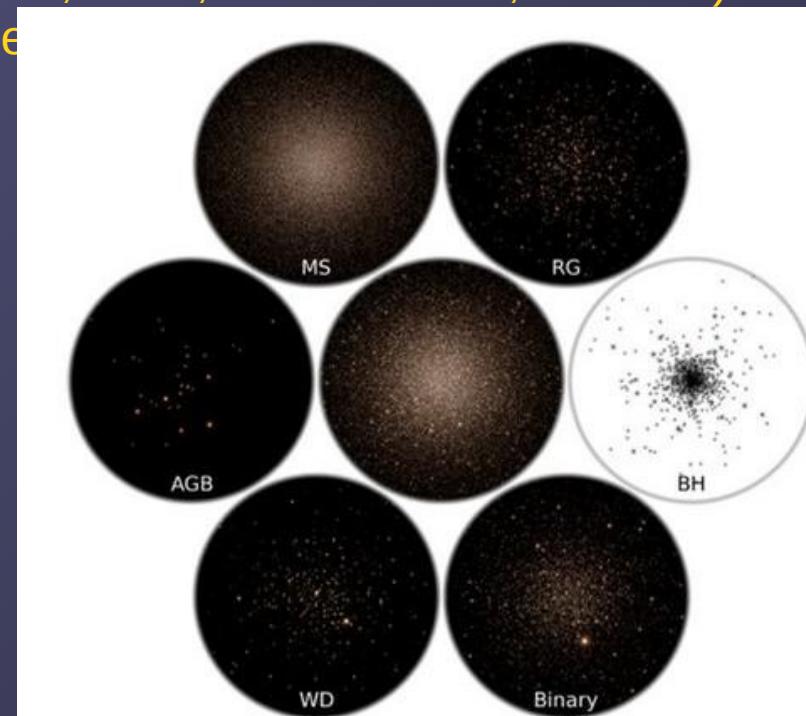
Massive Star Clusters:

- Direct N-Body Simulations of star clusters give LIGO/Virgo Sources (are consistent with them, it does not mean all sources are from star clusters)
- Necessary Input: single/binary stellar evolution / relativistic (PN) dynamics
- PN with $N > 2$
- Still very long computing time for few models (in spite of GPU, Lumi, ...)

Monte Carlo Models (MOCCA, Warsaw, M. Giersz, CMC, Northwestern, F. Rasio)
needed to get good sweep of parameter space

Nuclear Star Clusters:

- Observable Tidal Disruption Events (TDE)
- Light Curves correspond
to dynamics of TDE
- Future Work: relativistic inspirals
- Future Work: pulsars
- Future Work: star disk



Kazakhstan - China - Korea (KCK) becomes Silk Road Conference

Kazakhstan-China-Korea meeting becomes Silk Road Conference



13th Silk Road Conference: Dali, China, June 23-27, 2025 (just after the IAU Symposium Compact Objects and Binaries in Dense Star Clusters to be held in Seoul, Korea, June 16-20, 2025)

14th Silk Road Conference: Samarkand, Uzbekistan, June 2026 (to be confirmed)

Exploring the Frontiers of Dynamical Astronomy with High-Performance Computing, Artificial Intelligence, and Leading-Edge Observational Techniques.

We continue the numbering of these conferences, which started as Korea-China meeting in 2009; later it became Kazakhstan-China-Korea meeting, and due to increasing interest and participation of colleagues from all over central and east Asia we decided now to rename it to Silk Road Conference.

Last Conference: 12th Kazakhstan-China-Korea meeting in Astana, Kazakhstan, May 20-24, 2024

**Scientific
Organizing
Committee**

Hyung Mok Lee (Seoul National Univ.) Chair

Rainer Spurzem (Univ. of Heidelberg) Co-Chair

Fred Rasio (Northwestern Univ.)

Di Li (National Astronomical Observatory, Chinese Academy of Sciences)

Michiko Fujii (University of Tokyo)

Sourav Chatterjee (TIFR)

Anna Lisa Varri (Edinburgh Univ.)

Antonio Milone (Padova University)

<https://gravity.snu.ac.kr/iaus398/>

**IAU Symposium 398:
Compact Objects and
Binaries in Dense Stellar
Systems**

**Local Organizing
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DRAGON simulations – globular and nuclear star clusters

- **DRAGON simulation:** PhD thesis **Long Wang**, KIAA/PKU, awarded for first realistic globular cluster simulation using **NBODY6++GPU** with one million stars and many binaries (**Wang, Spurzem, Aarseth, et al.**, MNRAS 2016).
- The Dragon-II simulations – Paper III. Compact binary mergers in clusters with up to 1 million stars: mass, spin, eccentricity, merger rate and pair instability supernovae rate (**Arca Sedda, M., Kamlah, A. W. H., Spurzem, R., et al.**) arXiv e-prints arXiv:2307.04807, Paper II: MNRAS 525, 429 (2023), Paper I: arXiv e-prints arXiv:2307.04805

NBODY6++GPU and more, current state:

- **Spurzem, R., Kamlah A.W.H.** Direct N-body simulations, in **Living Rev. in Comp. Astrophysics 9, id.3 (2023)** (NBODY7 see also Banerjee, Sambaran papers)

Direct Nuclear Star Cluster Models with SMBH and TDE:

- DRAGON simulation of the Galactic Center, PhD thesis of Taras Panamarev, ARI/ZAH Univ. of Heidelberg (**Panamarev, Just, Spurzem, Berczik, Wang, Arca Sedda**, MNRAS 2019), simple TDE
- Revisit the Rate of Tidal Disruption Events: The Role of the Partial Tidal Disruption Event (**Zhong, S., Li, S., Berczik, P., Spurzem, R.**) 933, 96 (2022), TDE improved 1
- **Marija Minzburg, Philip Cho:** Master Theses Heidelberg 2023, publication in progress, TDE improved 2 and 3.

Some other papers and collaborators may be mentioned:

- Rizzuto, Naab, Spurzem et al. (2021, 2022) precursor of DRAGON II but no GW kicks, no TDE
- Li, Zhong, Berczik, Spurzem, Chen, Liu (MNRAS 2023 and earlier): merging nuclei with TDE