



中国科学院国家天文台

NATIONAL ASTRONOMICAL OBSERVATORIES, CHINESE ACADEMY OF SCIENCES



University of Heidelberg, Germany



the SILK ROAD PROJECT at NAOC
丝绸之路计划

DRAGON star cluster simulations, and how to form an intermediate mass black hole

Rainer Spurzem, and Silk Road Team

(mainly: Kai Wu, Philip Cho, Vahid Amiri, Marcelo Vergara, Francesco Flammini Dotti, Shuo Li, ...)

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ARI Colloquium

July 24, 2025





Sverre Aarseth

20 July 1934 – 28 December 2024

<https://www.ast.cam.ac.uk/news/sverre-aarseth-20-july-1934-28-december-2024>



Picture:

Xi Shuang Banna,
Yunnan, SW China
(R.Sp.)

90's at Institute of Astronomy, Cambridge, UK

Left: Emanuel Vilkoviski,
Pavel Kroupa, Sverre Aarseth, R.Sp.

Below:
Chris Tout, E. Vilkoviski,
Sverre Aarseth





Binary Parallelization
in Norway 2001 – not
yet successful

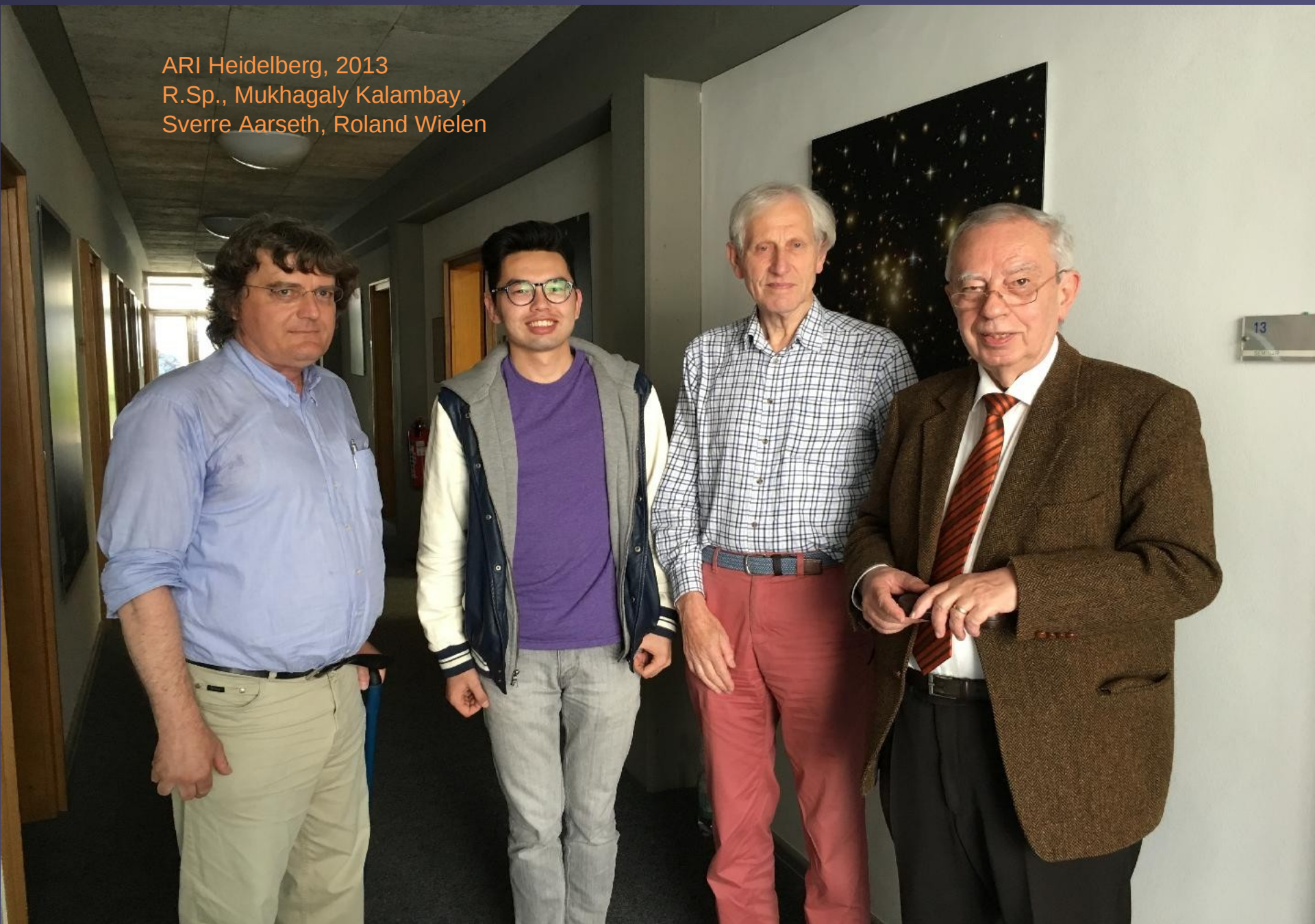


Sambaran Banerjee, Ann-Marie Madigan, Sverre Aarseth, Alexei Minz, Oliver Porth, R. Spurzem
Avetis Sadoyan, Matthew Benacquista

MODEST 8a
Heidelberg 2008



ARI Heidelberg, 2013
R.Sp., Mukhagaly Kalambay,
Sverre Aarseth, Roland Wielen





Further Development of Chain Regularization

IAU Symp. No. 246
Capri, Italy, 2007

Top:
Sverre Aarseth,
Seppo Mikkola

Right:
Sverre Aarseth
Pavel Kroupa
Paulina Assmann



Heidelberg 2018
Work Life Balance
(Pic: A. Just)



Aarseth, Henon, Wielen, 1974, A&A: A comparison of numerical methods for the study of star cluster dynamics

N-body Codes:

Von Hoerner
Wielen
Aarseth

independent

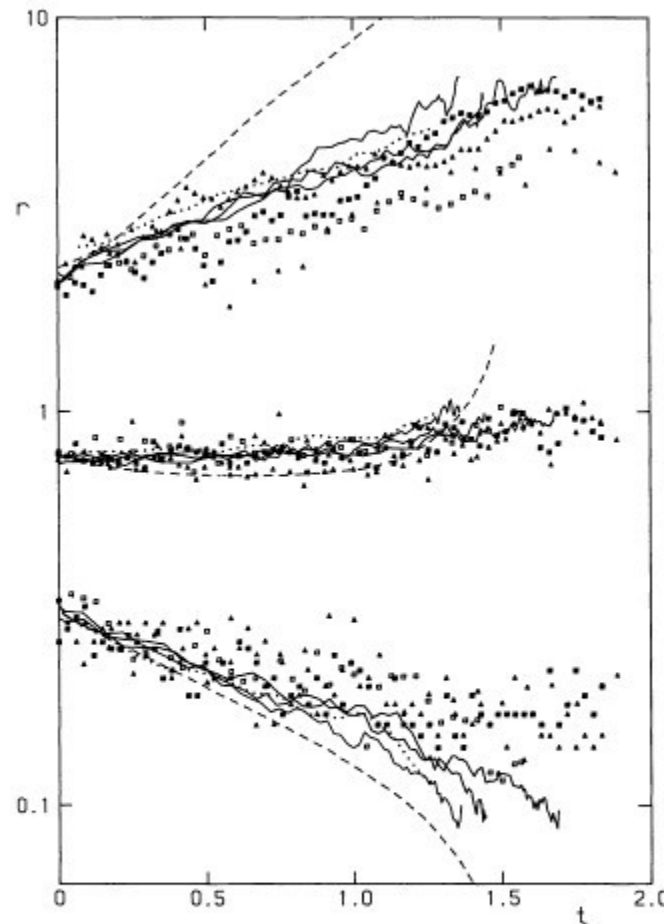


Fig. 1. Radii containing 10%, 50%, 90% of the mass, plotted versus time for a cluster with stars of equal masses. Open triangles and squares: N -body integrations with $N=100$ and $N=250$ (Wielen). Filled triangles and squares: N -body integrations with $N=250$ (Aarseth). Full lines: Monte Carlo models (Hénon). Dotted lines: Monte Carlo model (Shull and Spitzer). Dashed lines: fluid-dynamical model (Larson)

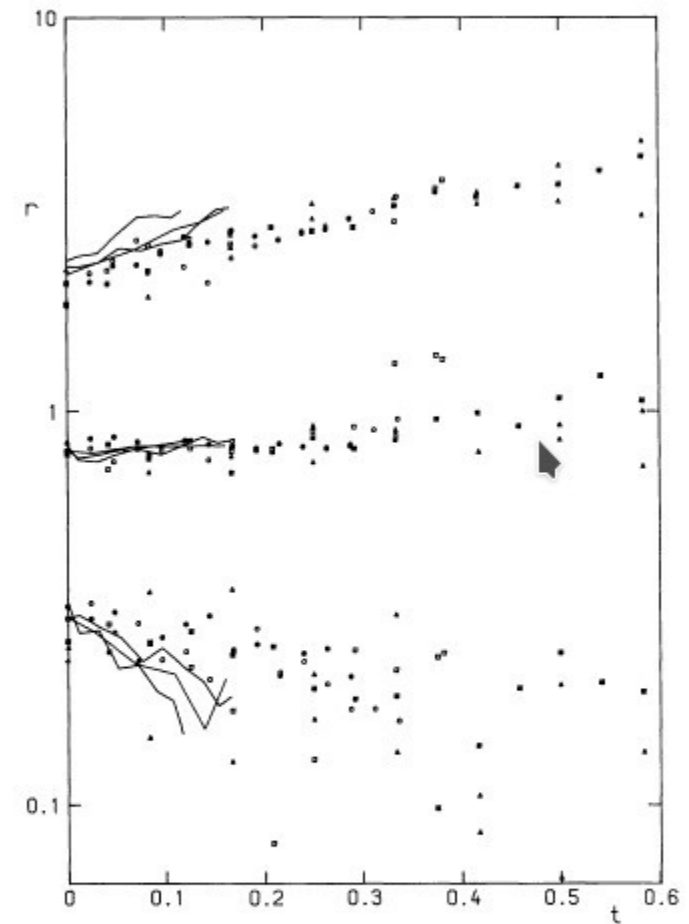


Fig. 2. Radii containing 10%, 50%, 90% of the mass, plotted versus time for a cluster with stars of different masses. Open and filled symbols: N -body integrations with $N=100$ (triangles), $N=250$ (squares), $N=500$ (circles) (Wielen). Full lines: Monte Carlo models (Hénon)

- 1) Star Cluster Simulations/
Black Holes/Grav. Waves
- 2) Nuclear Star Clusters
- 3) Code(s) and Hardware
- 4) Summary and References

Introduction

**Kavli Institute for Astronomy
and Astrophysics, Peking Univ..**



Founded 2008



Astronomisches
Rechen-Institut (ARI)
Univ. of Heidelberg,
Germany

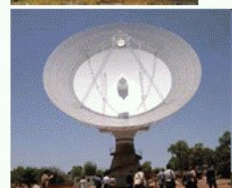
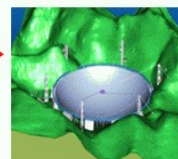
Founded May 10, 1700



中国科学院国家天文台

NATIONAL ASTRONOMICAL OBSERVATORIES, CHINESE ACADEMY OF SCIENCES

NAOC/
CAS



**Silk Road Project =
Computational Science Project...**

Top: NAOC Headquarter Beijing
Bottom: LAMOST Site



(Credit: X-ray: NASA/CfA/J. Grindlay et al.,
Optical: NASA/STScI/R. Gilliland et al.)

X-ray binaries

with neutron stars
and black holes

Globular Cluster 47 Tuc

~ one million stars

$$\vec{a}_0 = \sum_j Gm_j \frac{\vec{R}_j}{R_j^3} ; \quad \ddot{\vec{a}}_0 = \sum_j Gm_j \left[\frac{\ddot{\vec{V}}_j}{R_j^3} - \frac{3(\ddot{\vec{V}}_j \cdot \vec{R}_j)\vec{R}_j}{R_j^5} \right]$$

Dynamics of Star Clusters

- $N = \infty$

negative specific Heat

gravothermal Collapse

gravothermal Oscillations

- $N = 3$ ($N = 2, \dots, \approx 100$)

History

Exponential Instability

Chaos and Resonance

Regularisation

- $N = 10^6$ ($N = 10^4, 10^5$)

Post-Kollapse-Evolution

Binaries

Globular Clusters

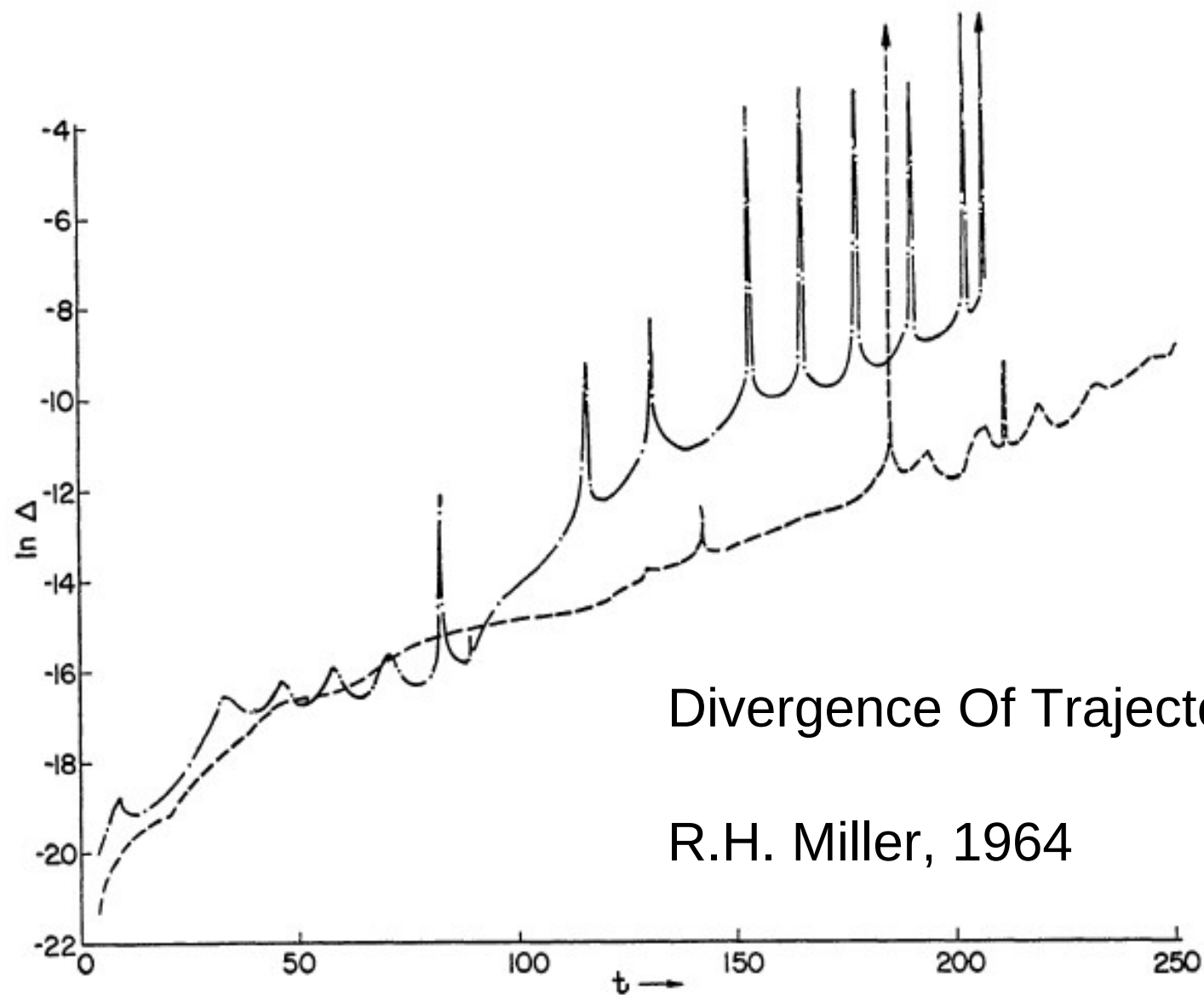


FIG. 2.—Tracks of the separation of phase points for 8-particle stellar dynamical systems. The dashed curve was based on the initial conditions of Section II, while the solid curve results from a different calculation starting from a system that had relaxed from the original atypical condition. The repeated equidistant spikes between $t = 150$ and $t = 220$ are formed by a single long-lived binary pair, which is destroyed in a close collision of one of its members with another star at $t = 225$.

Concept of Shadow Orbits

Quinlan & Tremaine 1992

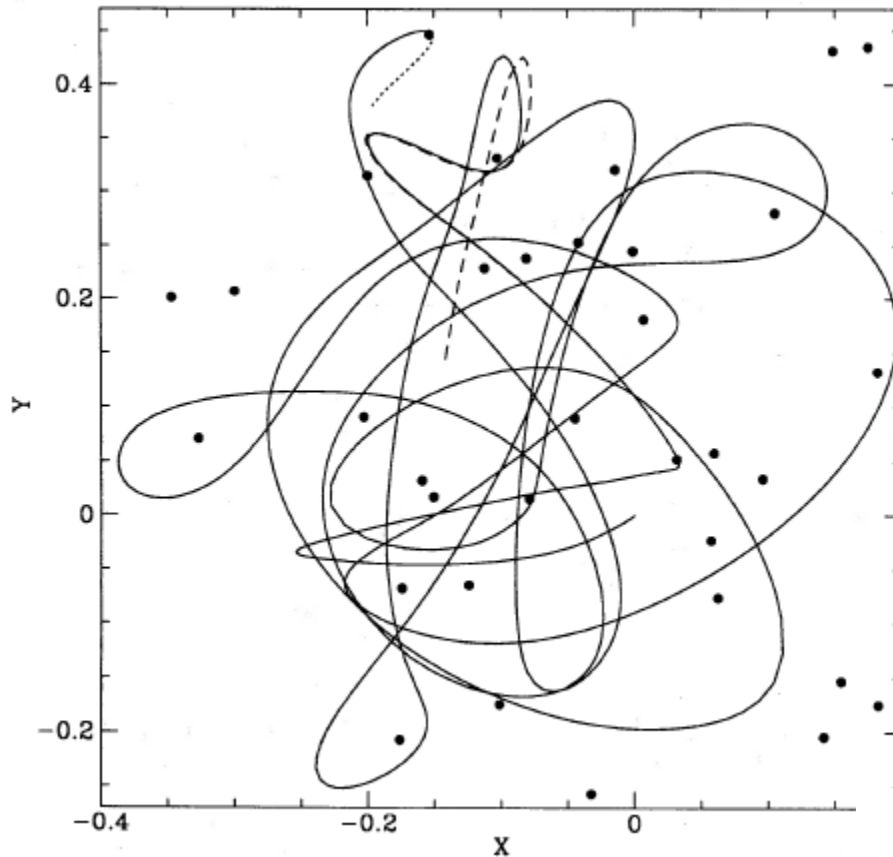


Figure 3. An example of a shadow orbit. The solid line is the projection on to the x - y plane of the noisy orbit of a particle starting at the origin of the cluster in Fig. 1 with velocity $(v_x, v_y, v_z) \approx (-0.57, -0.46, 0.56)$ integrated for 20.1 time units with accuracy parameter $\eta = 0.027$. The filled circles are the projected positions of the fixed particles in the cluster. The dashed line is an accurate Bulirsch-Stoer integration starting from the same initial conditions, integrated for 3.25 time units. The dotted line is the shadow orbit. The phase-space separation between the shadow and noisy orbits reaches a maximum of 0.042 (resulting mostly from the separation in velocity, not position) near the point $x = 0.033, y = 0.051$. The shadow orbit can be seen extending from the end of the noisy orbit at the top of the figure (because the shadow orbit has been plotted for a slightly longer time interval).

