

The issue about time symmetry

1) Excursion: "leap frog" - simplest form $O(N^2)$
"natural" time symmetry, Hamilton

Hamiltonian relative two-body motion, gravitational:

$$H = \frac{1}{2} \mu v^2 - \frac{GM\mu}{r}$$

$v = |\vec{v}_1 - \vec{v}_2|$ relative vel.
 $r = |\vec{x}_1 - \vec{x}_2|$ "distance"

Separable:

$$H = H_1(v) + H_2(r)$$

m, M two masses

$\mu = mM/(m+M)$ reduced mass

$$p = \mu v, \quad H = \frac{p^2}{2\mu} - \frac{GM\mu}{r}$$

Conservative Kin. En. System pot. Energy

$$H = H_1 + H_2 = \frac{1}{2} \mu v^2 - \frac{GM\mu}{r}$$

canonical equations $\dot{r} = \frac{\partial H}{\partial v} \stackrel{\text{Sep.}}{=} \frac{\partial H_1}{\partial v} = \mu \cdot v = p$

$$\dot{v} = - \frac{\partial H}{\partial r} = - \frac{\partial H_2}{\partial r} = - \frac{GM\mu}{r^2}$$

Newton's Law
"Leap Frog Scheme"

Time Step: t_0, t_1 ; $\Delta t = t_1 - t_0$

$$\frac{v_1 - v_0}{\Delta t} = - \frac{\partial H_2}{\partial r} = - \frac{GM\mu}{r_{1/2}^2}$$

$$v_1 = v_0 - \Delta t \cdot \frac{GM\mu}{r_{1/2}^2} + O(\Delta t^3)$$

because of centering $r_{1/2}$

	⋮	
t_0	$r_{1/2}$	t_1
v_0	v_1	

$$v_1 = v_0 - \Delta t \frac{GMm}{r_{1/2}^2} + \mathcal{O}(\Delta t^3) \text{ centering}$$

$$r_{3/2} = r_{1/2} + \Delta t \cdot v_1 + \mathcal{O}(\Delta t^3)$$

$$v_1 = v_0 - \Delta t \left. \frac{\partial H_2}{\partial r} \right|_{r=r_{1/2}}$$

$$r_{3/2} = r_{1/2} + \Delta t \cdot \left. \frac{\partial H_1}{\partial v} \right|_{v=v_1}$$



direct discretization of Hamiltonian equations.

Can only do Leap Frog

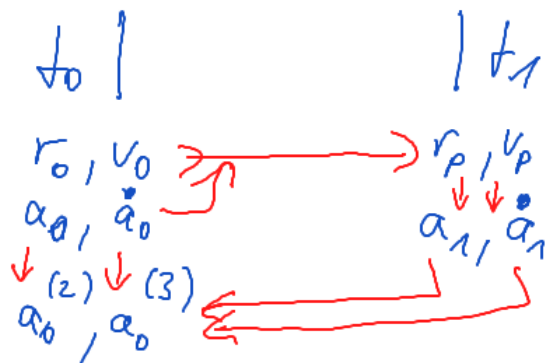
because H is separable!

Solutions of conservative Hamiltonian systems are time-symmetric!
 Numerical Leap Frog is exactly " " " " $\Delta t' = -\Delta t$

- Time Symmetry is nice feature of Leap Frog
- Numerical solution, using leap frog, is an exact solution of an approximate Hamiltonian $H = H_0 + \delta t \cdot H_1 + \dots$
 $t_0 \rightarrow t_1$; δt constant timestep

Heunike Scheme not time symmetric

has a global energy error
 Keep very small!



*- cluster simulations tolerated!

Trust

Q + Tremaine 1992

• Two possible ways out:

a) higher order "leap frog" or similar ("symplectic" integrator)
BIFROST Rautala, Naab et al. 2020, ---

b) make Hermite scheme symmetric!
Kohno + Makino (2004)

Classical Hermite corrector:

$$\vec{r}_c(t_1) = \vec{r}_p(t_1) + \frac{1}{24} \Delta t^4 \vec{a}_0^{(2)} + \frac{1}{120} \Delta t^5 \vec{a}_0^{(3)} + \mathcal{O}(\Delta t^6)$$

$$\vec{v}_c(t_1) = \vec{v}_p(t_1) + \frac{1}{6} \Delta t^3 \vec{a}_0^{(2)} + \frac{1}{24} \Delta t^4 \vec{a}_0^{(3)} + \mathcal{O}(\Delta t^5)$$

Hermite Formula $\vec{a}_0^{(2)}, \vec{a}_0^{(3)}$

Strictly
Time
Symmetry

$$\vec{r}_c(t_1) = \vec{r}_0(t) + \frac{1}{2} \Delta t (\vec{v}_c + \vec{v}_0) - \frac{1}{12} \Delta t^2 (\vec{a}_1 - \vec{a}_0) + \mathcal{O}(\Delta t^5)$$

$$\vec{v}_c(t_1) = \vec{v}_0(t) + \frac{1}{2} \Delta t (\vec{a}_1 + \vec{a}_0) - \frac{1}{12} \Delta t^2 (\vec{a}_1 - \vec{a}_0) + \mathcal{O}(\Delta t^5)$$

t_0
|
 r_0, v_0
 a_0, \dot{a}_0

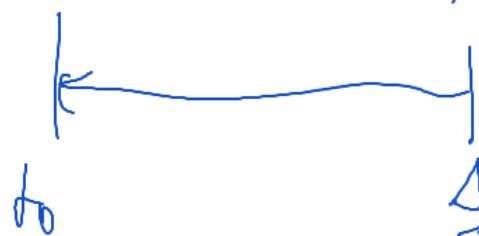
t_1
|
 a_1, \dot{a}_1
 r_c, v_c

$\Delta t' = -\Delta t$
Exchange 0, 1
 $r_0, r_c ; v_0, v_c$

- Was proposed for high accurate orbit integrations suitable for planetary systems, particles in protoplanetary disk.
- Caveat: Whenever time step varies \leftrightarrow symmetry is broken



$$\Delta t_{\rightarrow} = \Delta t(a_0, \dot{a}_0, a_0^{(2)}, a_0^{(3)})$$



$$\Delta t_{\leftarrow} = \Delta t(a_1, \dot{a}_1, a_1^{(2)}, a_1^{(3)})$$

Currently "time-symmetric" Hermite is not used: due to problems with variable time steps and Ahmad-Cohen neighbour scheme