Concept of Shadow Orbits

Quinlan & Tremaine
1992

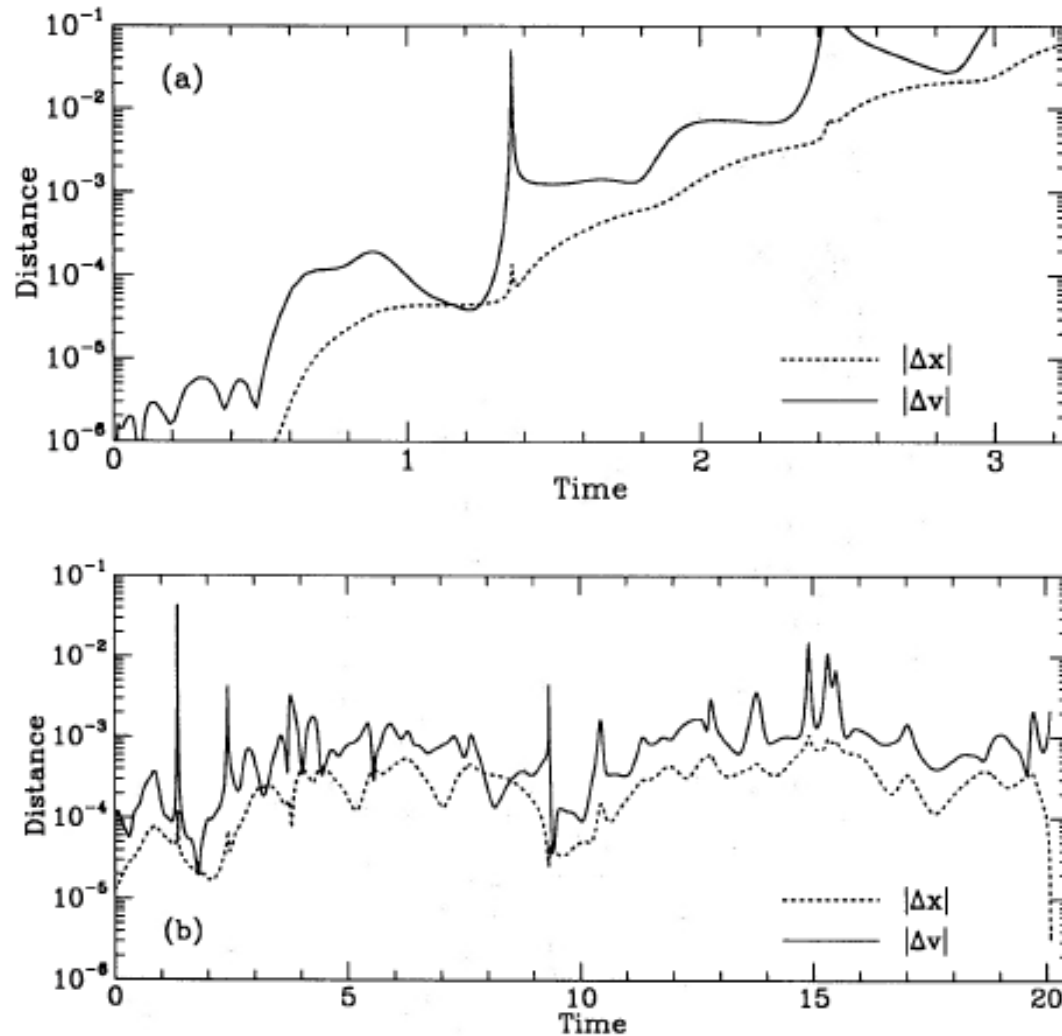
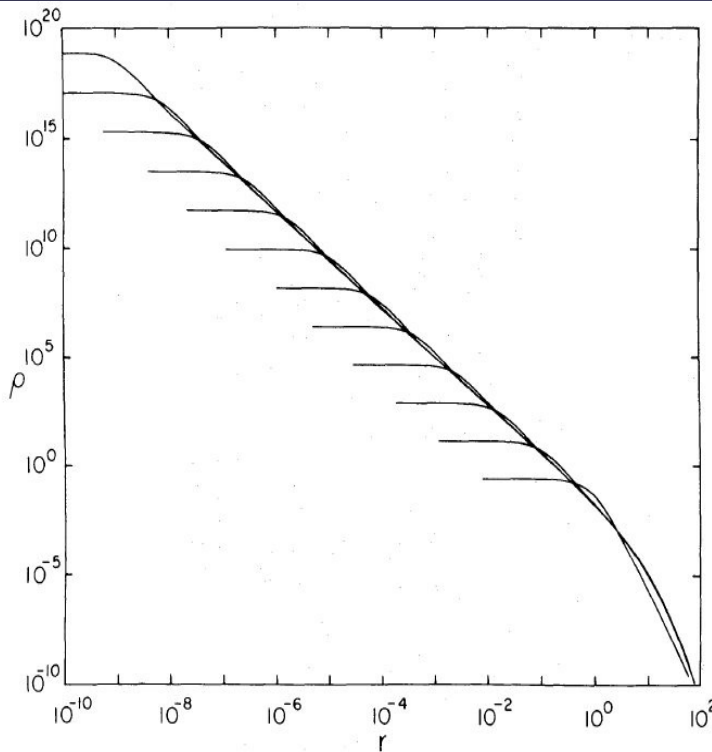


Figure 4. Phase-space separation between the orbits of Fig. 3: (a) separation between the noisy and accurate orbits; (b) separation between the noisy and shadow orbits.

Dynamics of Star Clusters: Gas Sphere Model



Cohn (1980): Direct Fokker-Planckj model

Core Collapse

Gravothermal Catastrophe

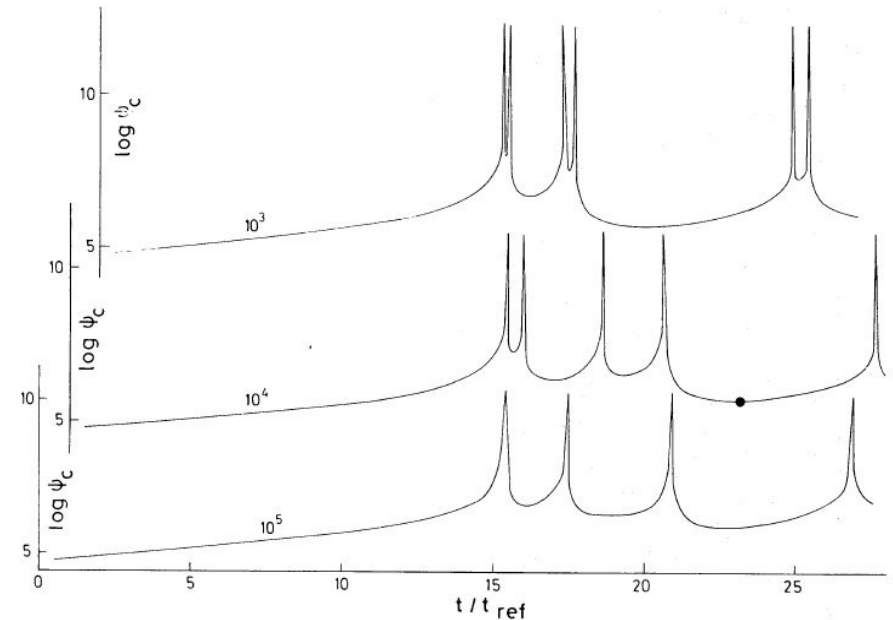


Figure 1. The 'central' density ψ_c is plotted against the non-dimensional time t/t_{ref} for $k = 2$ models with three different values of C as attached to each curve. Note, that if they were plotted with the same ordinate they would be close to each other despite the great differences in C . The model indicated with a filled circle will be compared with King's model in Section 4.2.

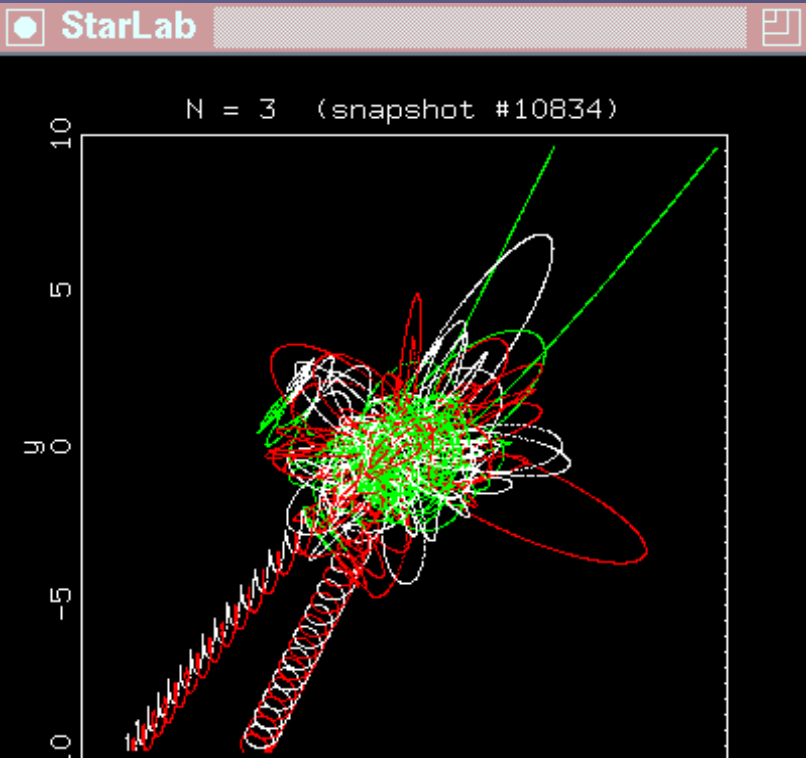
Bettwieser & Sugimoto 1984:
Gravothermal Oscillations by
energy generation from binaries
(cf. nuclear stellar energy generation)

3-body Encounters Starlab

Simulation (S.L.W. McMillan)

<http://www.physics.drexel.edu/~steve/>

-> Three-Body-Problem



Gravothermal Oscillations - Attractor in Phase Space

Spurzem 1994, Giersz & Spurzem 1994
Amaro-Seoane, Freitag & Sp. 2004

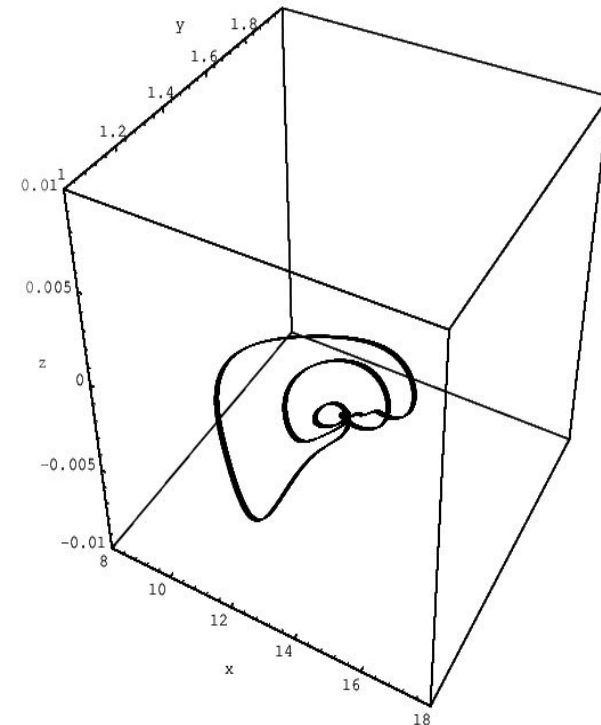


Fig. 3:

4-dimensional attractor for $N = 100.000$ system, $x = \sigma'_c$, $z = \xi$.

Gravothermal instability of anisotropic self-gravitating gas spheres: singular equilibrium solution

R. Spurzem^{1,2,3}

¹Institut für Theoretische Physik und Sternwarte, University of Kiel, Olshausenstraße 40, D-2300 Kiel, Germany

²Institut für Astronomie und Astrophysik, University of Würzburg, Am Hubland, D-8700 Würzburg, Germany

³Universitätssternwarte Göttingen, Geismarlandstraße 11, D-3400 Göttingen, Germany

**Follow-Up of Angeletti &
Giannone and Larson**

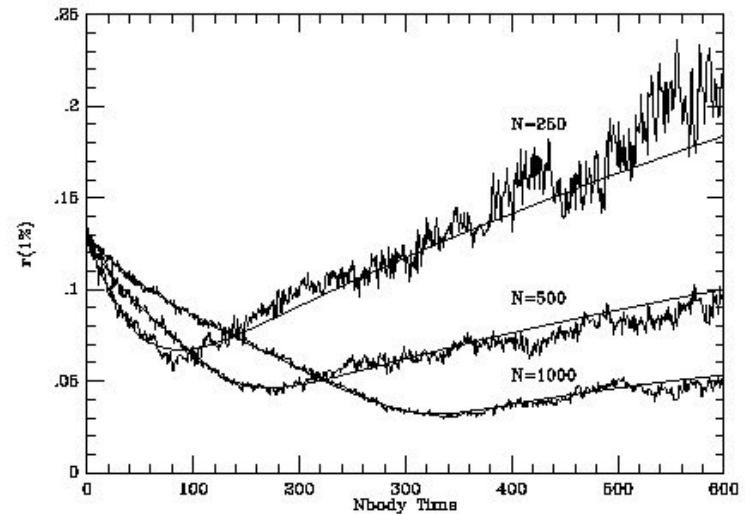
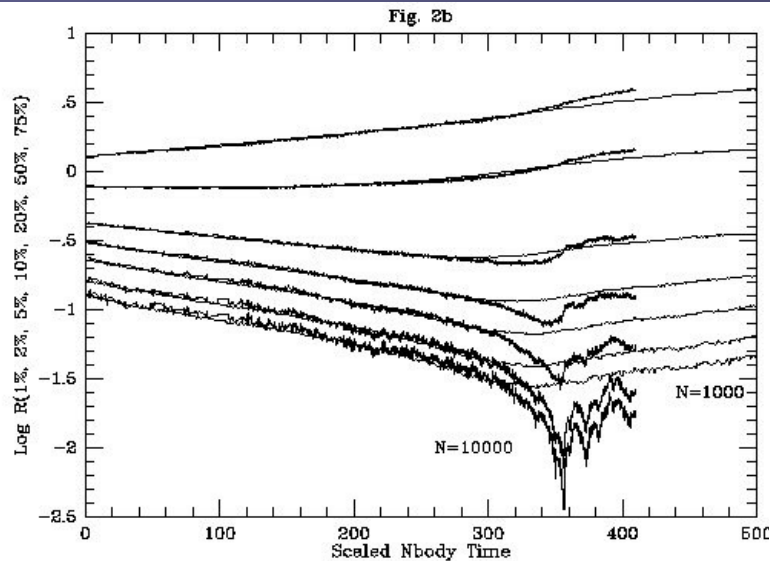
Dynamics of Star Clusters

(compare gaseous model with direct N-body integration)

(Spurzem & Aarseth 1996)

(Giersz & Spurzem 1994)

(now in Binney/Tremaine)



N-Body / N-Body

N-Body / Fokker-Planck

In spherical symmetry

...but...

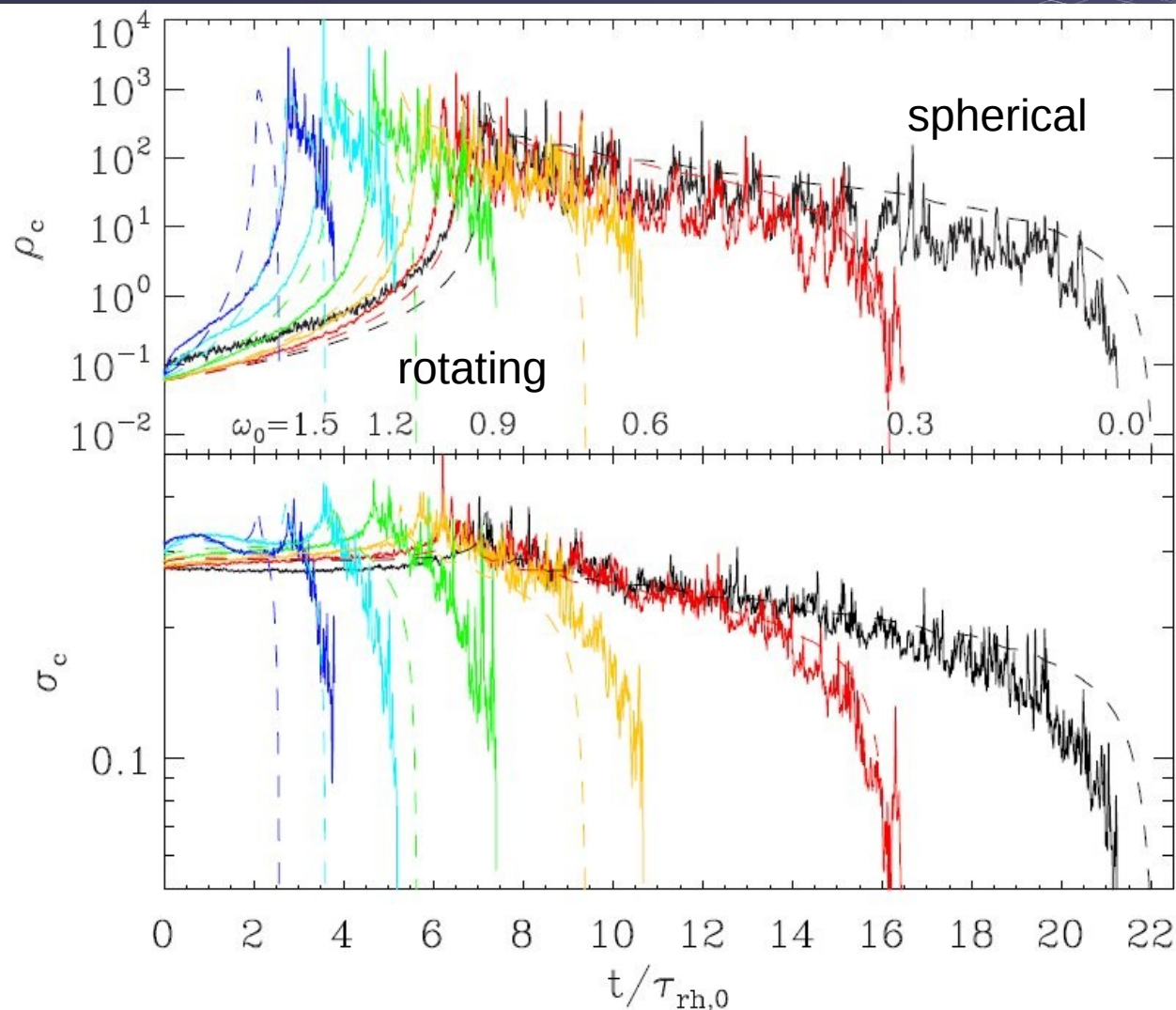
Dynamics of (Rotating) Star Clusters

Fokker-Planck N-Body Comparison

Dissolution of Star Cluster in Tidal Field

Kim, Yoon,
Lee, Spurzem,
2008, MNRAS

Hong, Kim,
Lee, Spurzem,
2013, MNRAS



Three Phases in
Cluster Dissolution:
1) Core Collapse
(Encounters)
2) Post-Collapse
Steady Evaporation
(Encount)
3) Dynamic
final dissolution

DRAGON I Simulation

DRAGON Simulation

<http://silkroad.zah.uni-heidelberg.de/dragon/>

<https://github.com/nbody6ppgpu>

Also in: <https://www.punch4nfdi.de/>



One million stars direct simulation,

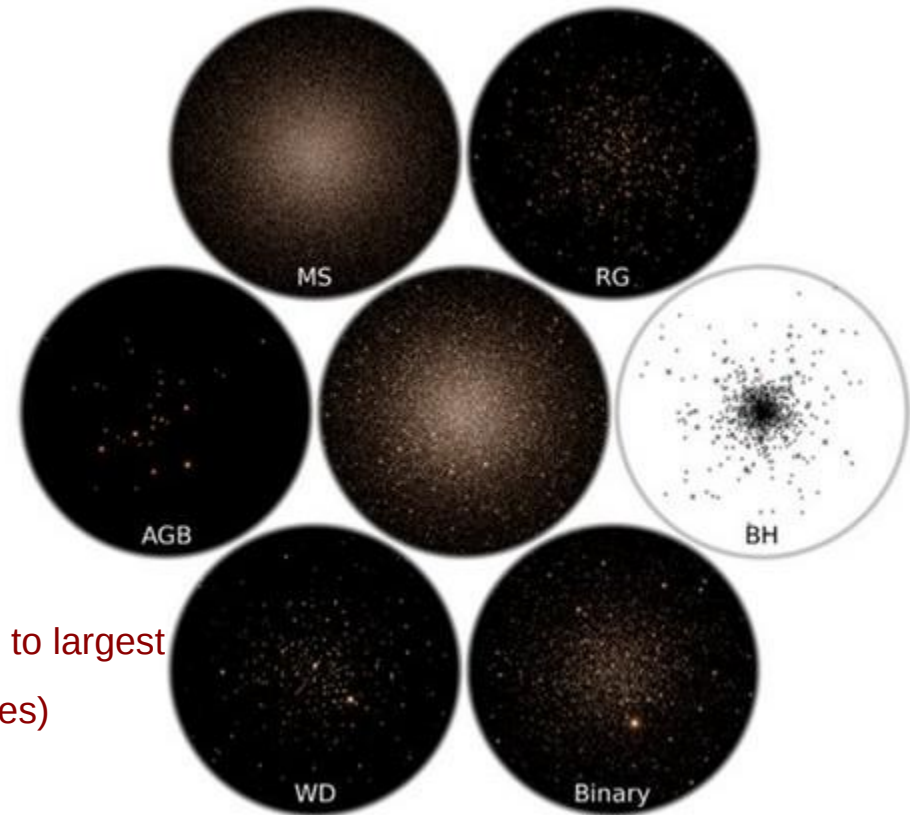
biggest and most realistic direct N-Body simulation of globular star clusters.

With stellar mass function, single and binary stellar evolution, regularization of close encounters, tidal field (NBODY6++GPU).

(NAOC/Silk Road/MPA collaboration).

Wang, Spurzem, Aarseth, Naab et al. MNRAS, 2015

Wang, Spurzem, Aarseth Naab, et al. MNRAS 2016

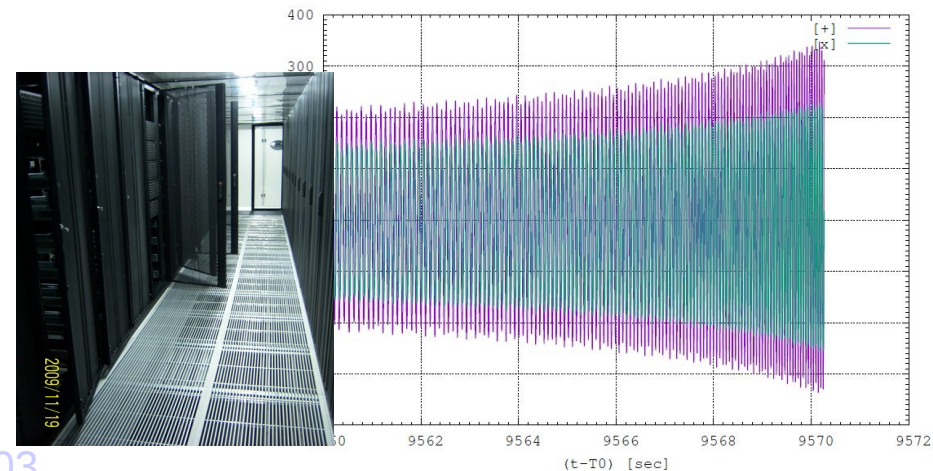
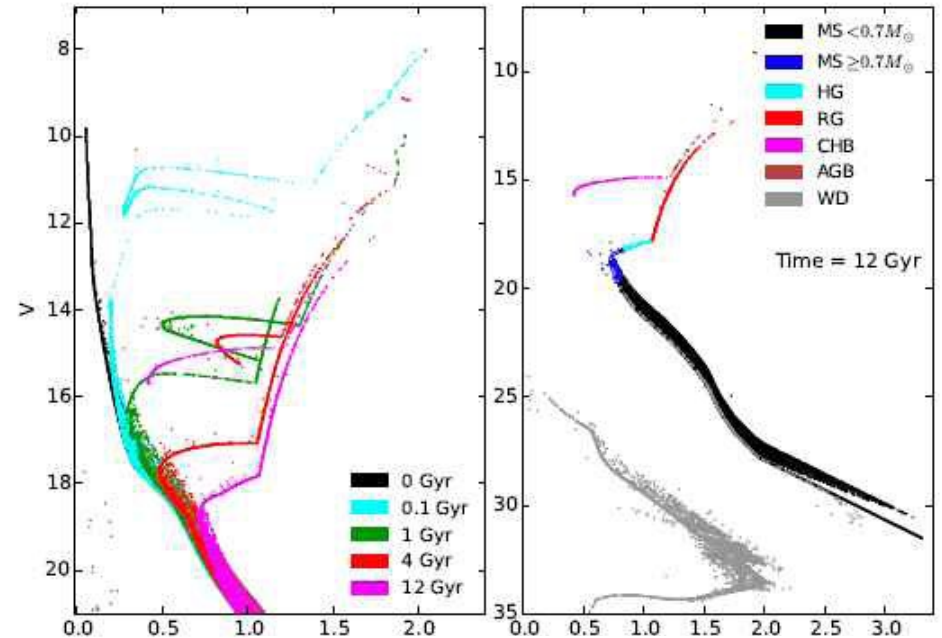


Number of Floating Point Operations (~1M bodies) similar to largest
Cosmological simulations (Millennium, Illustris, $\sim 10^{10}$ bodies)

天龙星团模拟：百万数量级恒星、黑洞和引力波

Dragon Star Cluster Simulations: Millions of Stars; black holes and gravitational waves

- First realistic globular star cluster model with million stars (*Wang, Spurzem, Aarseth, ..., Berczik, Kouwenhoven, ... MNRAS 2015, 2016*)
- Synthetic CMD (right side) with zero photometric errors, different ages shown
- Black hole binary mergers occur as observed by LIGO. Our grav. waveforms computed from simulation (right side). (Only inspiral plotted not ringdown.)
- GPU accelerated supercomputers laohu in NAOC and hydra of Max-Planck (MPCDF) in Germany needed!



DRAGON II Simulations

DRAGON-II Simulations – Paper III

using NBODY6++GPU

*Arca Sedda et al. 2023ab,
2024: MNRAS*

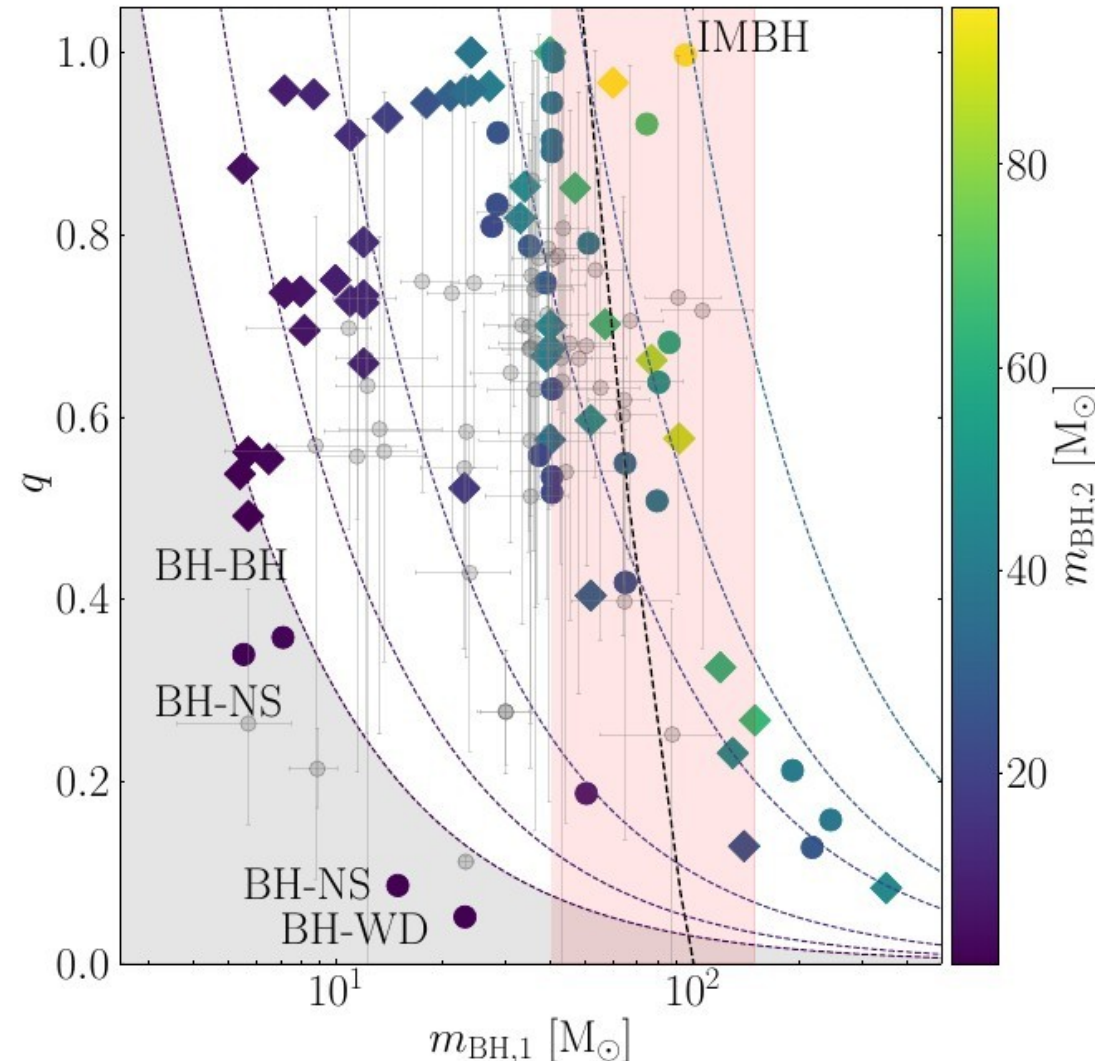
*19 models, up to 1 million stars,
up to 33% initial hard binaries*

Compact Object Mergers
Compared with LIGO-Virgo
GWTC-3 catalogue (grey
symbols)

Mass ratio q vs. primary mass
 m_1 ; colour code:
secondary mass m_2

grey shade: neutron star
involved

red Shade: mass gap



DRAGON I Simulation: Wang, Spurzem, Aarseth, et al. 2015, 2016

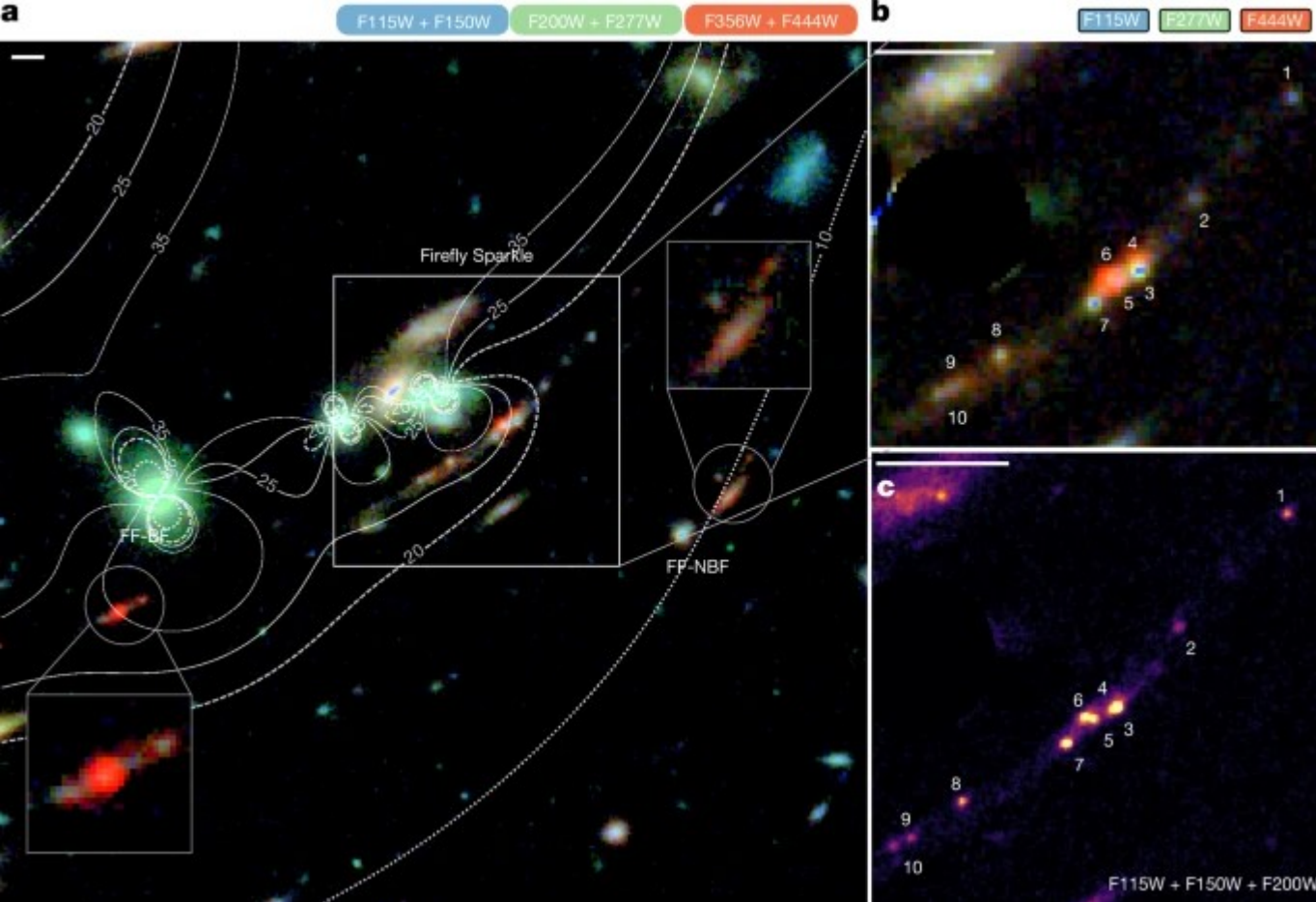
DRAGON II Simulation: Arca Sedda, Kamlah, Spurzem, et al. 2023, 2024ab

DRAGON III Simulation: Wu et al. 2025, in prep., and DRAGON Data Release

The Future - DRAGON III 1m – 8m (16m?)

ARI Colloquium Kai Wu, July 10, 2025

Young Dense Clusters



Globular Clusters in the Young Universe I

c, Combined short wavelength (F115W + F150W + F200W) image of the Firefly Sparkle, in which the distinct clusters can be seen. Scale bars, 1".

NATURE 636, pages 332–336 (2024), Mowla et al. Dec. 2024: The most distant galaxies detected were seen when the Universe was a scant 5% of its current age. At these times, progenitors of galaxies such as the Milky Way were about 10,000 times less massive. ... Here we present JWST observations of a strongly lensed galaxy at $z_{\text{spec}} = 8.296 \pm 0.001$, showing massive star clusters (the Firefly Sparkle)...The Firefly Sparkle exhibits traits of a young, gas-rich galaxy in its early formation stage. The mass of the galaxy is concentrated in 10 star clusters (49–57% of total mass), with individual masses ranging from $10^5 M_{\odot}$ to $10^6 M_{\odot}$. These unresolved clusters have high surface densities ($>10^3 M_{\odot} \text{ pc}^{-2}$), exceeding those of Milky Way globular clusters and young star clusters in nearby galaxies. The central cluster shows ... a top-heavy initial mass function. These observations provide our first spectrophotometric view of a typical galaxy in its early stages, in a 600-million-year-old Universe.

Work in Progress: Very Massive Stars, ... Black Holes, Relativistic Dynamics

Astronomy & Astrophysics manuscript no. output
March 23, 2025

Submitted A&A May 2025

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Rapid formation of a very massive star $>50000 M_{\odot}$ and subsequently an IMBH from runaway collisions

Direct N -body and Monte Carlo simulations of dense star clusters

Marcelo C. Vergara^{1*}, Abbas Askar^{2**}, Albrecht W. H. Kamlah^{3,1}, Rainer Spurzem^{1,4,5}, Francesco Flammini Dotti¹,
Dominik R.G. Schleicher⁶, Manuel Arca Sedda^{7,8,9,10}, Arkadiusz Hypki^{11,2}, Mirek Giersz², Jarrod Hurley^{12,13},
Peter Berczik^{2,14,15}, Andres Escala¹⁶, Nils Hoyer^{3,17,18}, Nadine Neumayer³, Xiaoying Pang^{19,20},
Ataru Tanikawa^{21,22}, Renyue Cen^{23,24}, and Thorsten Naab²⁵

(Affiliations can be found after the references)

Received September 15, 1996; accepted March 16, 1997

ABSTRACT

Context. Simulation of a massive young star cluster using Nbody6++GPU and MOCCA is presented. The cluster is initially more concentrated than previously published models, with one million stars, an initial mass of $5.86 \cdot 10^5 M_{\odot}$, and an initial half-mass radius of $r_{h,i} = 0.1$ pc.

Aims. Analyse the growth of a very massive star (VMS) and the characteristic properties of the collisions leading to its growth. It also seeks to understand the formation of an intermediate-mass black hole (IMBH) through direct stellar collisions in the central regions of a star cluster.

Methods. Both direct N -body and Monte Carlo simulations, include an updated of the standard single and binary stellar evolution packages (SSE and BSE) for massive and VMS, particularly regarding stellar radii, rejuvenation, and mass loss during stellar collisions, providing approximations for stellar radii, rejuvenation, and mass loss during collisions. These updates serve as reasonable extrapolations for very massive stars, though future detailed models may require further refinement.

Results. The simulations reveal that a VMS of approximately $50,000 M_{\odot}$ forms through a large number of direct stellar collisions within about 5 Myr in the central regions of the star cluster. This VMS eventually becomes an IMBH.

Conclusions. Our model shows that dense stellar environments are ideally suited to forming massive black holes through stellar collisions. The findings have implications for the formation of black hole seeds in star clusters and are relevant to recent JWST observations of young star clusters in the early Universe.

Key words. star clusters: globular clusters—stars: massive—methods: numerical

Work in Progress: Very Massive Stars, ... Black Holes, Relativistic Dynamics

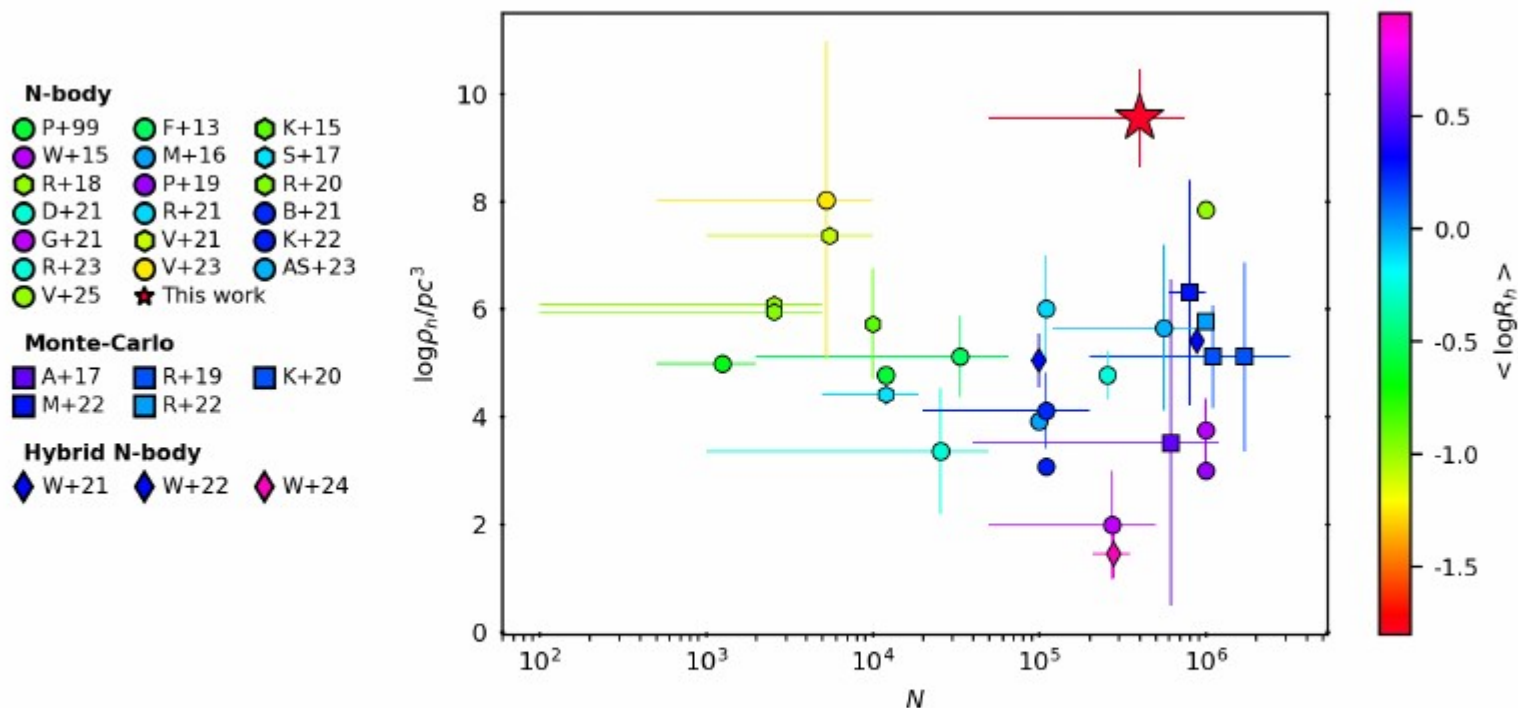
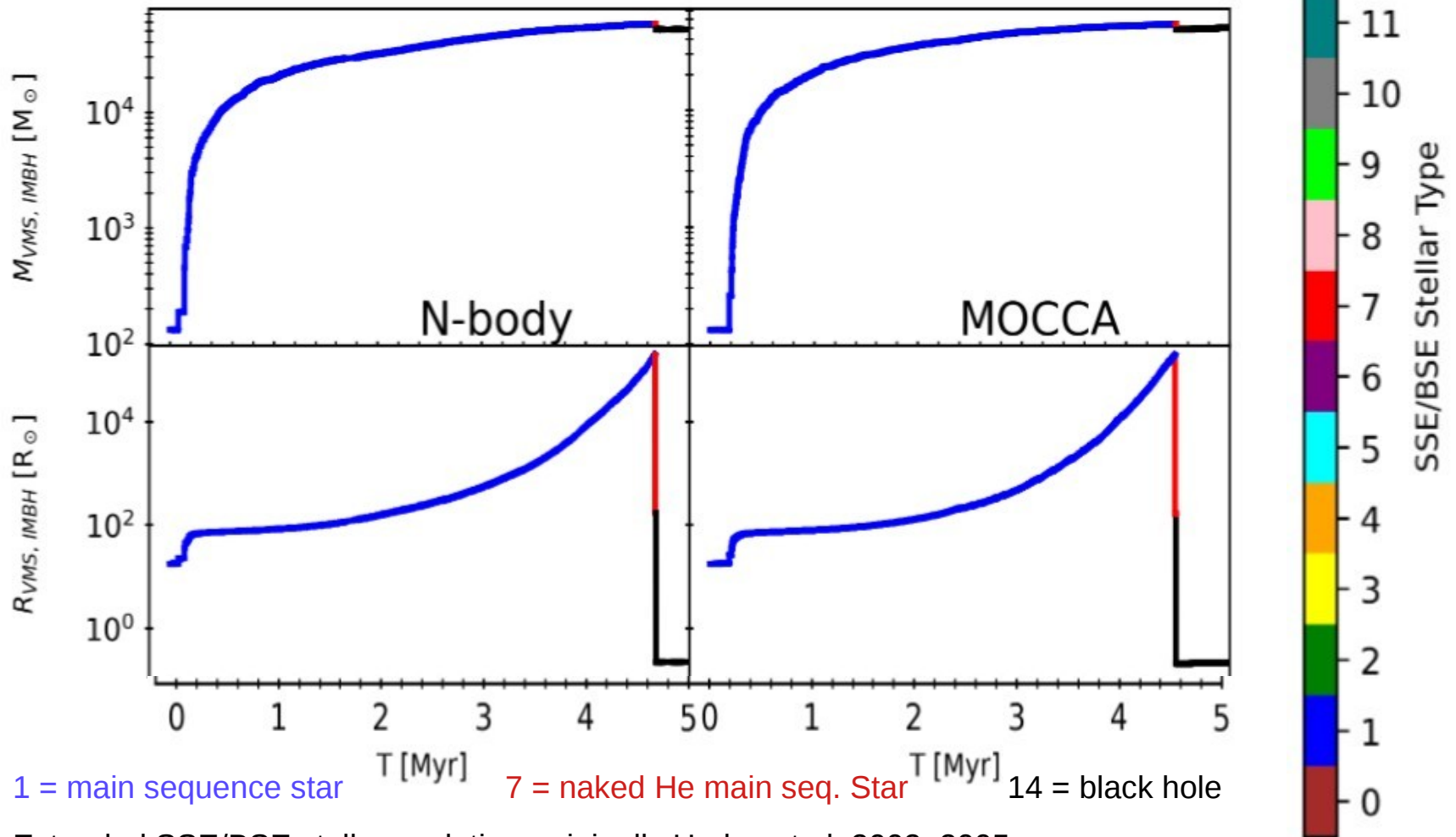


Fig. 1. Initial half-mass density ρ_h , computed at the initial half-mass radius r_h , is expressed as a function of the initial number of stars N . Figure has been reproduced from [Arca Sedda et al. \(2023, 2024a,b\)](#) and has been adapted to include a color bar with the logarithm of the average half-mass radius. Besides we include others simulations and the new ones presented in this paper.

Fast Intermediate Mass Black Hole Formation By Stellar Collisions in Dense Star Cluster

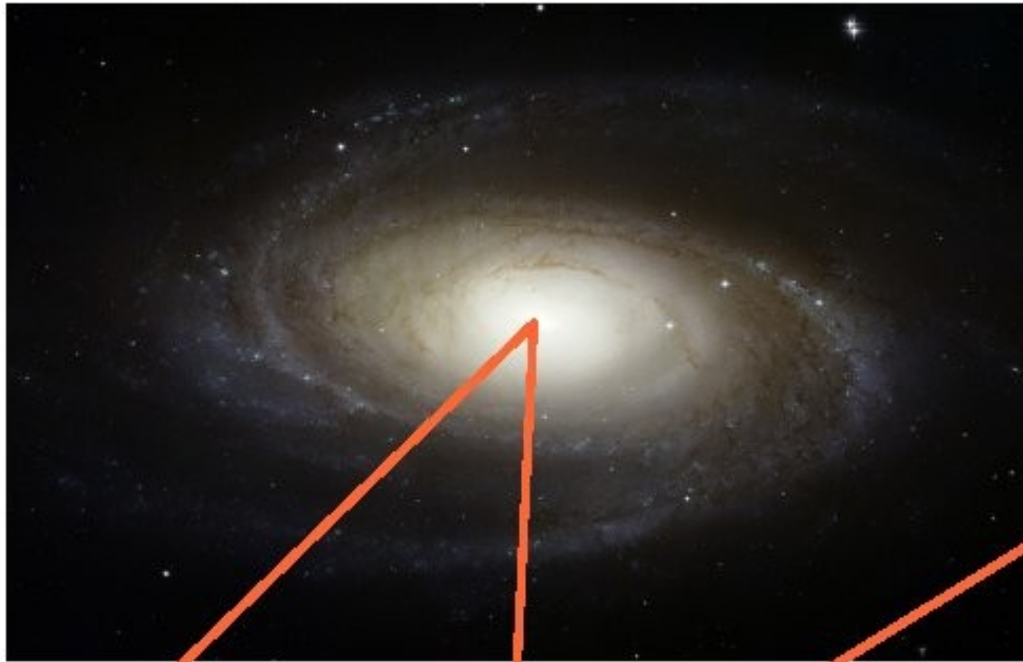
Vergara et al: Rapid formation of a VMS and subsequently an IMBH in a dense star cluster



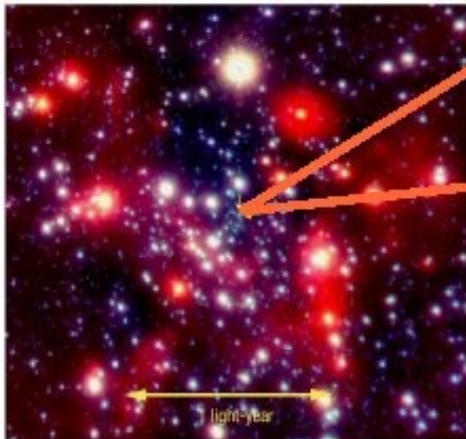
Extended SSE/BSE stellar evolution, originally Hurley et al. 2002, 2005

- 1) Star Cluster Simulations/
Black Holes/Grav. Waves
- 2) Nuclear Star Clusters**
- 3) Code(s) and Hardware
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Setting the stage: the galactic nucleus



Size ~ 10 Kpc
 Density $\sim 0.05 M_{\text{sun}} \text{ pc}^{-3}$
 Vel. Disp. ~ 40 Km/s
 Relaxation time $\sim 10^{15}$ yrs.



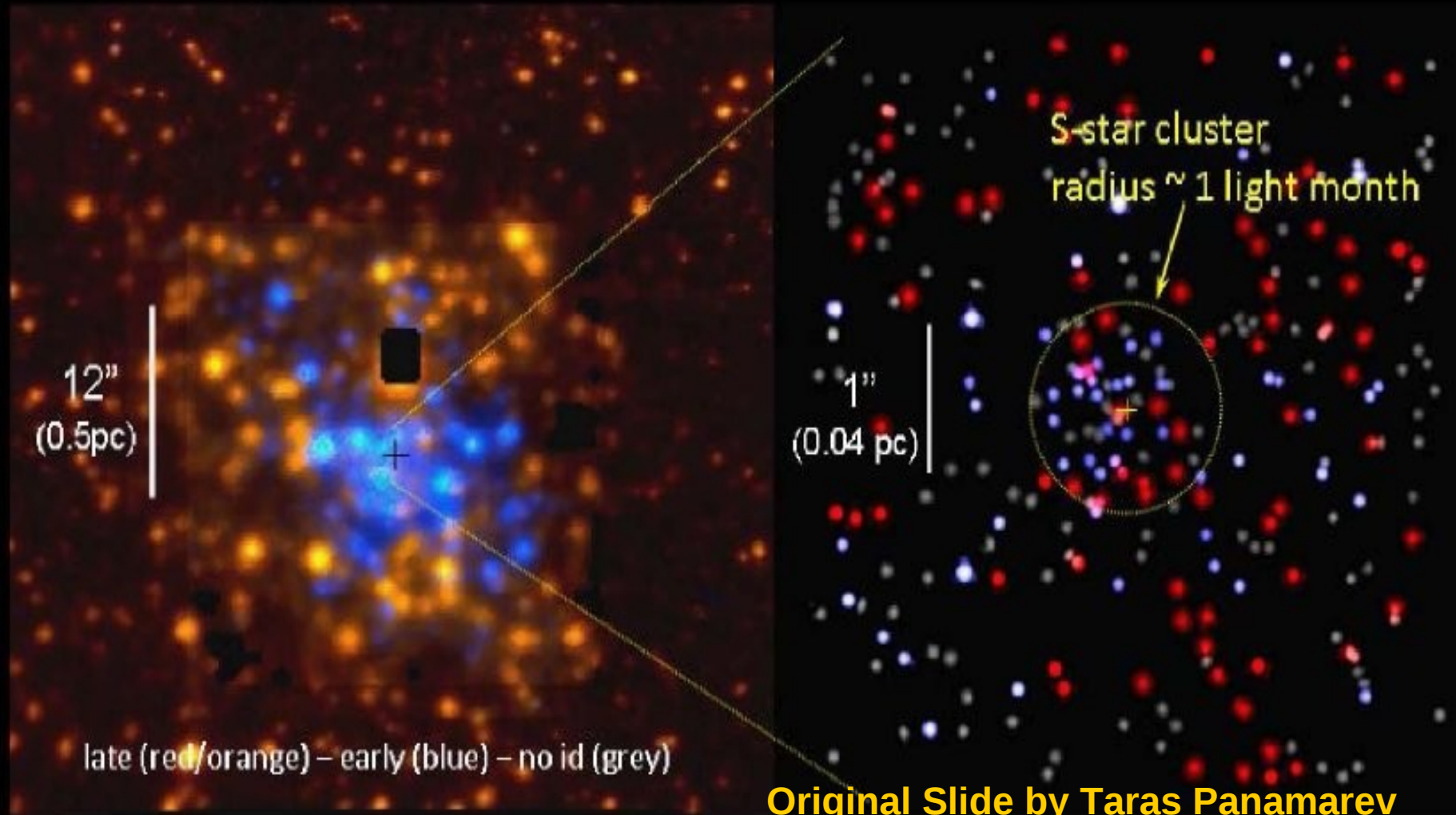
Size $\sim 1\text{-}10$ pc
 Density $\sim 10^{6-8} M_{\text{sun}} \text{ pc}^{-3}$
 Vel. Disp. $\sim 10^{2-3}$ Km/s
 Relaxation time $\sim 10^{8-9}$ yrs.

Size $\sim 10^{-7}\text{-}10^{-4}$ pc
 $R_s = 2G M_{\text{BH}} / c^2$
 $R_t \sim (\alpha M_{\text{BH}} / m_*)^{1/3} R_*$
 Loss cone aperture: θ

Slide:
Miguel
Preto

Distribution of stars

Model: Panamarev, Just, Spurzem, Berczik, Wang, Arca Sedda 2019, MNRAS
“DRAGON” of Galactic Center (one million bodies, SMBH, 1% hard binaries)



Original Slide by Taras Panamarev

Star Clusters around (supermassive) black holes

(From Frank & Rees 1976

See also:

Dokuchaev & Ozernoy 1977abc)

Predicted:
Char. Density profile
 $n(r) \sim r^{-7/4}$

r_s : Schwarzschild – rad.

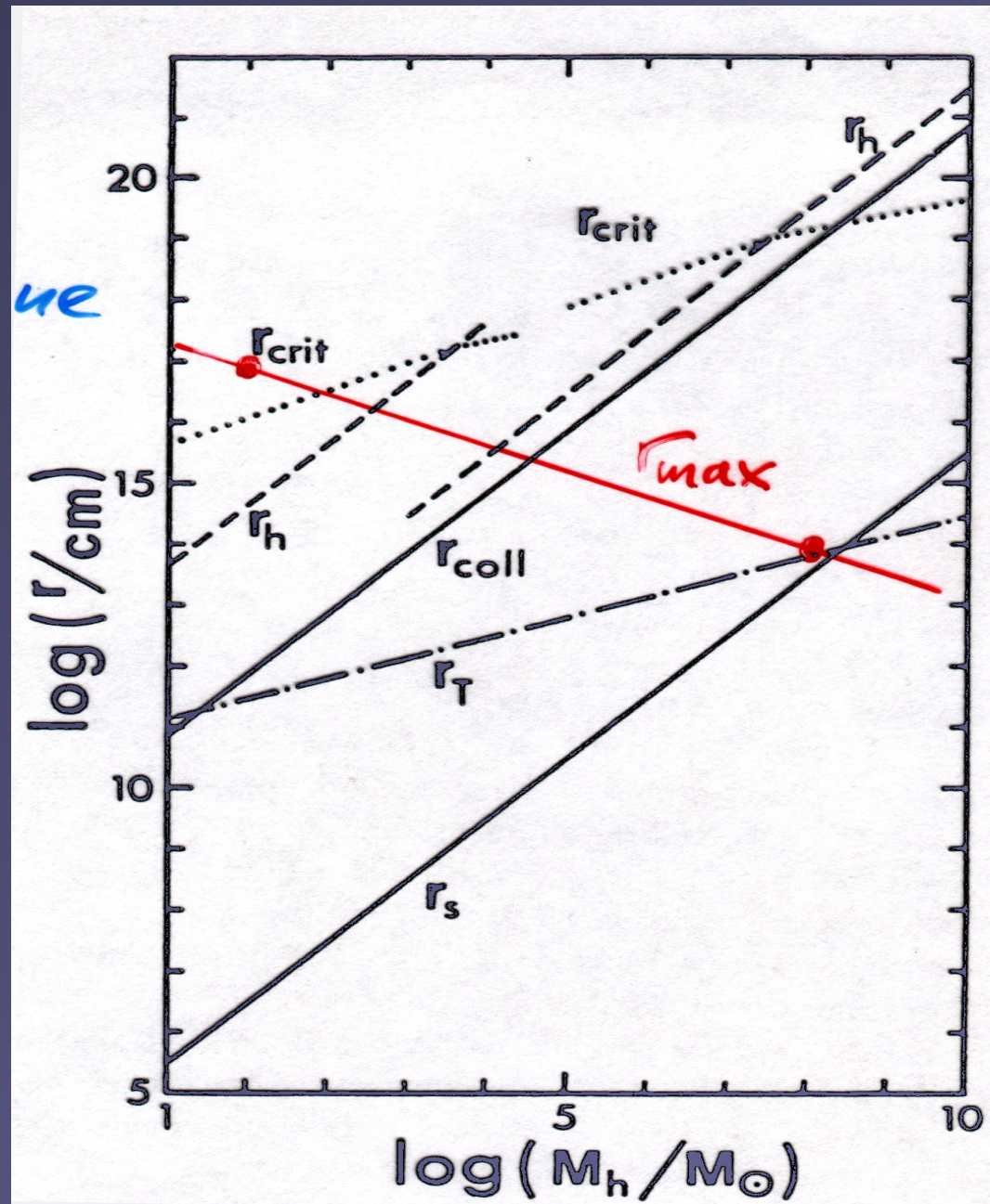
r_T : tidal disruption rad.
(solar type star)

r_{coll} : stellar collisions inside

r_h : grav. Influence rad.

r_{crit} : loss cone maximum
(spherical symmetry)

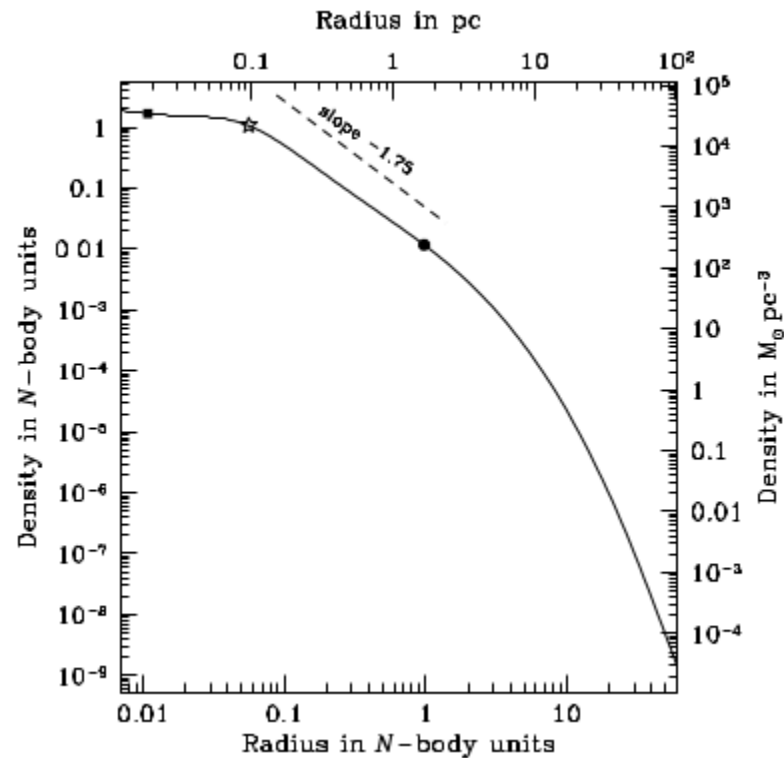
r_{max} : Oscillation in equipartition



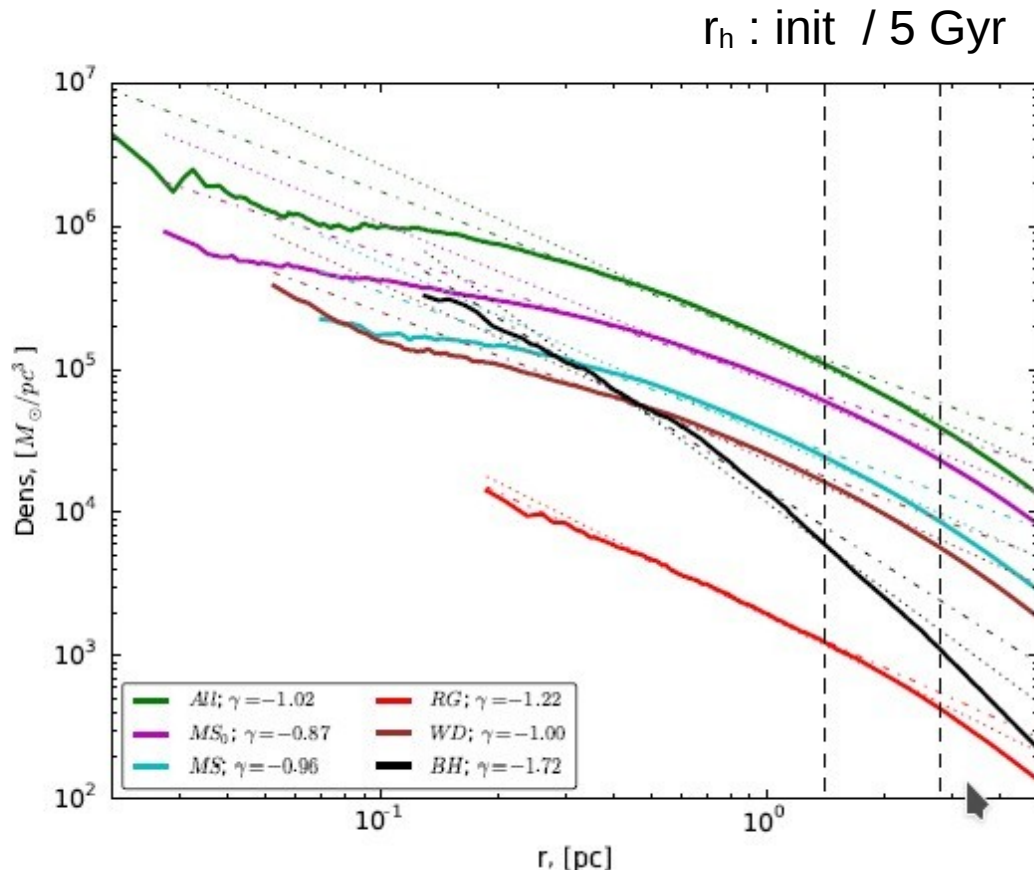
Accretion of stars on to a massive black hole: a realistic diffusion model and numerical studies

P. Amaro-Seoane,[★] M. Freitag[★] and R. Spurzem[★]

Astronomisches Rechen-Institut, Mönchhofstraße 12-14, Heidelberg, D-69120, Germany



DRAGON Galactic Center Simulations



Initial Data:

1 million stars

10% fixed SMBH mass

Zero age pop, 0.8 – 100

Spherical

Density Profiles of stars and
*-mass Black Holes in the
Galactic Center after 5 Gyr

Extremely simple

“accretion radius”

For ALL objects.

FIGURE 4.3: Stellar density profiles at $t = 5 \text{ Gyr}$ for different stellar types. Thick solid lines correspond to: All - all stars, MS_{low} - low mass main sequence stars, MS - main sequence, RG - red giants, WD - white dwarfs, BH - black holes. Corresponding power-law slopes fitted inside the initial and final influence radii of the SMBH are shown as dash-dotted and dotted lines of the same colour. The dashed vertical lines denote the initial influence radius ($r = 1.4 \text{ pc}$) and the influence radius at $t = 5 \text{ Gyr}$ ($r \sim 2.8 \text{ pc}$) of the SMBH. The power-law indices fitted inside $r = 1.4 \text{ pc}$ are shown in the legend.

Panamarev, Just, Spurzem, et al. 2019

Panamarev, ..., Just, Spurzem, 2018, MNRAS

DRAGON Galactic Center Simulations

TABLE 4.1: Properties of different stellar types at $t = 5$ Gyr.

Stellar type	N_{tot}	$\langle m \rangle$ (M_{\odot})	$\gamma(r < 1.4\text{pc})$	$\gamma(r < 2.8\text{pc})$	\dot{N}_{acc} (Gyr^{-1})	$\dot{N}_{\text{acc}}(r < 3.2\text{pc})$ (Gyr^{-1})
Low-mass main sequence	5.0×10^7	0.25	-0.87 ± 0.01	-1.11 ± 0.03	3278	-
Main sequence	4.3×10^6	0.92	-0.96 ± 0.02	-1.21 ± 0.05	658	-
Red giant	1.5×10^5	1.24	-1.22 ± 0.12	-1.34 ± 0.15	39	-
White dwarf	1.6×10^6	0.71	-1.00 ± 0.02	-1.23 ± 0.03	138	225
Black hole	1.8×10^4	10.05	-1.72 ± 0.04	-1.98 ± 0.07	2	3
All stars	5.8×10^7	0.33	-1.02 ± 0.02	-1.23 ± 0.03	4120	-

Panamarev, Just, Spurzem, et al. MNRAS, 2019

Panamarev, Berczik,..., Just, Spurzem, 2018, MNRAS

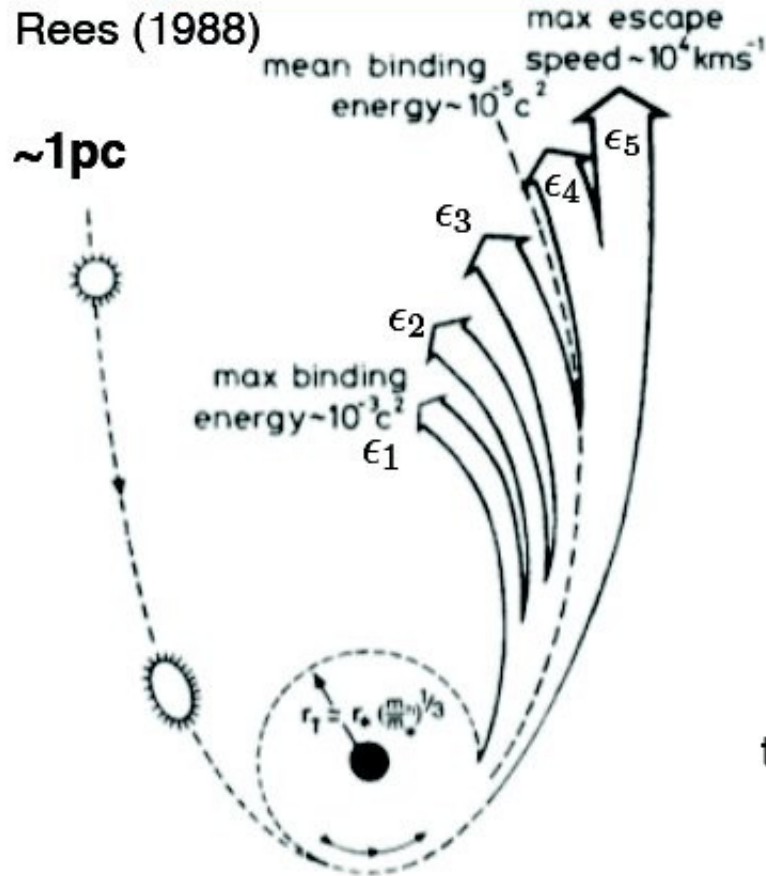
First Simulation of its kind – BUT!!

- One single accretion (“tidal”) radius for ALL objects!
- Rather outdated stellar evolution (winds/remnants/BH mergers)
- Neutron stars / pulsars lost (no gal. Potential, too high kicks)

Tidal Disruption Events for DRAGON-III (5 Improvements)

Tidal Disruption of a star by a SMBH

Standard Picture



Tidal disruption radius
(Tidal force=self-gravity force):

$$r_t = \left(\frac{M_{\text{BH}}}{m_*} \right)^{1/3} r_*$$

$\Delta\epsilon$: Spread in debris energy by tidal force

$$\Delta\epsilon = \frac{GM_{\text{BH}}}{r_t} \frac{r_*}{r_t}$$

ϵ : Debris specific energy

if $\epsilon \geq 0$ Stellar debris flies away from the black hole

if $\epsilon < 0$ Stellar debris is bounded by the black hole's gravity and falls back to black hole

t : Fallback time for most tightly bound debris

$$t_{\text{fall}} \sim 0.1 \text{ yr} \left(\frac{r_*}{R_{\odot}} \right)^{3/2} \left(\frac{m_*}{M_{\odot}} \right)^{-1} \left(\frac{M_{\text{BH}}}{10^6 M_{\odot}} \right)^{1/2}$$

What is the rate of mass fallback?

Mass fallback rate I.

$$\frac{dM}{dt} = \left| \frac{dM(\epsilon)}{d\epsilon} \right| \left| \frac{d\epsilon}{dt} \right| \quad (\epsilon < 0)$$

Specific energy:

$$\epsilon \approx -\frac{GM_{\text{BH}}}{2a}$$

Its time derivative:

$$\frac{d\epsilon}{dt} = -\frac{1}{3}(2\pi GM_{\text{BH}})^{2/3} t^{-5/3}$$

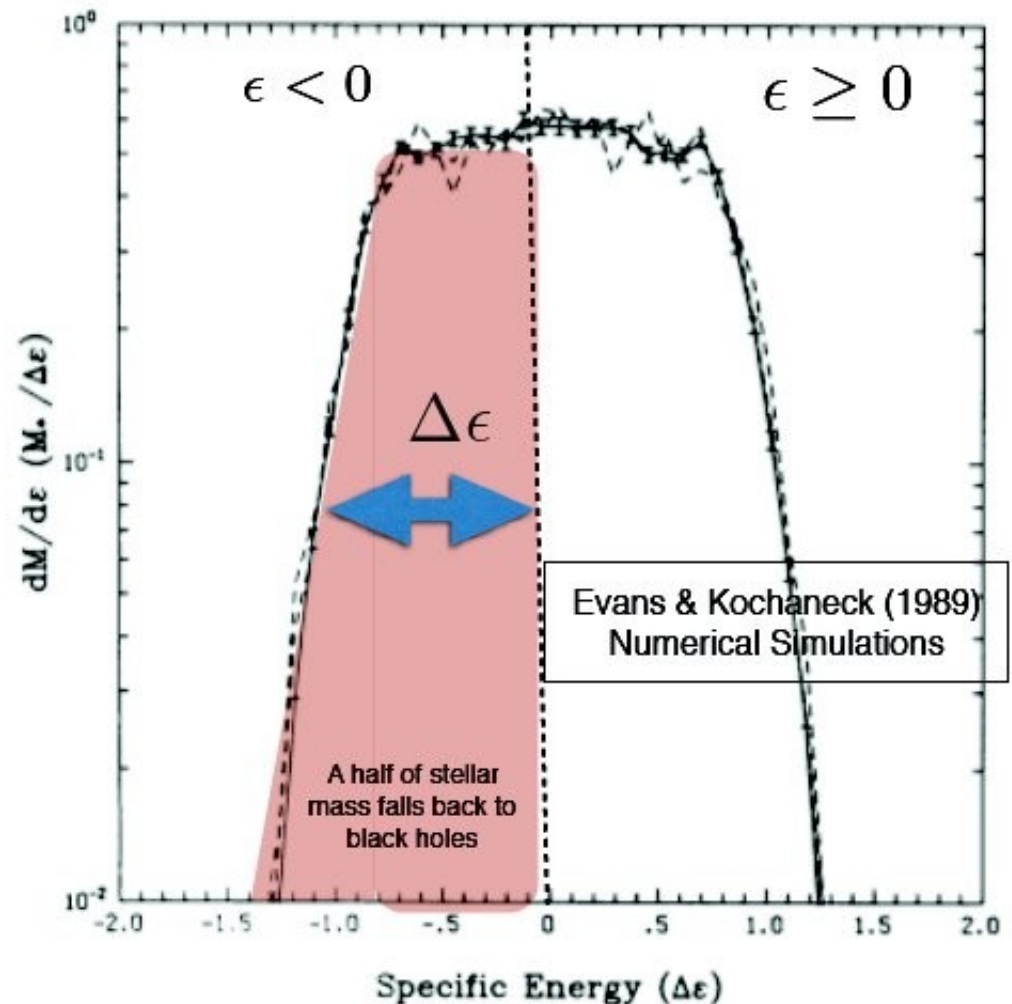
(by using Keplerian third law)



$$\frac{dM}{dt} \propto t^{-5/3}$$

Rees's conjecture (1988)

Differential mass-energy distribution of stellar debris



DRAGON III of Galactic Center

tidal disruption and partial accretion – improvement 1

(Cho Master Thesis 2024, in preparation for Cho et al. 2025, A&A;

see also Zhong, Li, Berczik, Sp. 2022, Hayasaki, Zhong, Hayasaki, Li, Berczik, Sp. 2023)

Classification

For eccentric orbits, if $\Delta\epsilon + \epsilon_{orb} \leq 0$, all the stellar debris are bound by the black hole after TDE.

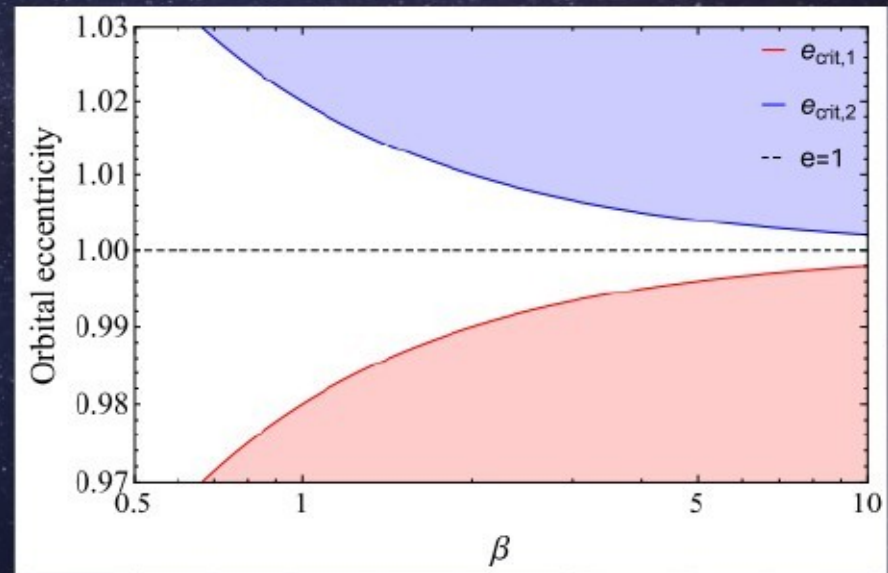
$$e_{crit,1} = 1 - \frac{2}{\beta} \left(\frac{m_{\star}}{M_{BH}} \right)^{1/3}$$

For a hyperbolic orbit, if $-\Delta\epsilon + \epsilon_{orb} \leq 0$,

$$e_{crit,2} = 1 + \frac{2}{\beta} \left(\frac{m_{\star}}{M_{BH}} \right)^{1/3},$$

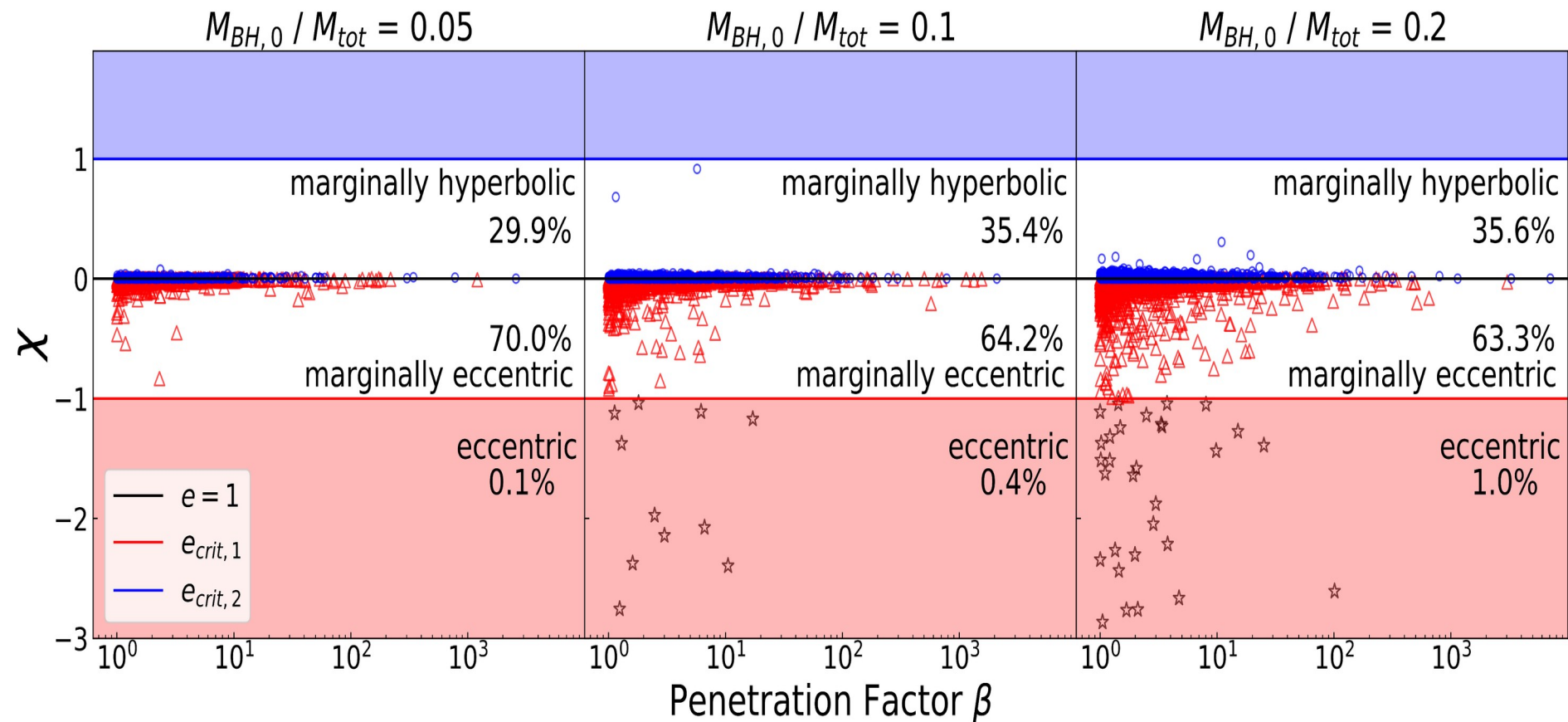
above which all the stellar debris should escape from BH.

(Hayasaki et al, 2018)



Hayasaki et al, 2018

Results – TDE (Slide by Philip Cho, Pilot Simulations)



$$\beta = r_t / r_p \quad ; \quad \chi = (e-1)/\Delta e \quad ; \quad \Delta e = (2/\beta)(m_*/M_{BH})^{(1/3)}$$

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tidal disruption and partial accretion – Improvement 1 (Cho)

Mass Fraction of a tidally disrupted star accreted to SMBH

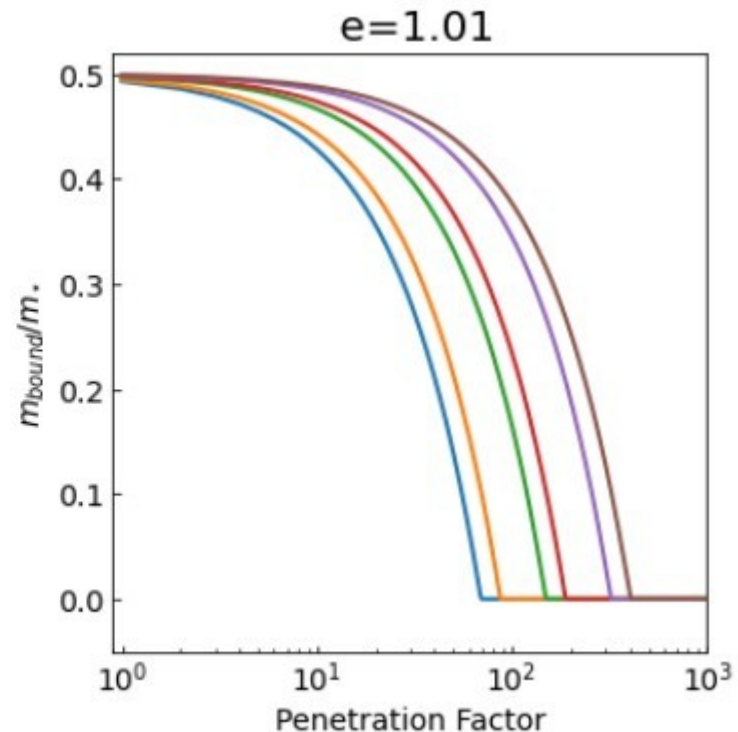
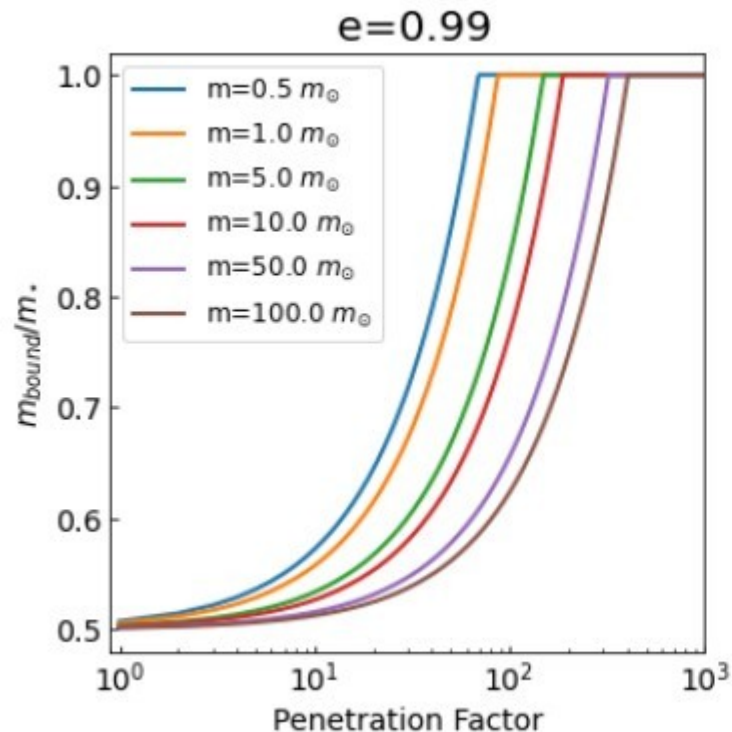
Using first order analytic model of energy profile in star (top-hat)

(see Hayasaki et al. 2013, 2018)

Eccentric TDE

Hyperbolic

TDE



DRAGON III of Galactic Center

tidal disruption and partial accretion – improvement 3 (Cho)

Preliminary Results: Percentage of Ecc./Hyp. Partial TDEs

(Relevant for TDE light curves)

Model	CMBH	r_{acc}	N_{TDE}	M_{acc}	M_{esc}	f_e	f_{me}	f_p	f_{mh}	f_h
1	3792 M_{\odot}	10^{-5}	245	235 M_{\odot}	107 M_{\odot}	2.1	86.5	0	11.4	0
2		10^{-4}	1210	1051 M_{\odot}	494 M_{\odot}	2.7	67.7	0	29.6	0
3		10^{-3}	3742	2381 M_{\odot}	1382 M_{\odot}	2.1	52.3	0	45.6	0
4	7583 M_{\odot}	10^{-5}	597	346 M_{\odot}	250 M_{\odot}	0.5	88.1	0	11.4	0
5		10^{-4}	2842	1784 M_{\odot}	1191 M_{\odot}	1.3	63.4	0	35.3	0
6		10^{-3}	6963	2883 M_{\odot}	2138 M_{\odot}	1.96	50.6	0	47.4	0.04
7	15168 M_{\odot}	10^{-5}	887	306 M_{\odot}	273 M_{\odot}	0.6	82.2	0	17.2	0
8		10^{-4}	4114	1584 M_{\odot}	1307 M_{\odot}	1.2	63.2	0	35.6	0
9		10^{-3}	9395	3248 M_{\odot}	2568 M_{\odot}	3.05	51.7	0	45.2	0.05

TABLE 3.1: Simulation Results. The model number is shown in the first column. The second and third column are the initial mass of CMBH, in solar mass unit, and the accretion radius in NB unit, respectively. The fourth column shows the total number of TDEs occurred. The fifth column shows the total accreted mass by the black hole, and the sixth column shows the total debris mass that escaped from the CMBH potential.

From seventh to last columns, the percentage of each type of TDEs are shown.

Taken from
Philip Seunghee Cho:
Master Thesis
Univ. Heidelberg
2023

DRAGON III of Galactic Center

Why is it important to know about TDE event rates from simulations?

Hundreds if not tens of thousands will be detected soon

(LSST, ZTF, ULTRASAT, eROSITA, ..., Szekerzczes, Ryu et al. 2024, A&A)

We need to know about parameters of host galaxies, nuclear clusters

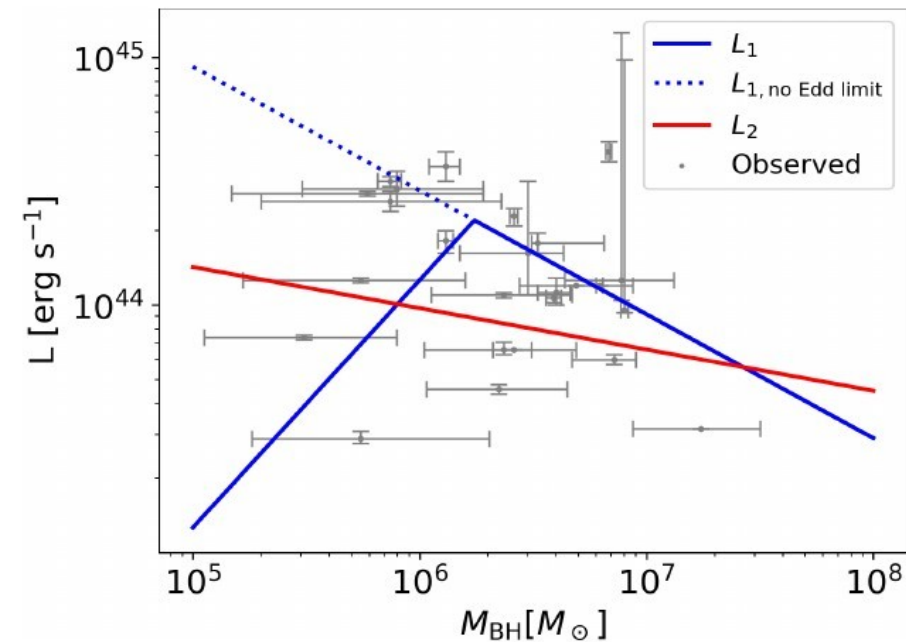


Fig. 2. Two luminosities of TDEs, L_1 (Eq. (6), Thomsen et al. 2022) and L_2 (Eq. (7), Ryu et al. 2020b) as a function of M_{BH} overlaid with observed optical TDEs using gray dots and error bars (see Table A.1 in Wong et al. 2022 and Table 1 in Ryu et al. 2020b, and references therein). The dotted blue line conveys L_1 as given by Eq. (6), whereas the solid blue line show the Eddington limited L_1 .

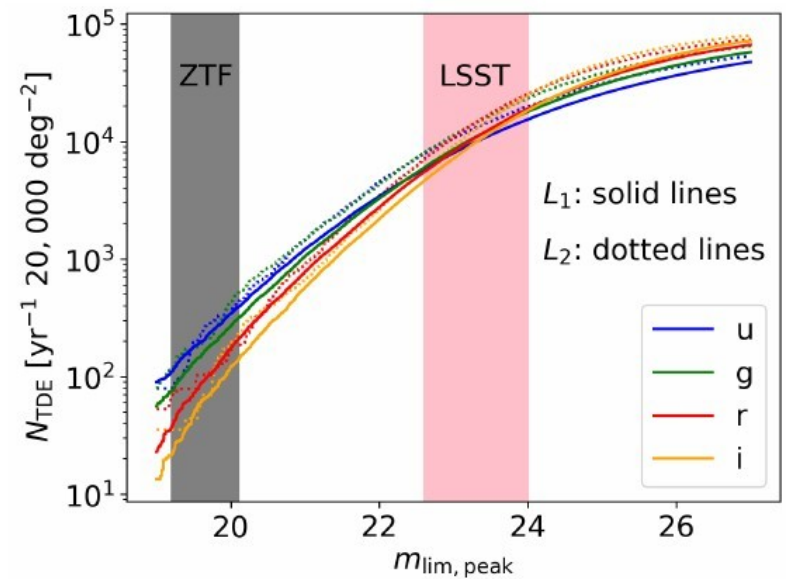


Fig. 6. Annual unlensed TDE detection rates as a function of the peak brightness, $m_{\text{lim,peak}}$, at $T = 2 \times 10^4$ K for L_1 in solid lines and for L_2 shown in dotted lines. The limiting magnitudes for the different filters of ZTF and LSST surveys are within the gray and pink shaded areas, respectively. These limiting magnitudes are 0.7 brighter than the survey limiting magnitude $m_{\text{lim,survey}}$ (i.e., $m_{\text{lim,peak}} = m_{\text{lim,survey}} - 0.7$), denoted by LSST1, in order for TDEs to have multiple detections near their peak brightness above the limit $m_{\text{lim,survey}}$.

DRAGON III of Galactic Center

accretion and disruption processes depend on type of object – improvement 2
(Minzburg Master Thesis to be published)

Improvement 2: Tidal Disruption Radius varies with * parameters (m,r);

Include relativistic inspirals of bh and ns (not plotted here)

Include direct capture of very low mass ms stars (not plotted here)

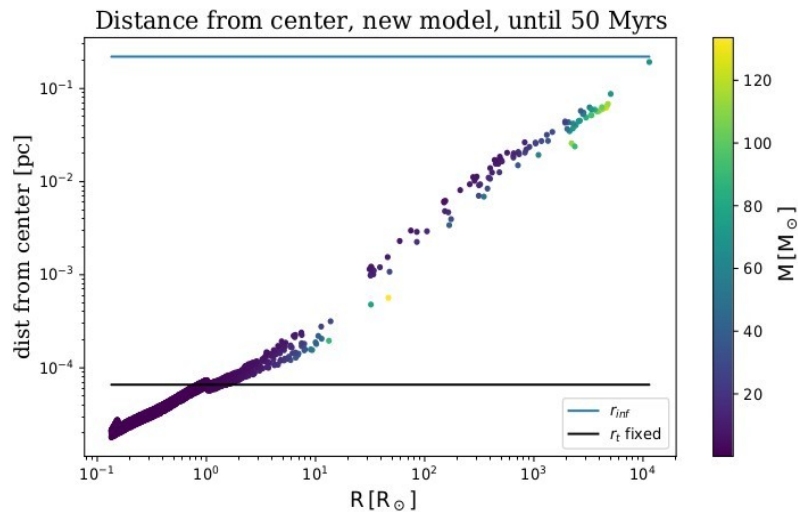


Figure 3.13: Distances of accreted objects vs their size in R_{\odot} colorcoded in mass until 50 Myrs. Accreted stars have sizes between 0.1 and 10000 R_{\odot} , most massive star has $\sim 130M_{\odot}$. The distances span 4 orders of magnitude between 10^{-5} and 10^{-1} pc.

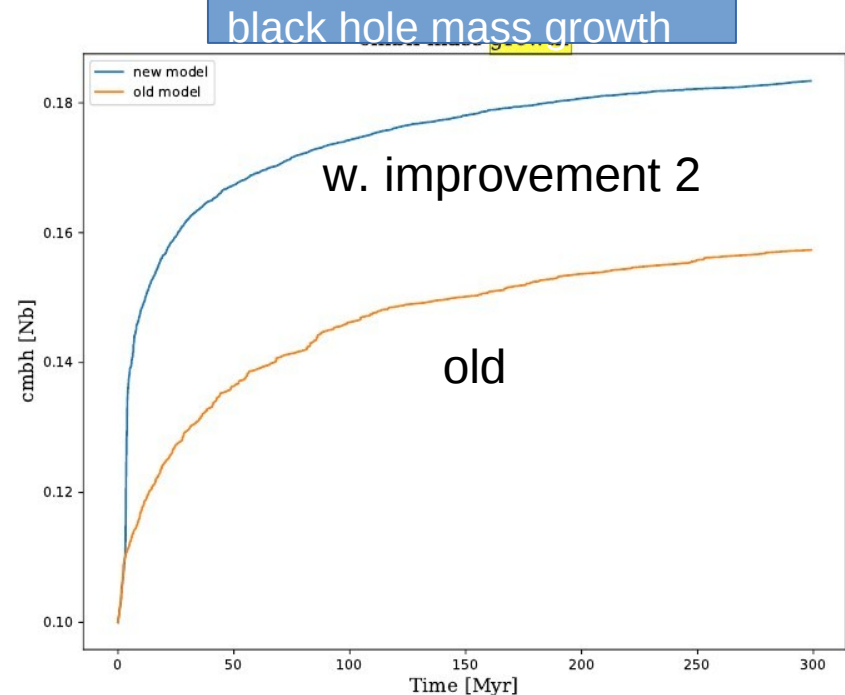


Figure 3.15: Mass growth of the central massive black hole throughout the simulation for the old model (orange) and new model (blue). While a steeper rise at the beginning of the simulation can be observed for the new model, the slopes appear similar for both models after 100 Myrs.

Improvement 3:

Partial and Full Tidal Disruption Events in Nuclear Star Cluster Simulations

From: Zhong, Li, Berczik, Spurzem, 2022, ApJ

THE ASTROPHYSICAL JOURNAL, 933:96 (13pp), 2022 July 1

Zhong et al.

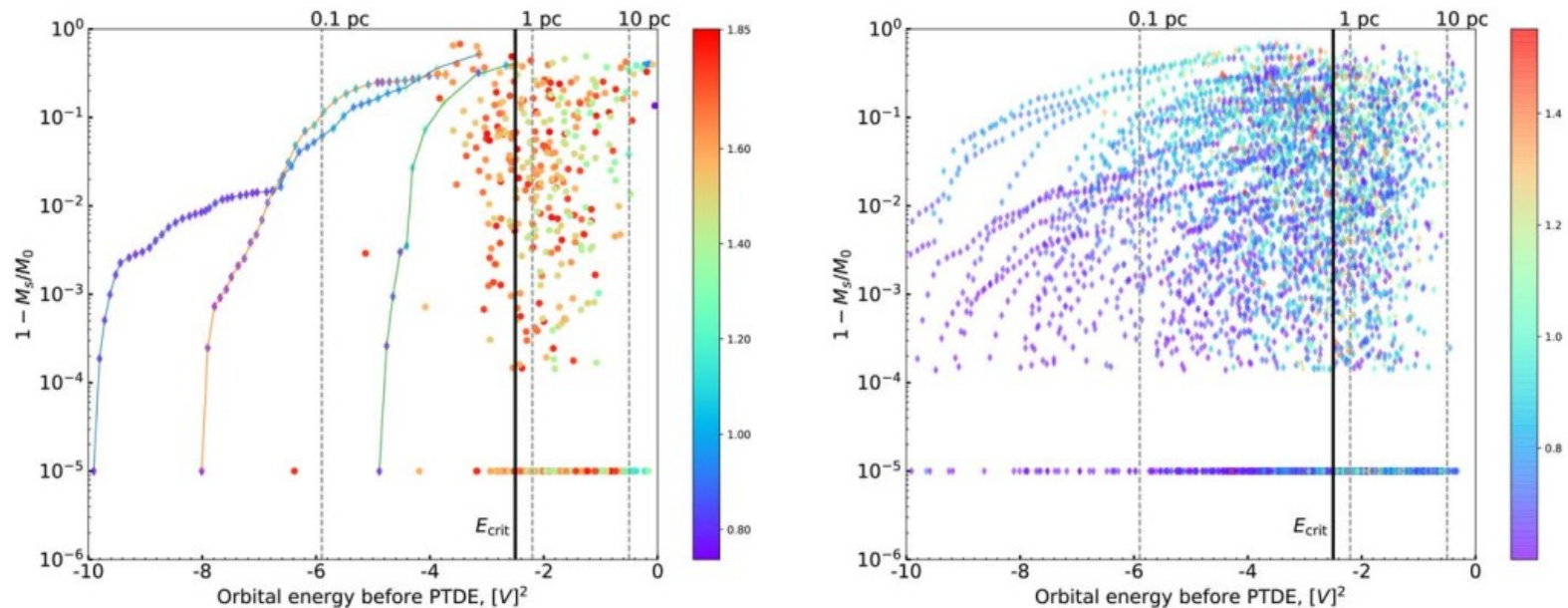
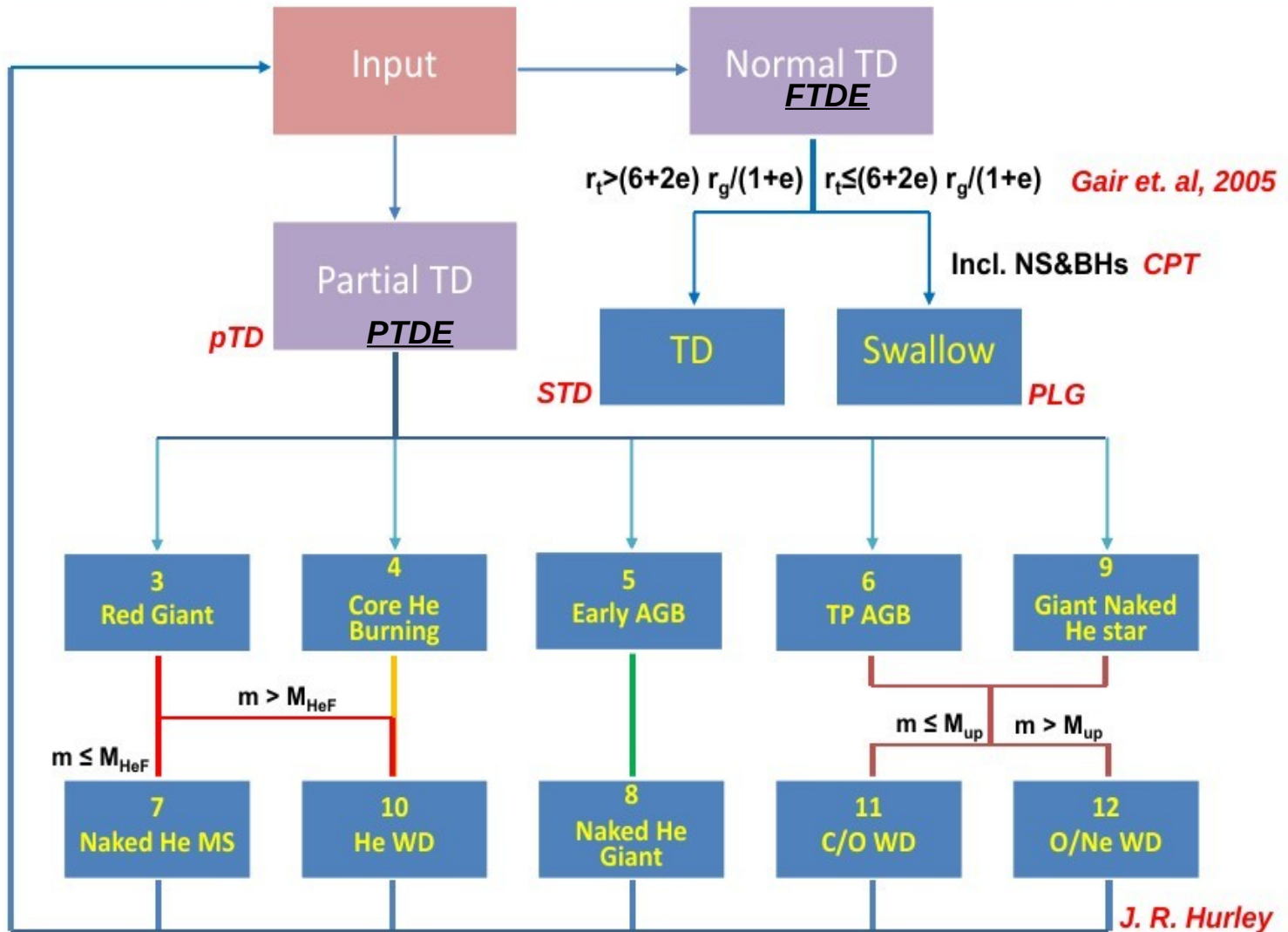


Figure 4. Distribution of stars that produce PTDEs in the parameter space spanned by the mass (m_s) and specific orbital energy (E_{tot}) of the stars. For every PTDE, the values of m_s and E_{tot} are measured at the moment immediately before the onset of the event. Colors indicate the penetration factor β . The vertical solid line marks the position of the critical energy, and the corresponding physical radius is roughly 0.7 pc. For clarity, the vertical axis shows the quantity $1 - m_s/m_0$, where $m_0 = 1/N[M]$ is the initial mass of the star, and is set to log scale. However, with this configuration, the PTDEs produced by the normal stars are not visible in the figure because $1 - m_s/m_0 = 0$. We have artificially reset the value of $1 - m_s/m_0$ to 10^{-5} for the PTDEs produced by normal stars. The left panel shows the PTDEs after which the leftover stars are ejected, while the right panel shows the PTDEs after which the leftover stars are retained in the star cluster. In the left panel, we also plot the historical PTDEs (diamonds connected by lines) for three individual stars. This figure is generated from the data of one realization of the fiducial model.

More PTDE than FTDE – orbits diverse!



Tidal Disruption Events in Merging Nuclear star clusters

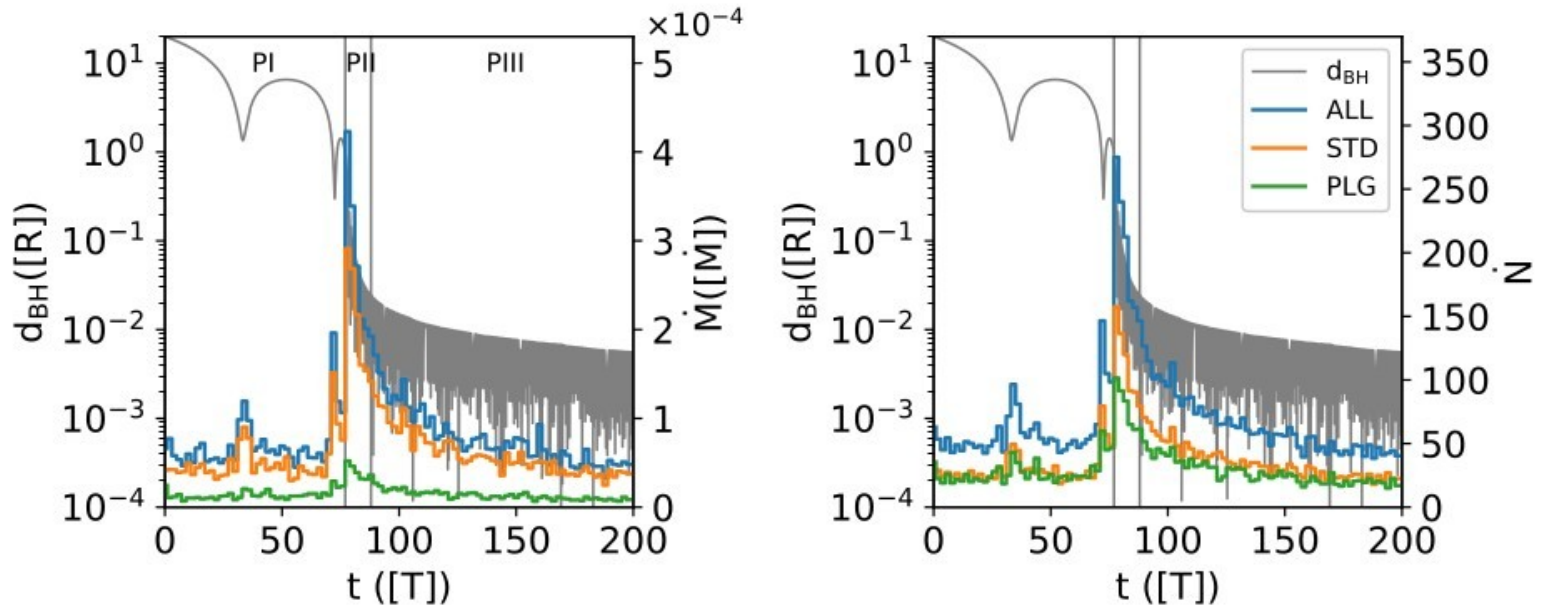


Figure 6. The tidal disruption/swallow evolution of MS stars. The left panel and right panel are, respectively, the evolution of the mass accretion rate and the event rate. The blue, orange, and green solid lines represent the results of all of MS stars, standard TDEs, and MS stars which directly plunge into the SMBH, respectively. The two vertical gray lines divide the evolution into three phases.

Li, Zhong, Berczik, Spurzem, Chen, Liu, MNRAS 2023

Formation and Evolution of Compact Binaries Containing Intermediate Mass Black Holes in Dense Star Clusters

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⁸Astronomy and Space Science Institute, Daejeon 34055, Republic of Korea

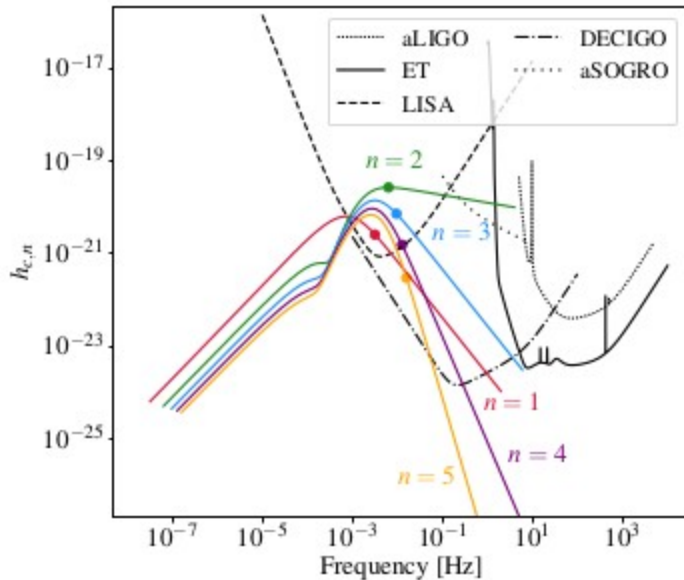


Figure 2. A sample illustration of the characteristic amplitude of the n -th harmonics, $h_{c,n}$, from $n = 1$ to 5 emitted during the **IMR**. The dashed black V-shaped curve represents LISA's sensitivity curve ($\sqrt{f}S_h(f)$; see Section 2.2), while two U-shaped curves indicate the sensitivities of aLIGO (densely-dotted line) and ET (solid line).

TDE in Nbody6++GPU - Improvement 5

<https://arxiv.org/abs/2503.22109>

- 1) Star Cluster Simulations/
Black Holes/Grav. Waves
- 2) Nuclear Star Clusters
- 3) Code(s) and Hardware
- 4) Summary and References

Code(s)
Nbody6++GPU

Physical and Numerical Methods: Direct Simulations

Direct: high accuracy / active-inactive particles

The Hermite Scheme: 4th Order on two time points

$$\vec{a}_0 = \sum_j Gm_j \frac{\vec{R}_j}{R_j^3} \quad ; \quad \vec{\ddot{a}}_0 = \sum_j Gm_j \left[\frac{\vec{V}_j}{R_j^3} - \frac{3(\vec{V}_j \cdot \vec{R}_j) \vec{R}_j}{R_j^5} \right] ,$$

$$\begin{aligned} \vec{x}_p(t) &= \frac{1}{6}(t - t_0)^3 \vec{\ddot{a}}_0 + \frac{1}{2}(t - t_0)^2 \vec{a}_0 + (t - t_0) \vec{v} + \vec{x} , \\ \vec{v}_p(t) &= \frac{1}{2}(t - t_0)^2 \vec{\ddot{a}}_0 + (t - t_0) \vec{a}_0 + \vec{v} , \end{aligned}$$

Repeat Step 1 at t_1 using predicted $\mathbf{x}, \mathbf{v} \rightarrow \mathbf{a}_1, \mathbf{\ddot{a}}_1$

NBODY6++GPU: <https://github.com/nbody6ppgpu/>

Physical and Numerical Methods: Direct Simulations

$$\frac{1}{2}\vec{a}^{(2)} = -3\frac{\vec{a}_0 - \vec{a}_1}{(t - t_0)^2} - \frac{2\vec{a}_0 + \vec{a}_1}{(t - t_0)}$$
$$\frac{1}{6}\vec{a}^{(3)} = 2\frac{\vec{a}_0 - \vec{a}_1}{(t - t_0)^3} - \frac{\vec{a}_0 + \vec{a}_1}{(t - t_0)^2} ,$$

The Hermite Step
Get Higher Derivatives

$$\vec{x}(t) = \vec{x}_p(t) + \frac{1}{24}(t - t_0)^4\vec{a}_0^{(2)} + \frac{1}{120}(t - t_0)^5\vec{a}_0^{(3)} ,$$
$$\vec{v}(t) = \vec{v}_p(t) + \frac{1}{6}(t - t_0)^3\vec{a}_0^{(2)} + \frac{1}{24}(t - t_0)^4\vec{a}_0^{(3)} .$$

The Corrector Step – this is not time symmetric!

Physical and Numerical Methods: Direct Simulations

P(EC)ⁿ Scheme, n=1, Kokubo et al. 1998, Hut et al. 1995

$$\mathbf{x}_{p,j} = \mathbf{x}_j + \mathbf{v}_j(t - t_j) + \frac{\mathbf{a}_j}{2}(t - t_j)^2 + \frac{\dot{\mathbf{a}}_j}{6}(t - t_j)^3,$$

1 higher order in
x-prediction

$$\mathbf{v}_{p,j} = \mathbf{v}_j + \mathbf{a}_j(t - t_j) + \frac{\dot{\mathbf{a}}_j}{2}(t - t_j)^2,$$

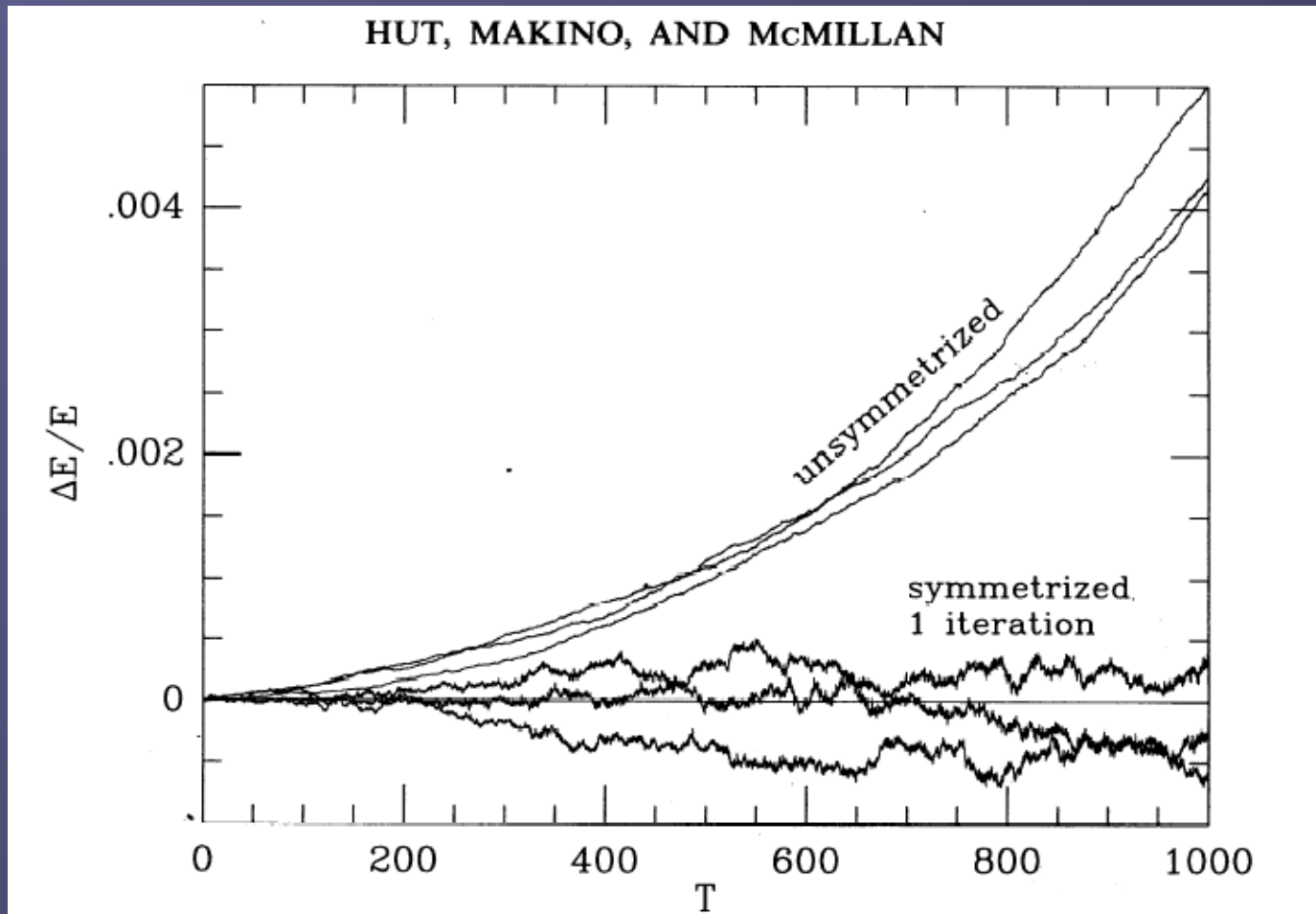
$$\mathbf{x}_1 = \mathbf{x}_0 + \frac{1}{2}(\mathbf{v}_1 + \mathbf{v}_0)\Delta t - \frac{1}{10}(\mathbf{a}_1 - \mathbf{a}_0)\Delta t^2 + \frac{1}{120}(\dot{\mathbf{a}}_1 + \dot{\mathbf{a}}_0)\Delta t^3$$

$$\mathbf{v}_1 = \mathbf{v}_0 + \frac{1}{2}(\mathbf{a}_1 + \mathbf{a}_0)\Delta t - \frac{1}{12}(\dot{\mathbf{a}}_1 - \dot{\mathbf{a}}_0)\Delta t^2.$$

This is time-symmetric (exchange 0,1, change sign(v),sign($\dot{\mathbf{a}}$))!
(But pred/corr diff. in x,v of field particles neglected)

By iteration n>1 it improves further.

Physical and Numerical Methods: Direct Simulations



1995

$N=100$

1000
orb.times

Physical and Numerical Methods: Direct Simulations

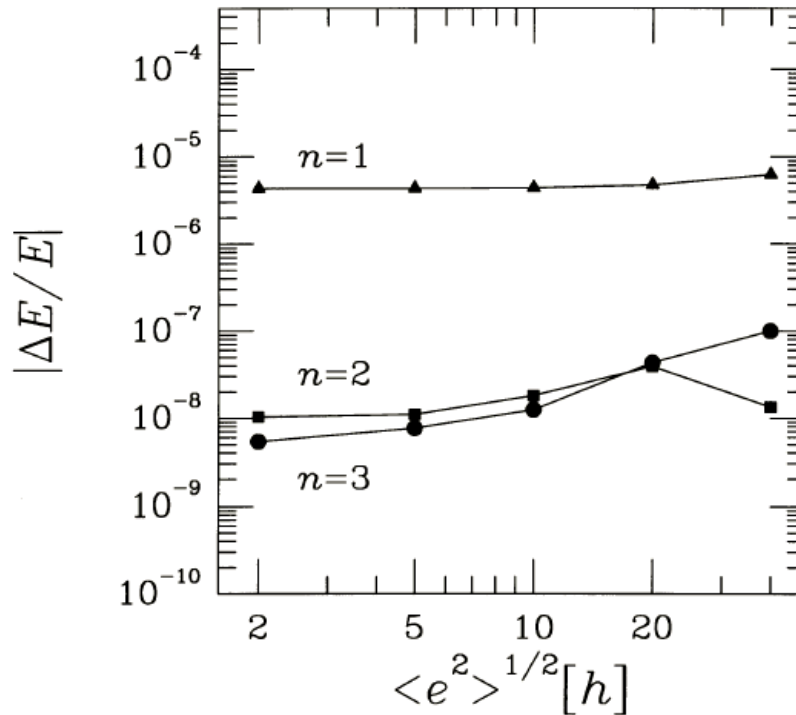


Figure 4. The relative error of the total energy of a planetesimal system that consists of 100 equal-mass ($m = 10^{25}$ g) bodies after $1000T_K$ as a function of the initial $\langle e^2 \rangle^{1/2}$ for the P(EC) n ($n = 1, 2, 3$) Hermite schemes with the hierarchical time-step scheme. The triangles show the result of the $n = 1$ scheme, the squares $n = 2$ and the circles $n = 3$.

Kokubo et al. 1998

Yoshinaga et al. 1999

Kokubo & Makino 2004

Makino, Hut et al. 2006

Glaschke, Amaro-Seoane, Sp, 2014a

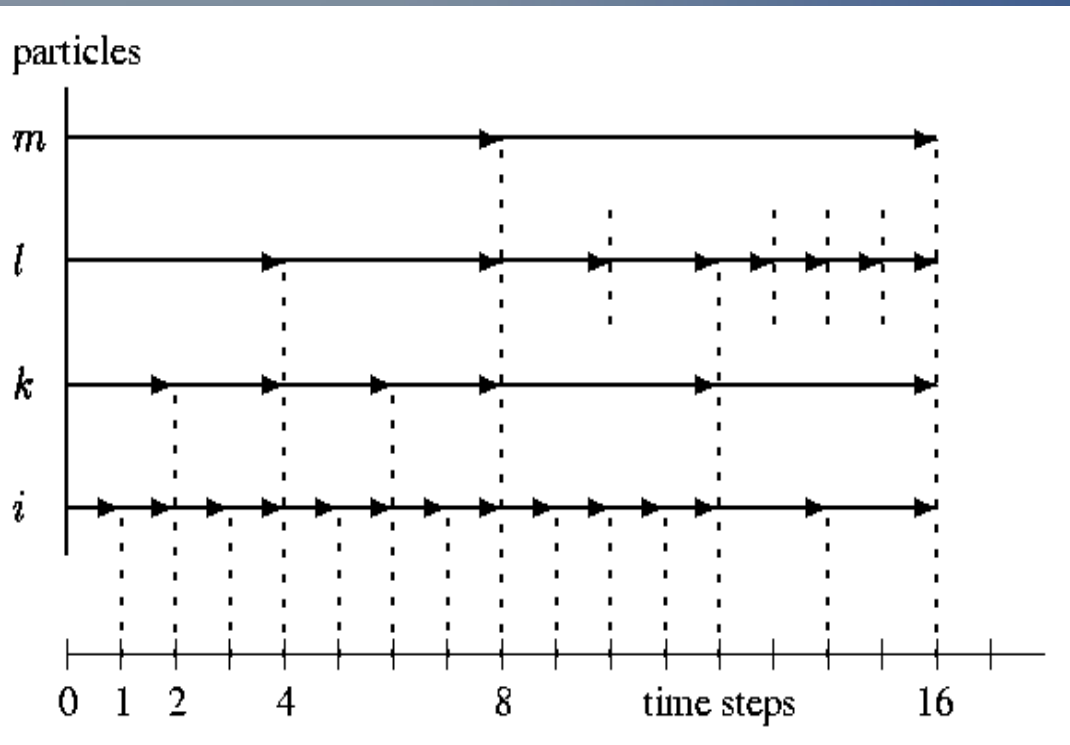
Amaro-Seoane, Glaschke, Sp, 2014b

Software

S.J.Aarseth, S. Mikkola

(ca. 20.000 lines):

- Hierarchical Block Time Steps
- Ahmad-Cohen Scheme
- Regularisations
- 4th order Hermite scheme
- NBODY6 (Aarseth 1999)
- NBODY6++ (Spurzem 1999)
- MPI
- NBODY6++GPU (Wang, Spurzem, Aarseth et al. 2015)



Hierarchical Block Time Steps

$$\Delta t = \sqrt{\eta \frac{|\vec{a}| |\vec{a}^{(2)}| + |\vec{a}|^2}{|\vec{a}| |\vec{a}^{(3)}| + |\vec{a}^{(2)}|^2}}.$$

NBODY1 – NBODY7: “The Growth of an Industry” (Aarseth 1999)

	ITS	ACS	KS	HITS	PN	AR	CC	MPI	GPU
NBODY1	✓								
NBODY2		✓		✓					
NBODY3	✓		✓						
NBODY4			✓	✓					
NBODY5	✓	✓	✓						
NBODY6		✓	✓	✓					
NBODY6GPU		✓	✓	✓				✓	
NBODY6++		✓	✓	✓			✓		
NBODY6++GPU		✓	✓	✓	✓		✓	✓	✓
NBODY7		✓	✓	✓	✓	✓			✓

ITS: Individual time-steps [107]

ACS: Ahmad-Cohen neighbour scheme [109]

KS: KS-regularization of few-body subsystems [104]

HITS: Hermite scheme integration method combined with hierarchical block time-steps [111]

PN: Post-Newtonian [150,125,151]

AR: Algorithmic regularization [125]

CC: Classical chain regularization [114]

MPI: Message Passing Interface, multi-node multi-CPU parallelization [139]

GPU: use of GPU acceleration [138] (if also MPI: multi-node many GPU [144])

Berczik, Spurzem, et al., LNCS 2013; Table from: Spurzem, Kamlah 2023, LRCA

NBODY6++GPU: <https://github.com/nbody6ppgpu/>

Part of <https://www.punch4nfdi.de/> PUNCH4NFDI Consortium w. Jülich

Table 5.1: Table showing important algorithmic, hardware and software development stepping stones in the development of direct N -body codes. The table is adapted from Aarseth (1999a), corrected in some places, but expanded to more recent developments.

Year	Keyword	Reference
1961	Force polynomial	(Aarseth, 1963)
	Individual time steps	(Aarseth, 1963)
	Gravitational softening	(Aarseth, 1963)
1966	Spherical harmonics	(Aarseth, 1967)
1969	Two-body regularization	(Kustaanheimo & Stiefel, 1965)
1972	Three-body regularization	(Aarseth & Zare, 1974)
1973	Global regularization	(Heggie, 1974)
	Neighbor scheme	(Ahmad & Cohen, 1973)
1978	Co-moving coordinates	(Aarseth, 1979)
1979	Regularized AC	(Aarseth, 1985)
1980	Planetary formation	(Lecar & Aarseth, 1986)
1986	Hierarchical block-time steps	(McMillan, 1986)
1989	Chain regularization	(Mikkola & Aarseth, 1990)
1990	Particle in box scheme	(Aarseth, Lin, & Palmer, 1993)
1991	Collisional tree code	(McMillan & Aarseth, 1993)
1992	Chain N -body interface	(Aarseth, 1994)
1993	Hermite integration	(Makino, 1991)
1995	Synthetic stellar evolution	(Tout et al., 1997)
	Tidal circularization	(Mardling, 1995a, 1995b)
	Slow chain regularization	(Mikkola & Aarseth, 1998)
1996	Hierarchical stability	(Mardling & Aarseth, 1999)
1998	Evolution of hierarchies	(Mardling & Aarseth, 1999)
	Stumpff KS method	(Mikkola & Aarseth, 1998)
1999	HARP-6 procedures	(Aarseth, 1999a)
	Symplectic integrators	(Mikkola & Tanikawa, 1999b, 1999a)
	Nbody6++ SPMD / MPI acceleration	(Spurzem, 1999)
2000	Single stellar evolution - SSE	(Hurley, Pols, & Tout, 2000)
2002	Binary stellar evolution - BSE	(Hurley, Tout, & Pols, 2002)
2003	GRAPE-6 procedures	(Makino et al., 2003)
2006	2.5PN in Nbody5	(Kupi, Amaro-Seoane, & Spurzem, 2006)
2007	direct N -body GPU acceleration	(Portegies Zwart, Belleman, & Geldof, 2007)
2008	AR with Post-Newtonian terms	(Mikkola & Merritt, 2008)
2010	Updated AR for few-body problems	(Hellström & Mikkola, 2010)
2012	Nbody codes GPU acceleration	(Nitadori & Aarseth, 2012)
2013	MPI acceleration on GPU clusters / PHIGPU	(Berczik et al., 2013)
	3.5PN in Nbody6	(Brem, Amaro-Seoane, & Spurzem, 2013)
2015	SSE/AVX acceleration on GPU clusters	(Wang et al., 2015)
2017	Symplectic integrators (FSI)	(Dehnen & Hernandez, 2017)
2020	P^3T with SDAR in PeTAR	(Wang, Nitadori, & Makino, 2020a; Wang et al., 2020a)
2021	Minimum spanning tree MSTAR/BiFROST	(Rantala, Naab, & Springel, 2021)

From:
 Spurzem,
 Kamlah,
 Living
 Review
 Computational
 Astrophysics
 Vol. 9, pp. 3-109
 2023

Parameterized stellar evolution tracks (IFMR for neutron stars and white dwarfs) (SSE++/BSE++ from Kamlah et al. 2022 and Spurzem & Kamlah, Living Reviews Comp. Astroph. 2023)

- updated metallicity dependent core-collapse SNe, their remnant masses and fallback (Fryer et al. 2012; Banerjee et al. 2020),
- updated electron-capture supernovae (ECSNe), accretion-induced collapse (AIC) and merger-induced collapse (MIC) remnant masses and natal kicks (Nomoto 1984, 1987; Nomoto & Kondo 1991; Saio & Nomoto 1985, 2004; Kiel et al. 2008; Gessner & Janka 2018)
- (P)PISNe remnant masses (Belczynski et al. 2010, 2016; Woosley 2017),
- updated fallback-scaled natal kicks for NSs and BHs (Fuller et al. 2003; Scheck et al. 2004; Fryer 2004; Fryer & Kusenko 2006; Meakin & Arnett 2006, 2007; Fryer & Young 2007; Scheck et al. 2008; Fryer et al. 2012; Banerjee et al. 2020),
- and BH natal spins (see also Belczynski et al. (2020); Belczynski & Banerjee (2020)) from
 - Geneva model (Eggenberger et al. 2008; Ekström et al. 2012; Banerjee et al. 2020; Banerjee 2021b),
 - MESA model (Spruit 2002; Paxton et al. 2011, 2015; Banerjee et al. 2020; Banerjee 2021b),
 - and the Fuller model (Fuller & Ma 2019; Fuller et al. 2019; Banerjee et al. 2020; Banerjee 2021b).

ECSN = electron capture
Supernova

AIC = accretion induced
collapse

MIC = merger induced
Collapse

PISN = pair instability
Supernova

PPISN = pulsating PISN

NS = neutron star

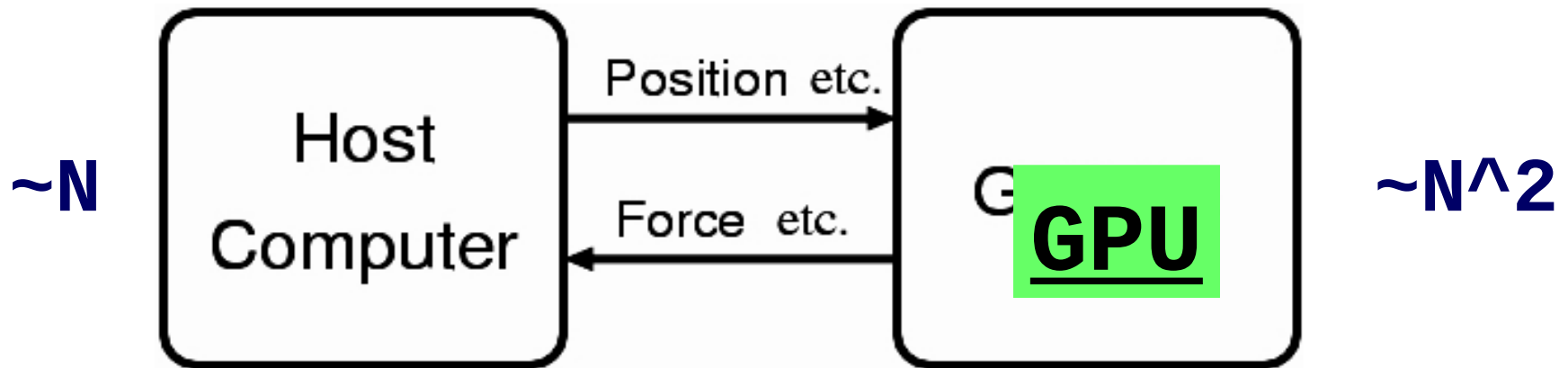
BH = black hole

MESA = recent stellar
evolution model

Codes

Parallelization
Supercomputing
GPU Computing

Our own ϕ GRAPE/GPU N-body code



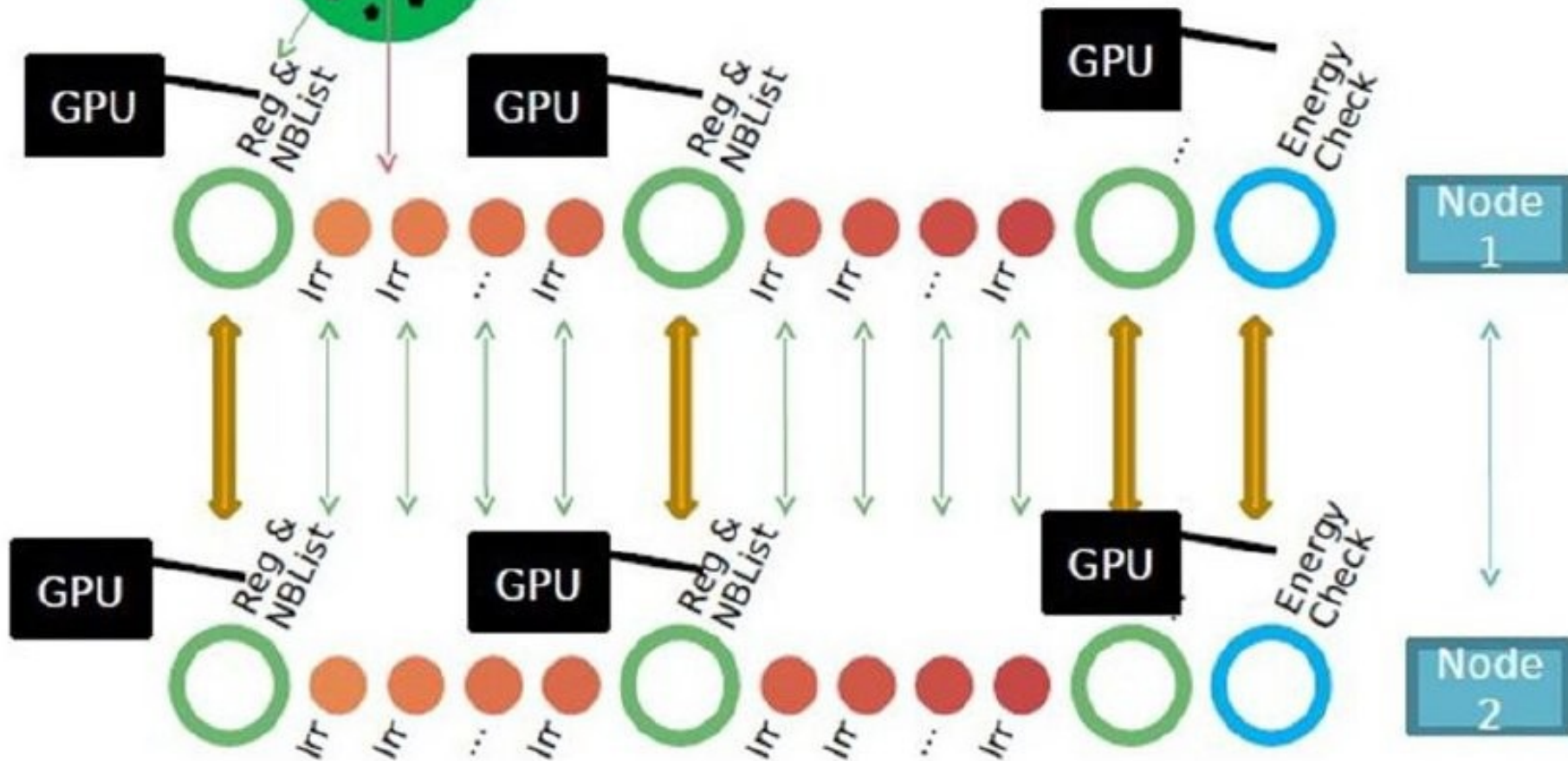
$$\vec{a}_i = \sum_{j=1; j \neq i}^N \vec{f}_{ij} \quad \vec{f}_{ij} = - \frac{G \cdot m_j}{(r_{ij}^2 + \varepsilon^2)^{3/2}} \vec{r}_{ij}$$

Our CPU/GPU N-body (AC) code

NBODY6++GPU

Wang, Spurzem, Aarseth, et al. 2015, 2016

→ Exaflop/s Huang, Berczik, Spurzem 2016



<https://github.com/lwang-astro/betanb6pp>

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Year	Keyword	Reference	
1961	Force polynomial	(Aarseth, 1963)	=
	Individual time steps	(Aarseth, 1963)	-
	Gravitational softening	(Aarseth, 1963)	
1966	Spherical harmonics	(Aarseth, 1967)	
1969	Two-body regularization	(Kustaanheimo & Stiefel, 1965)	
1972	Three-body regularization	(Aarseth & Zare, 1974)	
1973	Global regularization	(Heggie, 1974)	
	Neighbor scheme	(Ahmad & Cohen, 1973)	
1978	Co-moving coordinates	(Aarseth, 1979)	
1979	Regularized AC	(Aarseth, 1985)	
1980	Planetary formation	(Lecar & Aarseth, 1986)	
1986	Hierarchical block-time steps	(McMillan, 1986)	
1989	Chain regularization	(Mikkola & Aarseth, 1990)	
1990	Particle in box scheme	(Aarseth, Lin, & Palmer, 1993)	
1991	Collisional tree code	(McMillan & Aarseth, 1993)	
1992	Chain N -body interface	(Aarseth, 1994)	
1993	Hermite integration	(Makino, 1991)	
1995	Synthetic stellar evolution	(Tout et al., 1997)	
	Tidal circularization	(Mardling, 1995a, 1995b)	
	Slow chain regularization	(Mikkola & Aarseth, 1998)	
1996	Hierarchical stability	(Mardling & Aarseth, 1999)	
1998	Evolution of hierarchies	(Mardling & Aarseth, 1999)	
	Stumpff KS method	(Mikkola & Aarseth, 1998)	
1999	HARP-6 procedures	(Aarseth, 1999a)	
	Symplectic integrators	(Mikkola & Tanikawa, 1999b, 1999a)	
	Nbody6++ SPMD / MPI acceleration	(Spurzem, 1999)	
2000	Single stellar evolution - SSE	(Hurley, Pols, & Tout, 2000)	
2002	Binary stellar evolution - BSE	(Hurley, Tout, & Pols, 2002)	
2003	GRAPE-6 procedures	(Makino et al., 2003)	
2006	2.5PN in Nbody5	(Kupi, Amaro-Seoane, & Spurzem, 2006)	
2007	direct N -body GPU acceleration	(Portegies Zwart, Belleman, & Geldof, 2007)	
2008	AR with Post-Newtonian terms	(Mikkola & Merritt, 2008)	
2010	Updated AR for few-body problems	(Hellström & Mikkola, 2010)	
2012	Nbody codes GPU acceleration	(Nitadori & Aarseth, 2012)	
2013	MPI acceleration on GPU clusters / PhiGPU	(Berczik et al., 2013)	
	3.5PN in Nbody6	(Brem, Amaro-Seoane, & Spurzem, 2013)	
2015	SSE/AVX acceleration on GPU clusters	(Wang et al., 2015)	
2017	Symplectic integrators (FSI)	(Dehnen & Hernandez, 2017)	
2020	P^3T with SDAR in PeTAR	(Wang, Nitadori, & Makino, 2020a; Wang et al., 2020a)	
2021	Minimum spanning tree $\text{MSTAR}/\text{BiFROST}$	(Rantala, Naab, & Springel, 2021)	

From:
 Spurzem,
 Kamlah,
 Living
 Review
 Computational
 Astrophysics
 Vol. 9, pp. 3-109
 2023

Parameterized stellar evolution tracks (IFMR for neutron stars and white dwarfs)

(SSE++/BSE++ from Kamlah et al. 2022 and Spurzem & Kamlah, Living Reviews Comp. Astroph. 2023)

- updated metallicity dependent core-collapse SNe, their remnant masses and fallback (Fryer et al. 2012; Banerjee et al. 2020),
- updated electron-capture supernovae (ECSNe), accretion-induced collapse (AIC) and merger-induced collapse (MIC) remnant masses and natal kicks (Nomoto 1984, 1987; Nomoto & Kondo 1991; Saio & Nomoto 1985, 2004; Kiel et al. 2008; Gessner & Janka 2018)
- (P)PISNe remnant masses (Belczynski et al. 2010, 2016; Woosley 2017),
- updated fallback-scaled natal kicks for NSs and BHs (Fuller et al. 2003; Scheck et al. 2004; Fryer 2004; Fryer & Kusenko 2006; Meakin & Arnett 2006, 2007; Fryer & Young 2007; Scheck et al. 2008; Fryer et al. 2012; Banerjee et al. 2020),
- and BH natal spins (see also Belczynski et al. (2020); Belczynski & Banerjee (2020)) from
 - Geneva model (Eggenberger et al. 2008; Ekström et al. 2012; Banerjee et al. 2020; Banerjee 2021b),
 - MESA model (Spirou 2002; Paxton et al. 2011, 2015; Banerjee et al. 2020; Banerjee 2021b),
 - and the Fuller model (Fuller & Ma 2019; Fuller et al. 2019; Banerjee et al. 2020; Banerjee 2021b).

ECSN = electron capture
Supernova

AIC = accretion induced
collapse

MIC = merger induced
Collapse

PISN = pair instability
Supernova

PPISN = pulsating PISN

NS = neutron star

BH = black hole

MESA = recent stellar
evolution model

PeTar: a high-performance N -body code for modeling massive collisional stellar systems

Long Wang,^{1,2★} Masaki Iwasawa,^{2,3} Keigo Nitadori² and Junichiro Makino^{2,4}

¹Department of Astronomy, School of Science, The University of Tokyo, 7-3-1 Hongo, Bunkyo-ku, Tokyo, 113-0033, Japan

²RIKEN Center for Computational Science, 7-1-26 Minatojima-minami-machi, Chuo-ku, Kobe, Hyogo 650-0047, Japan

³National Institute of Technology, Matsue College, 14-4, Nishi-ikuma-cho, Matsue, Shimane 690-8518, Japan

⁴Graduate School of Science, Kobe University, 1-1 Rokko, Nada-ku, Kobe, Hyogo 657-8501, Japan

Slide 50

Accepted XXX. Received YYY; in original form ZZZ



New competition 1:
PeTar: MNRAS 2020
Many more papers...

ABSTRACT

The numerical simulations of massive collisional stellar systems, such as globular clusters (GCs), are very time-consuming. Until now, only a few realistic million-body simulations of GCs with a small fraction of binaries (5%) have been performed by using the NBODY6++GPU code. Such models took half a year computational time on a GPU based super-

FROST: a momentum-conserving CUDA implementation of a hierarchical fourth-order forward symplectic integrator

Antti Rantala^{1★}, Thorsten Naab¹, Volker Springel¹

¹Max-Planck-Institut für Astrophysik, Karl-Schwarzschild-Str. 1, D-85748, Garching, Germany

Slide 50

New competition 2:
FROST: MNRAS 2021
BIFROST: MNRAS 2022

Accepted XXX. Received YYY; in original form ZZZ

ABSTRACT

We present a novel hierarchical formulation of the fourth-order forward symplectic integrator and its numerical implementation in the GPU-accelerated direct-summation N -body code FROST. The new integrator is especially suitable for simulations with a large dynamical range due to its hierarchical nature. The strictly positive integrator sub-steps in a fourth-

NBODY6++GPU with up to 16M particles

(Benchmarks on raven at MPCDF)

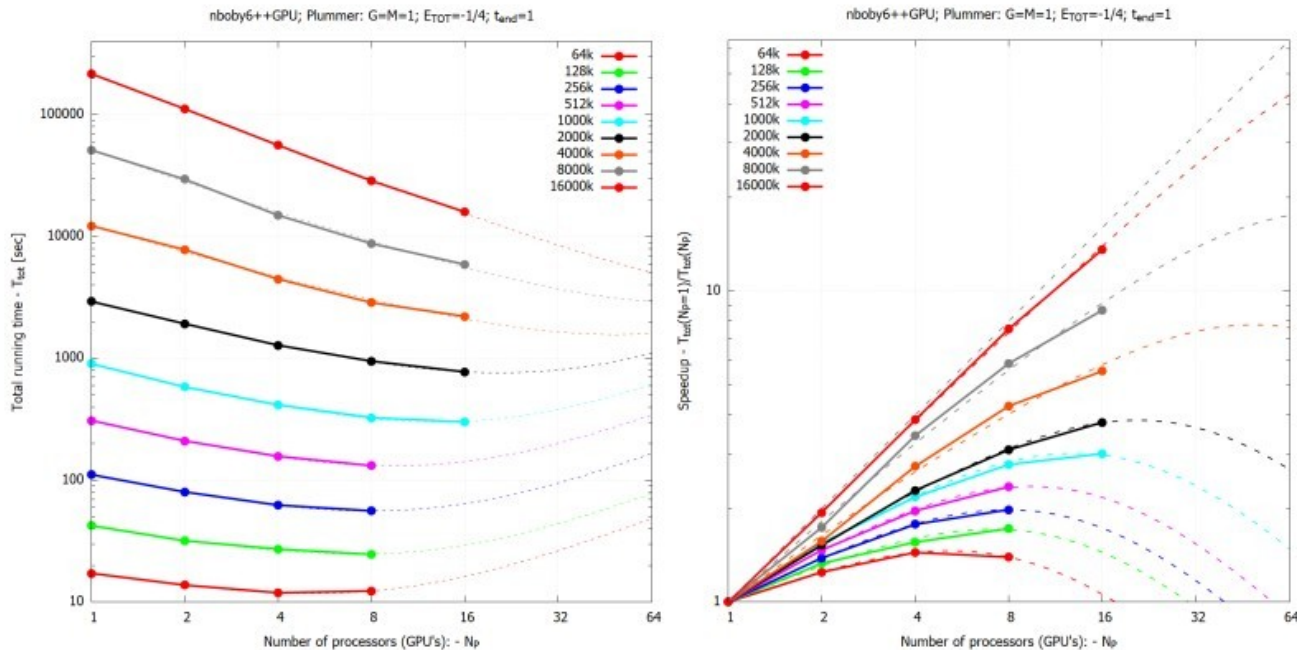


Figure 6: Benchmark results and extrapolated scaling for NBODY6++GPU, initial Plummer model, on the raven cluster at MPCDF, see main text. **Left:** Total time for one NBODY model unit in secs; **Right:** Speed-Up compared to using one GPU only. In both cases different curves for particle numbers from 64k to 16m. Ideal Speedup is the diagonal dashed lines, other dashed lines extrapolations from the timing model.

Supercomputers

A special-purpose computer for gravitational many-body problems

Nature 1990

Daiichiro Sugimoto^{*}, Yoshihiro Chikada[†], Junichiro Makino^{*}, Tomoyoshi Ito^{*}, Toshikazu Ebisuzaki^{*} & Masayuki Umemura[‡]

^{*} Department of Earth Science and Astronomy, College of Arts and Sciences, University of Tokyo, 3-8-1 Komaba, Meguro-ku, Tokyo 153, Japan

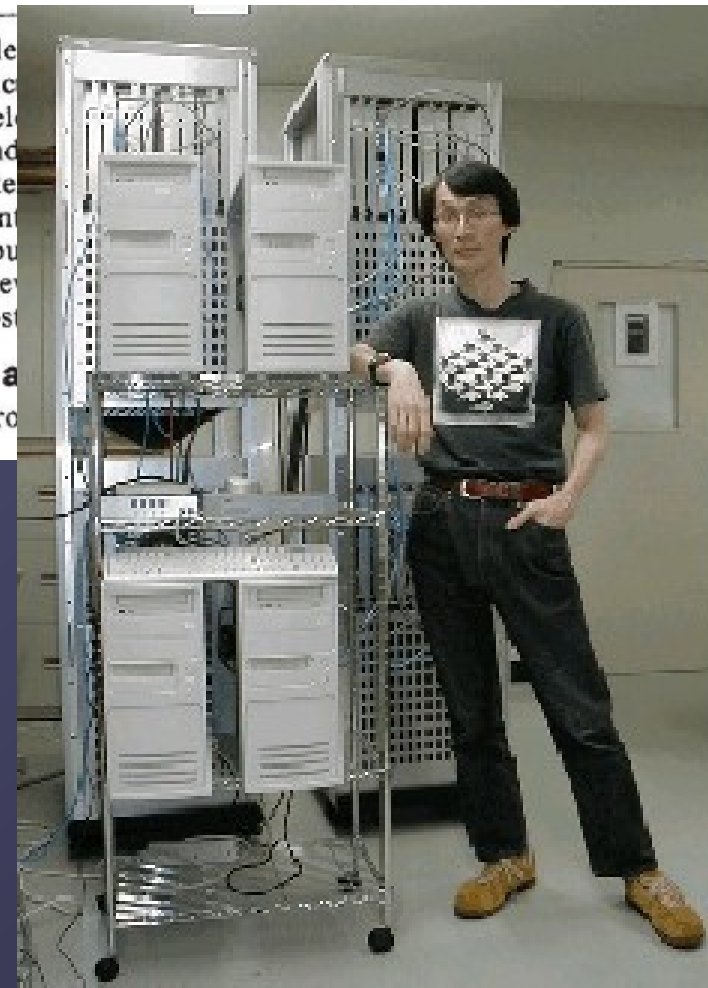
[†] Nobeyama Radio Observatory, Minamimaki-mura, Minamisaku-gun, Nagano 384-13, Japan

[‡] National Astronomical Observatory, Mitaka, Tokyo 181, Japan

A processor has been constructed using a 'pipeline' architecture to simulate many-body systems with long-range forces. It has a speed equivalent to 120 megaflops, and the architecture can be readily parallelized to make teraflop machines a feasible possibility. The machine can be adapted to study molecular dynamics, plasma dynamics and astrophysical hydrodynamics with only minor modifications.

(very-large-scale computer is constructed. The problem is to develop construction and of physical system to these apparent in the computer GRAPE-1 achieves XMP/1 at a cost

The pipeline architecture
The N -body problem



Jun Makino
with GRAPE6 →
(2002)

GRAPE-6 Gravity/Coulomb Part

- G6 Chip: 0.25μ 2MGate ASIC, 6 Pipelines
- at 90MHz, 31Gflops/chip
- 48Tflops full system (March 2002)
- Plan up to 72Tflops full system (in 2002)
- Installed in Cambridge, Marseille, Drexel, Amsterdam, New York (AMNH), Mitaka (NAO), Tokyo, etc..

NAOC laohu cluster 64 Kepler K20

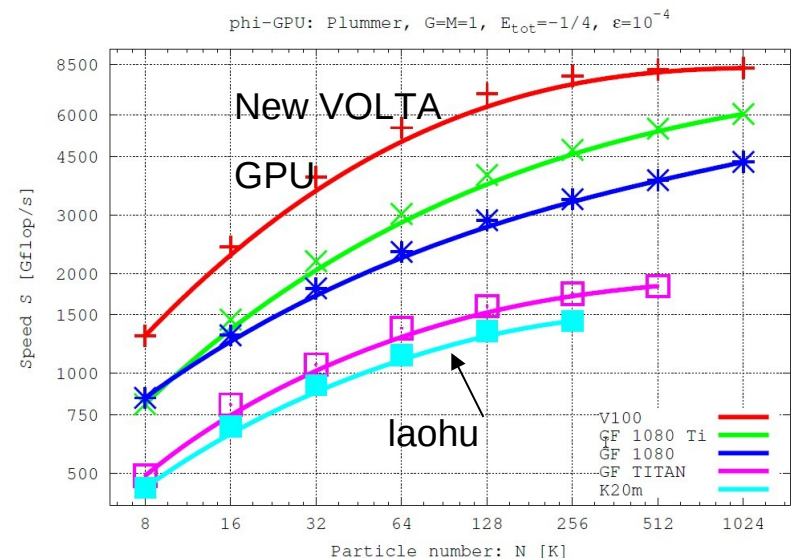


Laohu: 2009/2015

(Kepler GPU)

100 Tflop/s 150k cores

New GPUs 5-6 times
faster... (see below)



JUWELS Booster 936 nodes (AMD CPU, 4x Ampere GPU)

~450.000 AMD cores, 25 million NVIDIA Ampere GPU cores

~ 70 Pflop/s SP ~ 44 Pflop/s DP

No. 12 in top500 list No. 25 in green500 list

Jülich Wizard for European Leadership Science



LUMI

Supercomputer, Kajaani, Finland

Using only
Hydroelectric
Power and its
Heat used for heating
buildings.

No. 3 in top500

No. 7 in green500

2.2 million cores

10.000 AMD GPUs



EuroHPC and LUMI consortium:

Finland, Belgium, Czech Republic, Denmark, Estonia, Iceland, Norway, Poland, Sweden, and Switzerland.

China New Computing Technology

V100的全面国产替代 ——海光DCU Z100L 面向人工智能的GPGPU加速卡



简单介绍一下海光的这款产品

DCU (Deep Computing Unit) Z100L 是一款面向人工智能、科学计算的高性能全功能GPGPU加速卡。

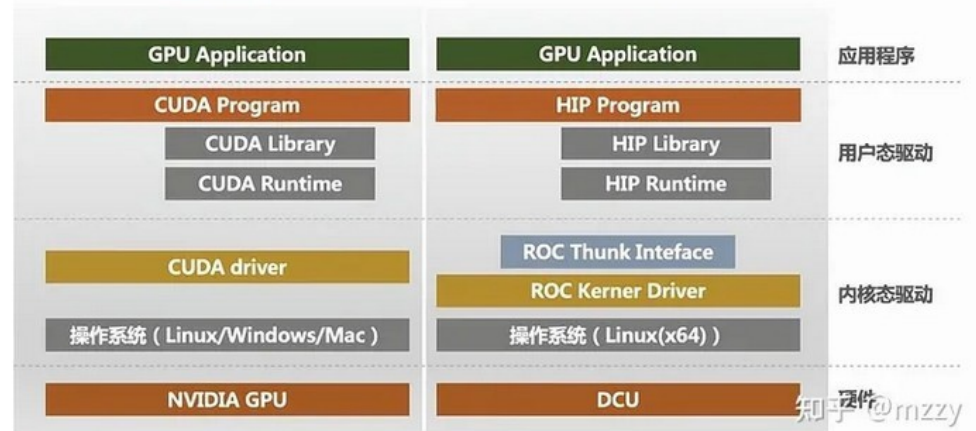
规格参数		Z100L及Z100L-LP	
应用场景		全精度支持	AI优化
适配服务器		2U/4U风冷服务器	
构架核心	计算核心	3840	
	频率	1319MHz	1600MHz
性能指标	FP64	10.1TFlops	——
	FP32	10.1TFlops	12.2TFlops
	FP16	20.2TFlops	24.5TFlops
	INT8	40.5TFlops	49.1TFlop
存储规格	显存容量	32GB HBM2	
	显存带宽	1024 GB/s	
接口类型	PCIe 接口	PCIe 4.0 x16	
	GPU 互联	xHMI 互联 184GB/s	
最大功耗	TDP	250W	280W
尺寸规格	尺寸	全高全长双宽	

Hygon CPU/DCU
(Deep Computing Unit)

~ AMD EPYC/AMD GPU
HIP programming language

Cuda

DTK



CUDA数学库	DTK数学库	数学库功能
cuda	miopen	深度学习基础数学库
cublas	rocblas/hipblas	基础矩阵运算数学库
nccl	rccl	通信库



FESENKOV
ASTROPHYSICAL
INSTITUTE

China-Aid HPC Solution for FAI

China – Kazakhstan government funded science collaboration (several disciplines)

Astrophysics: Silk Road NAOC with FAI

50 GPU Servers (~400 GPUs), Storage, Network, ...

Use Chinese Hardware

Racks Layout (AC-units are not shown)

16x - racks

50x - GPU-servers

20x - 400G Switches

2x - storage servers

2x - management

servers

17 - UPSes



- 1) Star Cluster Simulations/
Black Holes/Grav. Waves
- 2) Nuclear Star Clusters
- 3) Code(s) and Hardware
- 4) **Summary and References**

DRAGON I Simulation: Wang, Spurzem, Aarseth, et al. 2015, 2016

DRAGON II Simulation: Arca Sedda, Kamlah, Spurzem, et al. 2023, 2024ab

DRAGON III Simulation: Wu et al. 2025, in prep., and DRAGON Data Release

<https://github.com/nbody6ppgpu>

The Future - DRAGON III 1m – 8m (16m?)

- Compare and validate, also with other methods (MOCCA, CMC, Fokker-Planck, ...)
- Globular Cluster (GC) next issues: Rotation, Populations, Tides and Tails, ...
- Nuclear Star Cluster (NSC) Simulations: TDE, EMRI, direct capture, partial TDE, partial accretion to SMBH.
- Will legacy codes survive with good support? (Ease of use, data releases, interfaces, new modules...)
- AI assisted coding, how far will it go?

the SILK ROAD PROJECT

丝绸之路计划

IN ANCIENT TIMES ... THERE WAS A BRIDGE
BETWEEN CULTURES AND CONTINENTS ...
... TODAY THERE IS A PROJECT
OF ASTROPHYSICS IN CHINA ON THE MOVE ...
... TO BUILD AN INTERNATIONAL BRIDGE
FOR COMPUTATIONAL SCIENCE

HOME RESEARCH ▾ PEOPLE ▾ SEMINARS ▾ CONFERENCES ▾ BOARD

A A A Login Register     

Silk Road Conference (former KCK - Kazakhstan, China, Korea, Uzbekistan Meeting)

14th Silk Road Conference

2nd Meeting of the Silk Road Astronomy Research and Education Alliance, Silk Road AREA



Exploring the Frontiers of Dynamical Astronomy with High-Performance Computing, Artificial Intelligence, and Leading-Edge Observational Techniques.

We continue the numbering of these conferences, which started as Korea-China meeting in 2009; later it became Kazakhstan-China-Korea meeting, and due to increasing interest and participation of colleagues from all over central and east Asia we decided now to rename it to Silk Road Conference.

New Uzbekistan University, Tashkent, Uzbekistan, May 18-22, 2026.

Last Conference: 13th Silk Road / 1st Silk Road AREA meeting

Scientific Organizing Committee*: <https://astro-silkroad.eu/conferences/silkroad14>

Zhongmu Li (Dali U), Rainer Spurzem (ARI/ZA, NAOC/CAS, KIAA/PKU)

Bobomurat Ahmedov (ITP, IAS, IFAR, UBAI, HAI)

Kanat Baigarim (ECL/NU), Chingis Omarov (FAI),

Ernazar Abdikamalov (ECL/NU), Bekdaulet Shukirgaliyev (ECL/NU)

Hyung Mok Lee (SNU)

Thijs (M.B.N.) Kouwenhoven (XJTLU)

Team & Collaborators as below and further:

Thorsten Naab (MPA), Mirek Giersz (CAMK),
Ataru Tanikawa (Tokyo U), Nadine Neumayer (MPIA)

DRAGON simulations – globular and nuclear star clusters

- **DRAGON simulation: PhD thesis Long Wang, KIAA/PKU**, awarded for first realistic globular cluster simulation using **NBODY6++GPU** with one million stars and many binaries (**Wang, Spurzem, Aarseth**, et al., MNRAS 2016).
- **The Dragon-II simulations – Paper III. Compact binary mergers in clusters with up to 1 million stars: mass, spin, eccentricity, merger rate and pair instability supernovae rate** (**Arca Sedda, M., Kamlah, A. W. H., Spurzem, R., et al.**) Paper I,II,III (MNRAS 2023, 2024)

NBODY6++GPU and more, current state:

- **Spurzem, R., Kamlah A.W.H. Direct N-body simulations, in Living Rev. in Comp. Astrophysics 9, id.3 (2023) (NBODY7 see also Banerjee, Sambaran papers)**

Direct Nuclear Star Cluster Models with SMBH and TDE:

- **DRAGON simulation of the Galactic Center**, PhD thesis of Taras Panamarev, ARI/ZAH Univ. of Heidelberg (**Panamarev, Just, Spurzem, Berczik, Wang, Arca Sedda**, MNRAS 2019), simple TDE
- **Revisit the Rate of Tidal Disruption Events: The Role of the Partial Tidal Disruption Event** (**Zhong, S., Li, S., Berczik, P., Spurzem, R.**) 933, 96 (2022), TDE improved 3
- **Marija Minzberg, Philip Cho**: Master Theses Heidelberg 2023, publication in progress, TDE improved 1 and 2, paper in preparation.

Recent or in prep.:

- Vergara et al. 2024, A&A, 2025, to be submitted A&A – Formation of IMBH
- Li, Zhong, Berczik, Spurzem, Chen, Liu (MNRAS 2023 and earlier): merging nuclei with TDE
- Amiri et al., fractal initial models, primordial mass segregation
- Kamlah, Wu, Liu, Tanikawa, Neumayer, Sp, et al., PopIII star clusters
- Cho, Wu, Sp et al., DRAGON-III of nuclear star clusters